## Dissertation

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# The progenitors of early-type dwarf galaxies in the Virgo cluster 

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## Zusammenfassung

In der vorliegenden Doktorarbeit untersuche ich die möglichen Vorgänger von elliptischen Zwerggalaxien im Virgo Galaxienhaufen. Zur Analyse der Galaxien des späten morphologischen Typs wurde ein Breitband-Datensatz in den $u$, $g$, $r$, i und $z$ Filtern des Sloan Digital Sky Survey benutzt. Die photometrischen und strukturellen Eigenschaften (z.B. Helligkeit und Radius) der einzelnen morphologischen Unterklassen wurden mit einem bestehenden Datensatz von elliptischen Zwerggalaxien verglichen, um mögliche evolutionäre Verbindungen zu untersuchen. Innerhalb des Datensatzes zeigen Galaxien vom Typ der blauen kompakten Zwerge - bei vorheriger Entfernung des Einflusses der Sternentstehungskomponente - eine bemerkenswerte Ähnlichkeit mit den elliptischen Zwerggalaxien. Diese Übereinstimmung bezüglich der photometrischen Eigenschaften könnte auf eine mögliche evolutionäre Verbindung hindeuten.

Die Resultate der Photometrie des gesamten Datensatzes wurden als Eingabewerte für einen Evolutionssynthese-Modelle benutzt, mit dem die ungestörte und gestörte zeitliche Entwicklung der photometrischen Parameter der Galaxien simuliert wurden. Durch den Vergleich der Modellergebnisse mit den Beobachtungsdaten heutiger elliptischer Zwerggalaxien bin ich in der Lage, die möglichen Vorgängertypen auf bestimmte morphologische Unterklassen einzugrenzen.
Innerhalb dieser Doktorarbeit konnte ich zeigen, dass neben den heutigen blauen kompakten Zwerggalaxien, auch die zukünftigen photometrischen Eigenschaften von späten Spiralgalaxien des Typs Sd und Sm sehr gut mit denen heutiger elliptischer Zwerggalaxien übereinstimmen.


#### Abstract

In the here presented thesis, I investigate the possible progenitors of elliptical dwarf galaxies in the Virgo galaxy cluster. For the analysis of the galaxies of the late morphological type, a broadband data set in the $\mathrm{u}, \mathrm{g}, \mathrm{r}, \mathrm{i}$ and z filters of the Sloan Digital Sky Survey was used. The photometric and structural properties (e.g. luminosity and radius) of the single morphological subclasses was compared with an existing data set of elliptical dwarf galaxies to explore the possible evolutionary connections. Within the sample, galaxies of the type blue compact dwarf - if the contribution of the starburst component is removed - show a remarkably good similarity to the elliptical dwarf galaxies. This agreement in the photometric properties could point to a possible evolutionary connection. The results of the photometry of the entire sample were used as input parameters for an evolutionary synthesis code, which simulates the undisturbed and disturbed evolution of the photometric parameters of the galaxies. By comparing the results of the simulation with the ones of the observations of today's elliptical dwarf galaxies, I am able to restrict the possible progenitors to certain morphological subclasses. In the course of this thesis, I was able to show that apart from blue compact dwarf galaxies, also the prospective photometric properties of late-type spirals of the type Sd and Sm are in good agreement with today's elliptical dwarf galaxies.


## Declaration

This thesis contains no material that has been accepted for the award of any other degree or diploma.

Hagen Thilo Meyer

## Contents

1 Introduction ..... 3
1.1 Galaxies ..... 3
1.2 Galaxy clusters and the transformation of galaxies ..... 8
1.3 The observation of galaxies ..... 11
1.4 Thesis outline ..... 13
2 BCDs ..... 15
2.1 Introduction ..... 15
2.2 Sample selection and data reduction ..... 20
2.2.1 Sample selection ..... 20
2.2.2 SDSS data ..... 21
2.3 Surface photometry and decomposition ..... 22
2.3.1 LSB-component: exponential model ..... 24
2.3.2 LSB-component: inner flattening ..... 25
2.3.3 LSB-component: outer tail ..... 26
2.3.4 Starburst component ..... 27
2.3.5 Determination of the apparent axis ratio ..... 28
2.3.6 Additional tests: comparison of SBPs for different methods ..... 29
2.4 Results ..... 29
2.4.1 Decomposition examples ..... 29
2.4.2 Comparison of integral photometric properties ..... 31
2.4.3 Colour-magnitude diagrams ..... 37
2.4.4 Sizes of BCDs ..... 38
2.4.5 Properties of the starburst regions ..... 41
2.4.6 Apparent axis ratio of the LSB-component ..... 43
2.4.7 Colour gradients within BCDs ..... 44
2.4.8 Where are galaxies with uncertain morphological classification? ..... 47
2.5 Discussion ..... 53
2.5.1 Difference between certain and possible cluster member BCDs ..... 53
2.5.2 Distance dependence of the results ..... 53
2.5.3 Comparison between BCDs and the inner components of dEs ..... 55
2.5.4 Comparison with other BCDs ..... 58
2.5.5 Evolutionary connection between different dwarf galaxies ..... 61
3 Late-types ..... 67
3.1 Introduction ..... 67
3.2 The sample and data reduction ..... 70
3.2.1 Sample selection ..... 70
3.2.2 Remarks on the distances ..... 72
3.2.3 Classification ..... 72
3.2.4 Data reduction ..... 74
3.2.5 Data analysis ..... 76
3.3 Results ..... 79
3.3.1 Distribution within the Virgo cluster ..... 79
3.3.2 Photometric properties ..... 84
3.3.3 Colour-colour diagrams ..... 95
3.3.4 The luminosity function ..... 99
3.4 Discussion ..... 103
3.4.1 Comparison of the different types ..... 103
3.4.2 Application of the GOLDMine distance ..... 107
3.4.3 Reclassification ..... 112
3.4.4 Difference between certain and possible cluster members ..... 118
3.4.5 Comparison with semi analytical models ..... 118
3.4.6 Comparison with the literature ..... 120
3.5 Summary and outlook ..... 123
4 GALEV ..... 125
4.1 Introduction ..... 125
4.2 Stellar population synthesis models ..... 127
4.2.1 Why GALEV? ..... 128
4.2.2 Input parameters ..... 129
4.2.3 Starbursts and truncations ..... 129
4.3 Fitting the observed SEDs with GALEV models ..... 132
4.3.1 GAZELLE ..... 132
4.3.2 The grid ..... 135
4.4 Results ..... 140
4.4.1 The output ..... 140
4.4.2 Number distribution of the galaxies among the model grid ..... 141
4.4.3 Physical properties of the morphological types ..... 142
4.4.4 Colour-colour diagram ..... 146
4.4.5 The time evolution of the CMDs ..... 152
4.4.6 The evolution of the surface brightness profiles ..... 167
4.5 Discussion ..... 174
4.5.1 Classification by models ..... 174
4.5.2 Models vs. observation ..... 174
4.5.3 The possible progenitors of dEs ..... 177
4.6 Outlook ..... 182
4.6.1 GALEV in the zoom ..... 182
4.6.2 New observation ..... 183
5 Conclusions and outlook ..... 187
5.1 The possible progenitors of early-type dwarf galaxies ..... 187
5.2 Outlook ..... 188
6 Appendix ..... 207
6.1 Abbreviations ..... 207
6.2 Structural parameters of the sample ..... 208
6.3 Results of GAZELLE ..... 218
6.4 Results of the GAZELLE surface brightness profiles ..... 261

## Chapter 1 Introduction

### 1.1 Galaxies

Since thousand of years humans are fascinated by the starry sky. Apart from the twinkling stars a milky band was apparent in the sky, which was therefore called the "Milky Way" or in grecian "Galaxy". With the emerge of the telescope in the 17th century it was Galileo Galilei who first pointed it to the nightly sky and discovered the famous moons of the planet Jupiter, which in turn revolutionised our understanding of the Universe. But he also pointed his telescope to the milky band in the sky and discovered that it actually consists of thousands and thousands of stars.
Centuries after the discoveries of Galilei our view of the Milky Way has changed from an accumulation of stars on a sphere to an disc-like object with a spacial extension. However, the real physical extent of the Galaxy was under debate until the beginning of the last century and found its highlight in "Great Debate". Within the Great Debate it was devotedly discussed whether objects like the Andromeda nebula are objects within our own Galaxy or if they are other galaxies far away. With the measurement of the distance of variable stars within the Andromeda nebula (M31) by Edwin Hubble (e.g. Hubble, 1922, 1929) it became clear that M31 is far away from our Galaxy and is indeed another galaxy in the Universe. Thus, there is not only our own Galaxy, which forms the Universe, but there are other objects like the Galaxy. Other examples of galaxies in our neighbourhood are the Small and Large Magellanic Clouds (SMC and LMC), which are even visible with the naked eye. Figure 1.1 shows images of M31, LMC and SMC at visible wavelength.

Thus, starting from this we have to ask what actually is a galaxy? A naive answer to the question would be that galaxies are a vast number of stars in a limited volume.
However, detailed studies over the last decades reveal a more complicated view


Figure 1.1: Optical images of the Andromeda Nebula (left), LMC (middle) and SMC (right) (Credit: R. Gendler (left and middle); J. Hambsch and R.Gendler (right)).
of galaxies. First of all, galaxies do not only contain stars, but huge amounts of other ingredients. One ingredient in galaxies is gas in different phases and different abundance. A part of these gas clouds represents the birth place of new born stars and are therefore called star forming regions. One of the prominent representative of a gas cloud in our Galaxy is the Orion Nebula, also visible with the naked eye. In galaxies we found atomic gas as well as molecular gas. As we will see in this thesis, gas plays an important role on how a galaxy is classified today and how it would evolve with time.
In the last century the movement of the stars in galaxies was investigated and the outcome of these studies provided a remarkable result. The stars of a galaxies do not move as they should in the sense of the classical Keplerian physics, where an object nearby the centre of gravity moves faster than an object at larger distance, given by the following formula:

$$
\begin{equation*}
\mathrm{V} \sim \frac{1}{\sqrt{R}} \tag{1.1}
\end{equation*}
$$

where V is the velocity of the object and R is the radius from the centre of gravity (e.g. Binney and Tremaine, 2008). Moreover, stars at the outer rim of some special kind of galaxy move almost with the same velocity as stars at intermediate distance (Rubin and Ford, 1970). To overcome this problem a large amount of invisible matter, which is only measurable through its gravitational influence, was introduced and was called "Dark Matter" (DM). As we will see, DM is one of the key parameters in the description of the Universe and the formation of galaxies.

Other ingredients of galaxies are for example dust and of course the stars. Dust has only a minor contribution to the total mass of a galaxy, however it has the ability to change the appearance of a galaxy due to absorption and scattering of the star light. The mass range of galaxies spans several orders of magnitudes from several million solar masses $\left(\mathrm{M}_{\odot}\right)$ up to $10^{12} \mathrm{M}_{\odot}$ for the most massive ones (Willman et al., 2006; Misgeld and Hilker, 2011). Since the total mass of a galaxy is not easy to measure, it is common to use the luminosity of the galaxies instead of the mass. In astronomy the luminosity is often expressed in terms of magnitude $m$, which is defined as

$$
\begin{equation*}
m=-2.5 \cdot \log \left(f / f_{0}\right) \tag{1.2}
\end{equation*}
$$

where $f$ is the measured flux of the object and $f_{0}$ is the flux of a reference object (e.g. the star Vega).
Looking at images of galaxies one immediately realises that galaxies do not form a uniform class of objects, but show a large variety in morphology from galaxies with a rather smooth appearance to galaxies with an obvious spiral structure. In the framework of the famous study by Hubble (1926) the galaxies were classified regarding their morphology, starting with smooth elliptical galaxies with different degrees of flattening to the prominent spiral galaxies (see Fig. 1.2). The class of elliptical galaxies is also called "early-type" galaxies (ETGs), which are characterised by relatively red optical colours and almost no sign of star formation. On the other hand there are the spirals, which are also called "late-type" galaxies. Galaxies in the transition between ellipticals and spirals are called lenticular galaxies (S0 galaxies) and are associated to the early-type galaxies, too. It has to be pointed out that the terms "early-types" and "late-types" do not give any hint to the stage of evolution of the galaxies, moreover the light of late-type galaxies is dominated by young massive stars, while the light of early-type galaxies is dominated by an old stellar population.
Hubble furthermore subdivided the spiral galaxies into the groups of spirals with and without a stellar bar, which is denoted by a capital " B " in the morphological class like SBa in Fig. 1.2. These two groups of spirals are then further classified regarding the pitch angle, which describes the opening angle of the spiral arms and the ratio of the size of the central component (the bulge) to the size of the surrounding disc. Apart from morphological differences to the ETGs, spiral galaxies also differ to ETGs in the ability to form new stars. The stellar disc of spiral galaxies is a location of enhanced star formation. As already mentioned, stars are formed in gas clouds, therefore the disc of spiral galaxies must contain gas to fuel the star formation. To determine the amount of stars formed within a unit time - the star formation rate (SFR) - different


Figure 1.2: The Hubble Tuning Fork. (Credit: NASA \& ESA (adapted by R. Kotulla))
star formation laws are used. The simplest law uses a power law correlation between the surface density of the gas and the SFR (Schmidt, 1959; Kennicutt, 1998). Although the light of the young and massive stars in the spiral arms dominates the total light of the galaxy at the visible wavelength, the contribution of these stars is less to the total mass of the galaxy. The majority of the total mass of the disc of spiral galaxies comes from an older stellar population.
It was soon realised that some galaxies do not belong to any of the above described classes of ellipticals and spirals, respectively. Among these are faint galaxies with an irregular shape, which are therefore classified as "irregular galaxies". With the improvement of observational techniques it became clear that large and bright elliptical and spiral galaxies of the Hubble sequence are a minority. Most of the galaxies, which are observed today, are faint dwarf galaxies, making the bulk of the galaxy population (Phillipps et al., 1998; Trentham and Hodgkin, 2002). This overabundance of dwarf galaxies makes the study of them even more interestingly and necessary.
The dwarf galaxy population on its own can also be subdivided in almost the same manner as the bright ellipticals and spirals. At first glance early-type dwarf galaxies (ETDGs) - like their massive counterparts - have a smooth appearance and relatively red colours. However, studies by e.g. Sandage and Binggeli (1984), Jerjen et al. (2000), Lisker et al. (2007) and others revealed several subclasses with disc- and spiralfeatures, central nuclei and central star formation, thus making our view of dwarf galaxies more complicated.

Another group of the dwarf galaxies are the late-type dwarfs. In contrast to ETDs these galaxies significantly form new stars and may show a high order of distortion. In the class of the late-type dwarf galaxies one finds the "dwarf irregular galaxies" (dIs) and the "Blue Compact Dwarfs" (BCDs). The particular class of "Blue Compact Dwarfs" (BCDs) was first described in the early work of e.g. Zwicky (1965) and Searle and Sargent (1972) and their physical properties are still under debate. First believed to be very young objects with primordial chemical composition, Loose (1985) showed that they are composed of two components. One component exhibits a phase of very strong star formation, producing very massive and young stars, while the other component consists of an old stellar population. Therefore the BCDs cannot be young galaxies, making them even more interesting due to their particularly low chemical abundance, which is originally more common for young galaxies in the early Universe (Izotov et al., 1999b).
The difference between dIs and BCDs is commonly given by a higher SFR, compact optical size and higher surface brightness of the BCDs compared to dI galaxies (Thuan and Martin, 1981). However, there is no well defined limit for the division, resulting in the presence of transition types. Furthermore, the problem of a not well defined limit is not only valid in case of late-type dwarfs, but between late-types and early-type dwarf galaxies as well (Dellenbusch et al., 2008).
In general, the magnitude is used to divide galaxies into dwarf and giant galaxies. The limit, where a galaxy is a dwarf or a giant, is relatively arbitrary set to e.g. $\mathrm{M}_{\mathrm{B}}=-18$ mag (e.g. Thuan and Martin, 1981; Mo et al., 2010) ${ }^{1}$, indicating the absolute magnitude in the blue filter at 4360 Å. However, this limiting magnitude holds some problems, since there are overlap regions between low-luminosity spirals and dwarf galaxies. Furthermore, when applying this limit one also has to care about the distance of the galaxy, which is needed to calculate the absolute magnitude. Thus, measuring a wrong distance, may shift a gigantic galaxy into the regime of the dwarfs, even though the physical properties may be very different.
This problem of defining the dwarf- and non-dwarf galaxies is still not solved and an actual matter of discussion among the scientific community.

[^0]
### 1.2 Galaxy clusters and the transformation of galaxies

As we saw in the last section, there is a wealth of galaxies with very different properties. Surveys of large areas in the sky found that galaxies are not homogeneous distributed in the sky, but that there are regions with a high galaxy density ( $\rho$ ) and regions with a lower $\rho$. The regions with a lower $\rho$ and low number of galaxies are called "the field", whereas regions with a high $\rho$ and high frequency of galaxies are called "clusters" and "groups". In a galaxy cluster there are thousands of galaxies in a volume of a few megaparsec ${ }^{2}$ and prominent examples are the Perseus-, Coma- and Virgo cluster. The difference between groups and clusters is not well defined.
As shown by many studies over the last decades, galaxy clusters do not just contain galaxies. Zwicky (1937) showed by the study of the movement of galaxies in the Coma cluster that galaxy clusters need additional non-visible mass to be stable. He was the first who introduced the term "Dark Matter". Apart from galaxies and DM one also finds gas in the cluster, which is called intra cluster medium (ICM). As we will see in this thesis, this gas plays an important role in the evolution of the galaxies in the dense environment of a cluster.
In this thesis I will focus on the Virgo galaxy cluster, which has a distance of 16.52 Mpc (Mei et al., 2007) and a mean velocity of about $1200 \mathrm{~km} / \mathrm{s}$ (Binggeli et al., 1987). The Virgo cluster is also the nearest galaxy cluster to our Galaxy, therefore detailed investigations and surveys exist (e.g. Sandage and Binggeli, 1984; Sandage et al., 1985; Binggeli et al., 1985; Côté et al., 2004; Fritz and Hevics Collaboration, 2011). In the early study of Binggeli et al. (1985) there were 2096 galaxies listed in the Virgo Cluster Catalog (VCC) and 388 of them are in the focus of my study.
It was first noted by Oemler (1974) and Dressler (1980) that the distribution of galaxy types in a galaxy cluster is not homogeneous. In the local Universe the fraction of early-type to late-type galaxies is lower in the inner regions of a galaxy cluster and vice versa for the outer region. This behaviour is called the morphology-density relation (Postman and Geller, 1984; Whitmore et al., 1993; Goto et al., 2003). Naturally the question arises why do we have this morphology-density relation and what drives the evolution of galaxies?
There are several mechanisms acting on a late-type galaxy, which are able to transform it. Some of them are internal others are external mechanisms and the influence depends on several parameters, for instance the mass and velocity of the galaxy and the density

[^1]of the ICM within the galaxy cluster.
One of the internal mechanisms, which is able to influence a galaxy, is caused by the natural evolution of the stars within a galaxy. Due to stellar evolution, the lifetime of stars is limited and depends on the initial mass and chemical composition of the single star. A star like the sun has a lifetime of about 10 Gyr on the main sequence, while the lifetime is getting shorter with increasing mass and/or decreasing fraction of heavy elements (Sparke and Gallagher, 2000). Massive stars with masses about 100 $\mathrm{M}_{\odot}$ will live only for several million years. In their final stage they will end up in a super nova (SN) explosion, releasing a huge amount of energy into the galaxy and the gas within. Due to the shallow gravitational potential, less massive galaxies are more effected by the feedback of the SNe than the massive ones. In extreme cases, parts the gas of the late-type galaxy can be "blown out" by the simultaneous SNe explosions in a star forming region, resulting a galactic wind (Dekel and Silk, 1986; Izotov et al., 1996; Heckman et al., 2001).
There are also external forces acting on late-type galaxies which fall into a cluster. When the galaxy enters the cluster, the ISM of the galaxy interacts with the hot ICM of the cluster, resulting in the removal of the galaxy's ISM (van Zee et al., 2004; Boselli et al., 2008). In case of the removal of the gas, which is located in the envelope of the galaxies, the star formation does not stop immediately but decreases slowly until the gas is consumed (Larson et al., 1980). This scenario is also called starvation and strangulation, respectively.

On the other hand, when almost all gas is removed from the in-falling late-type galaxy, the star formation will decrease on very short time scales. This process is called ram pressure stripping (RPS, Gunn and Gott, 1972; Farouki and Shapiro, 1981). Figure 1.3 shows an observational example for a galaxy in the Virgo cluster, which loses its gas due to RPS (Kenney et al., 2004). It is also visible in the figure that the stellar component is not influenced by the RPS event. However, the effect of ram pressure on spiral galaxies and their star formation is still debated. Gavazzi et al. (1995) and Fujita and Nagashima (1999) argued that ram pressure can also induce new star formation in galaxy, entering the cluster for the first time, by compressing the gas, which is left over in the galaxy.

Starvation and RPS are processes that occur in galaxy clusters, but they are not limited to the inner dense core of cluster. Tonnesen et al. (2007) showed that they also play an important role even up to a distance of a virial radius, which has a value of 1.5 Mpc in


Figure 1.3: Combined image (R-band and HI) of NGC-4522 taken from Kenney et al. (2004). The HI gas was observed at radio wavelength and is over-plotted via a contour map.
case of the Virgo cluster (McLaughlin, 1999).
Another process, which influences the evolution of galaxies in a cluster, is galaxy harassment (Farouki and Shapiro, 1981; Moore et al., 1996; Mastropietro et al., 2005). In the harassment scenario high-speed encounters between cluster galaxies occur and in contrast to RPS and starvation, also the stellar component of the galaxy are affected. As a result, disc galaxies will lose a fraction of their stars and the remaining stars will be heated up, which thickens the stellar disc.
In summary, all the above described processes act on late-type galaxies, which fall into a galaxy cluster and result in a morphological transformation. This transformation will change the present-day appearance of the late-type galaxies. But the question is, how will the morphology change - to objects looking like today's dwarf elliptical galaxies or something completely different? At present day the question is not fully answered and there are also arguments against the evolution from late-types to dwarf ellipticals due to differences in their chemical composition (e.g. Grebel et al., 2003).
Therefore, the questions about the evolution of galaxies in clusters or groups, but also in isolation, is still controversially discussed among the community and there is no commonly accepted consensus in this field. To bring the discussion forward and shed more light on the evolution of galaxies, it is crucial to investigate all the transformation effects, but also the initial conditions, which are set by the currently observed galaxies.

### 1.3 The observation of galaxies

To investigate the properties of galaxies and galaxy clusters, informations from the light distributions of the objects are necessary. The Sloan Digital Sky Survey (SDSS, Adelman-McCarthy et al., 2007) provides a wealth of digital informations of the Virgo galaxy cluster. The SDSS uses a 2.5 m telescope (left panel of Fig. 1.4) and is located at Apache Point Observatory in New Mexico, USA (York et al., 2000; Gunn et al., 2006). To collect the light of the galaxies, the SDSS uses a 120 mega pixel CCD camera with a field of view of 1.5 square degree. The camera consists of 30 chips assembled in 6 columns with 5 filters (u,g,r,i and z) per column (right panel of Fig.1.4) and a pixel scale of $0.396 \operatorname{arcsec}^{(G u n n}$ et al., 1998). The filter characteristics ${ }^{3}$ are shown in Fig. 1.5, and the informations of the central wavelength and the "Full Width Half Maximum" (FWHM) are summarised in Tab. 1.1.
The SDSS observed the astronomical objects in drift-scan-mode, which enables the observation in each single filter at almost the same time. Therefore, one obtains a precise photometry for the objects. The exposure time in each filter amounts to 54 s . Due to the observation of the galaxies with the same telescope and the same observational technique, we are in the excellent situation to investigate the structural parameters and the evolutionary connections of the late-type and early-type galaxies of the Virgo galaxy cluster based on a very homogeneous sample.


Figure 1.4: The 2.5 m telescope (left) and the CCD-camera of the SDSS (right; both images from www.SDSS.org).

[^2]

Figure 1.5: Filter characteristic of the SDSS camera.

Table 1.1: Filter characteristic of the SDSS camera.
Filter Central wavelength FWHM

|  | $[\AA \AA]$ | $[\AA \AA]$ |
| :---: | :---: | :---: |
| u | 3540 | 570 |
| g | 4770 | 1370 |
| r | 6230 | 1370 |
| i | 7630 | 1530 |
| z | 9130 | 950 |

### 1.4 Thesis outline

In Chapter 2 the structural properties of Blue Compact Dwarf galaxies (BCDs) are studied in more detail, by decomposing the contribution of the star burst region and the old stellar population of the underlying host galaxy. The results of this analysis will be used for the discussion of the possible evolutionary connections between star forming dwarfs and early-type dwarf galaxies.
In Chapter 3 the entire late-type population of the Virgo galaxy cluster is investigated and the structural properties are derived. The results of the different morphological types will be compared with each other in the derived parameter-space and, if necessary, single galaxies will be reclassified to a new morphological class.
The results of the late-type galaxies will be the foundation of Chapter 4, where the structural parameters will be used to model the time evolution of the galaxies by the means of evolutionary synthesis models of GALEV. The outcome of these analyses will shed more light on the long debated possible evolutionary links between late-type and early-type galaxies.

# Chapter 2 What will blue compact dwarf galaxies evolve into? 


#### Abstract

Blue compact dwarf galaxies are objects with particular interesting properties due to their compact optical sizes, extreme star formation rates and low metallicities, which were therefore first believed to be young systems. With the improvement of the observational techniques, it was soon realised that there is an underlying component, consisting of old stars, which in turn raises the question whether these systems are really young. To investigate the underlying component of the blue compact dwarfs, the contribution of the starburst regions has to be subtracted from the total optical luminosity. The properties of this underlying component are in the focus of the following chapter. Furthermore, the derived results of the underlying components are compared to a sample of early-type galaxies within the Virgo cluster. This study sheds light on the possible evolutionary connections between the starbursting blue compact dwarfs and the "red and dead" early-type dwarf galaxies.


This study will be published together with Thorsten Lisker, Joachim Janz and Polychronis Papaderos.

### 2.1 Introduction

Galaxy clusters like the Virgo cluster are characterised by a large variety in galaxy morphology. This morphological variety depends on the local galaxy density, it therefore shows a clear trend with the cluster centric radius $\mathrm{R}_{\mathrm{CC}}$. The dependence of the morphology on $R_{C C}$ was first studied by Dressler (1980) and confirmed by many subsequent studies (e.g. Binggeli et al., 1990; Jerjen, 2012). At small $R_{C C}$ the dominant
galaxy types are early-type galaxies (ETGs): elliptical (E) and lenticular (S0), normal and nucleated dwarf ellipticals (dEs and $\mathrm{N}, \mathrm{dEs}$, respectively), dwarf lenticular (dS0), and dwarf spheroidal (dSph) galaxies, spanning a range from high to low luminosities. Traditionally, ETGs are associated with a smooth, regular appearance, with no signs of star formation. However, among the dEs is a particular class of galaxies with relatively blue cores, which indicate recent or still ongoing star formation at very low star formation rates (SFRs) (e.g. Lisker et al., 2006a).
Studies of ETGs showed that they are dominated by an old stellar population, resulting in relatively red colours. These properties indicated that ETGs are the oldest galaxy population in galaxy clusters. On the other hand, dynamically young galaxy clusters (e.g. the Virgo galaxy cluster) also contain late-type galaxies, which, in contrast to the ETGs, are located at larger $\mathrm{R}_{\mathrm{CC}}$. In its low-luminosity end this group is composed of star forming (SF) dwarf galaxies like blue compact dwarfs (BCDs) and dwarf irregulars (dIs) with properties strikingly different than those of ETGs. As shown by Vilchez (1995), the $\mathrm{H} \beta$ equivalent width (EW) of late-type galaxies in the Virgo cluster increases with increasing $\mathrm{R}_{\mathrm{CC}}$, echoing the strong impact of the cluster environment on galaxy evolution.
Tempting questions in this respect are i) how could this morphology-density relation be explained, ii) which mechanisms are responsible for the gradual transformation of late-type galaxies into ETGs as $\mathrm{R}_{\mathrm{CC}}$ decreases, and iii) are the descendants of today's late-type galaxies objects like ETGs in the Virgo cluster. In the literature several mechanisms have been proposed as drivers of galaxy transformations within the dense cluster environment, most notably ram pressure stripping (Gunn and Gott, 1972), tidal stirring (Mayer et al., 2001) and harassment (Moore et al., 1996; Mastropietro et al., 2005), all of which have in common the removal of gas (see e.g. Hensler, 2012, for a recent review).
Dwarf galaxies with a low central stellar density $\rho_{\star}$, such as dIs, are expected to be particularly prone to gas removal as they plunge into the hot intracluster medium (ICM), whereas high $-\rho_{\star}$ systems such as BCDs (cf. e.g. Papaderos et al., 1996b) might be able to retain some fraction of their gaseous reservoir down to a lower $\mathrm{R}_{\mathrm{CC}}$. But how would these dwarfish late-type galaxies in the cluster periphery look like after some billion years, once their gas content has been removed, in the course of one or several passages through the dense ICM core, and their ensuing long passive photometric evolution? Addressing this question is fundamental to the understanding of the morphological diversity of the dwarf galaxy population in clusters. Another
issue of special interest is, what impact has the initial contact of late-type dwarfs with the cluster periphery on their SF activity and whether, for certain conditions, starbursts can be ignited, transforming them into BCDs. If so, then how do these cluster-BCDs differ from the main population of field BCDs in their recent star formation history (SFH) and the morphological and metric properties of their SF component?
BCDs are low-luminosity galaxies ( $\mathrm{M}_{\mathrm{B}}>-18 \mathrm{mag}$ ) with a compact optical appearance and blue integral colours (Thuan and Martin, 1981). Many studies over the past decades have shown that BCDs are metal deficient with a median oxygen abundance $12+\log (\mathrm{O} / \mathrm{H}) \sim 8.0$ (e.g. Kunth and Östlin, 2000) and a small percentage of systems with a gas-phase metallicity as low as $7.0 \lesssim 12+\log (\mathrm{O} / \mathrm{H}) \lesssim 7.6$ (Searle and Sargent, 1972; Izotov et al., 1999a; Kunth and Östlin, 2000; Kniazev et al., 2004; Papaderos et al., 2008). These systems (XBCDs) are therefore the best nearby analogues of young low-mass galaxies in the early Universe (see e.g. the discussion in Papaderos et al., 2008). Apart from their low metallicity, BCDs also exhibit strong bursts of star formation, which are fed by a relatively large amount of gas (Thuan and Martin, 1981; Staveley-Smith et al., 1992; van Zee et al., 1998). Various arguments suggest, in line with evolutionary synthesis models, that starbursts in BCDs do not last longer than a few $10^{7}$ yr (Thuan, 1991; Krüger et al., 1995; Mas-Hesse and Kunth, 1999; Thornley et al., 2000) and have to be separated by long ( $\sim 1$ Gyr) quiescent phases. Dwarf irregulars, on the other hand, are characterised by prolonged low-level star formation over time scales of 450 Myr up to 1.3 Gyr , as shown by, e.g., McQuinn et al. (2010b, see also Skillman et al. (2012) for a review) through colour-magnitude diagram (CMD) studies of 20 of these systems.
The seminal study by Loose and Thuan (1985) has shown that BCDs are composed of two main stellar components. The first one displays the region of the ongoing star formation where young OB-stars are rapidly formed, it therefore has a low M/L-ratio. Due to the simultaneous formation of new massive stars in the starburst and their death within several Myr, metal-enriched galactic winds resulting from multiple supernovae explosions are expected to influence the chemical evolution of BCDs. However, studies have shown that in most cases starburst-driven feedback is insufficient for the expulsion of the entire ISM from a BCD (e.g. Silich and Tenorio-Tagle, 1998; Ferrara and Tolstoy, 2000; Tajiri and Kamaya, 2002; Recchi and Hensler, 2006). The starburst component contributes, on average, $\sim 50 \%$ of the optical emission of a BCD and, in some cases, up to $90 \%$ (Papaderos et al., 1996b; Noeske, 1999;

Cairós et al., 2001b; Amorín et al., 2009). Quite importantly, in several XBCDs ${ }^{1}$ nebular emission has been determined to be extraordinarily intense (EW ~ 1600 - 2000 $\AA$ ) and to contribute $30-50 \%$ of their total optical luminosity (Papaderos et al., 1998, 2002), in line with theoretical predictions for young starbursts (e.g. Krüger et al., 1995).

The second component of BCDs is dominated by an old population of M and K stars. This component is referred to in the following as the host galaxy or low-surface brightness ( $L S B)^{2}$ component and is characterised by a high M/L-ratio. Various studies indicate that the LSB-component dominates the baryonic-, and in some cases even the virial mass of BCDs within their Holmberg radius ${ }^{3}$, it therefore has to have a significant influence on the gas collapse characteristics and the starburst activity in these systems (e.g. Papaderos et al., 1996a; Östlin et al., 1999; Lelli et al., 2012).

According to Loose and Thuan (1986) BCDs can be classified in four main classes, based on the morphology of their SF- and LSB-component. Table 2.1 summarises the different morphological types of BCDs. Nuclear-elliptical (nE) and irregular-elliptical (iE) BCDs, both characterised by an extended circular or elliptical LSB-component, dominate the BCD population ( $\sim 90 \%$ ). Irregular (iI) BCDs comprise about $10 \%$ of local BCD population. Interestingly, their cometary (iI,C) subclass (see e.g. Noeske et al., 2000, for two nearby examples) is remarkably common among XBCDs (Papaderos et al., 2008), a fact pointing to a connection between gas-phase metallicity, morphology and evolutionary status. This is also indirectly suggested by the large frequency of cometary galaxies (also referred as tadpoles) among comparatively unevolved high-z galaxies (e.g. Elmegreen and Elmegreen, 2010). Whereas, the galaxies VCC-0802 and VCC-0274 are the most characteristic examples of the iI,C class in the Virgo cluster, a closer inspection of the VCC catalogue reveals several more such candidates. It is probable that the cometary morphology in cluster-BCDs has a different origin than in more isolated systems. Cometary morphology in field XBCDs has been proposed to result from unidirectional sequential SF activity with a typical velocity of the sound speed in the warm ISM ( $\sim 20 \mathrm{~km} / \mathrm{sec}$, Papaderos et al., 1998,

[^3]2008). This scenario is supported by stellar age gradients along the "comets tail" (e.g. Guseva et al., 2003). In the cluster periphery, however, cometary morphology may arise from extranuclear star formation that is triggered through the interaction with the ICM (see Hensler, 2012, for a recent comprehensive review), as, e.g. the impressive cases of two SFDGs in Abell 1367 (Gavazzi et al., 2001) demonstrate.
Apart from the classification scheme of Loose and Thuan (1986), there are also other proposed schemes in the literature based on photometric and/or spectroscopic properties of the BCDs (e.g. Salzer et al. (1989), Sung et al. (2002), Telles et al. (1997) and Cairós et al. (2001a)). In the study presented here, will shall use the classification schema by Loose and Thuan (1986).
The evolution of star forming dwarf galaxies (SFDGs) like BCDs to "red and dead" early-type galaxies is still under debate and there is no satisfying answer right as yet. Early studies by Thuan (1985) and Davies and Phillipps (1988) introduced the possible evolution of dIs into BCDs in several bursts and finally, after reaching a higher metallicity and the depletion of gas, the fading into dEs. Thuan (1985) concluded that the metallicities of BCDs and dEs are very different and that only periods of 3-10 bursts over the Hubble time could be able to produce the metallicity range of dEs. The study of Marquart et al. (2007) on the BCD He 2-10 showed that the stars in this galaxy have random motions and show signatures of a merger. They concluded that, due to the velocity dispersion of the stars, in the future this BCD may evolve into a nucleated dE. In comparison to this, the simulations of Bekki (2008) showed that dwarf-dwarf mergers are able to produce BCDs, but the further evolution of these BCDs into dEs with no gas is ruled out due to the extended gas discs of the simulated BCDs.
In contrast to BCDs, the early-type dwarf galaxies (ETDGs) show almost no evidence for ongoing SF. However, several studies e.g. by Vigroux et al. (1984), Gu et al. (2006) or Lisker et al. (2006a) found dEs with blue colours (so called dE(bc)s) in their central region. The study by Lisker et al. (2006a) of 476 dEs in the Virgo cluster has disclosed 23 dEs with central blue colours, indicating recent or ongoing SF. Since the debate of the evolutionary connection of ETDGs and BCDs is still ongoing, it is worth to compare the structural properties of these two types of galaxies (BCDs vs. ETDGs) and in particular the properties of the $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$ (BCDs vs. $\mathrm{dE}(\mathrm{bc}) \mathrm{s})$.

The goal of this study is to shed light on two questions. First, whether BCDs in the periphery of galaxy clusters share similar structural properties with field BCDs and, secondly, whether the former can evolve into dEs after the cessation of their starburst. Neither question has been previously addressed in the literature.

Table 2.1: Different subtypes of BCDs according to the morphology of the starburstand LSB-component.

| Subtype | Description |
| :---: | :--- |
| nE | nuclear star forming region and |
|  | elliptical LSB-component $(\sim 20 \%)$ |
| iE | irregular star forming region(s) |
|  | and elliptical LSB-component $(\sim 70 \%)$ |
| iI | irregular star forming region(s) |
|  | and LSB-component $(\sim 10 \%)$ |
| iI,C | iI with cometary shape |
| i0 | no detected LSB-component (very rare) |

This paper is structured as follows: Section 2 presents our sample of cluster BCDs and the classification criteria used. In Section 3 we discuss the method used to derive surface brightness profiles (SBPs) and their decomposition into the luminosity contribution from the starburst and the underlying LSB-component. Several examples of our decomposition methodology and an initial discussion of the derived structural properties are given in Section 4, and in Section 5 we present a comparative study of the structural properties of our BCD sample with field BCDs, dIs and dEs. Our results and conclusions are summarised in Section 6.

### 2.2 Sample selection and data reduction

### 2.2.1 Sample selection

Our sample is based on the Virgo Cluster Catalog (VCC) by Binggeli et al. (1985), which includes galaxies of all types within the Virgo cluster area. Due to incomplete velocity information for the VCC galaxies, the VCC includes certain and possible cluster members (updated by Binggeli et al., 1993), apart from background galaxies. Since new velocities have become available in the meantime, largely due to the Sloan Digital Sky Survey (SDSS, Adelman-McCarthy et al., 2007), the membership was revised by one of us (T.Lisker, see appendix of Weinmann et al. 2011) using the NASA/IPAC Extragalactic Database (NED). If a galaxy is listed as certain or possible member in the VCC, but has a velocity above $3500 \mathrm{~km} / \mathrm{s}$, it is considered as a background galaxy. We do not change possible to certain members or vice versa.

For our study we take into account both certain and possible members to a magnitude limit of $\mathrm{m}_{\mathrm{B}} \leq 18.0 \mathrm{mag}$, to which the VCC was found to be complete (Binggeli et al., 1985). When applying a constant distance modulus of $\mathrm{m}-\mathrm{M}=31.09 \mathrm{mag}(d=16.5$ Mpc ; Mei et al., 2007) to all galaxies, this corresponds to an absolute magnitude limit of $\mathrm{M}_{\mathrm{B}}<-13.09 \mathrm{mag}$.
Within these limits, the VCC contains 57 galaxies which have the term "BCD" included in their morphological type, 38 of them with BCD as the only or primary class (Tab. 2.2.2). Among the galaxies listed as "unknown" in the VCC, a visual inspection of the SDSS images ${ }^{4}$ suggests a reclassification of VCC-0429 and VCC-1713 to the transition type "Im / BCD", and of VCC-1411 to a possible "BCD?", thus adding up to 60 galaxies, and 39 with BCD as primary class. Since the VCC photographic plates were more sensitive in the blue, it needs to be kept in mind that this could have partly influenced the classification, as any underlying red stellar population would appear less prominent in the blue. Therefore, we treat the 39 galaxies with primary class BCD as the primary working sample for our analysis. The remaining 21 galaxies with uncertain morphological classification, but still possibly being a BCD, are treated as a separate sample, which will be compared to the BCDs in Section 2.4.8.

### 2.2.2 SDSS data

To analyse the Virgo BCDs we have used the SDSS Data Release Five (DR 5) in u, $\mathrm{g}, \mathrm{r}, \mathrm{i}$ and z band with an effective exposure time of 54 s . Due to insufficient sky subtraction of the SDSS pipeline for nearby galaxies of large apparent size, we used the sky-subtracted images of Lisker et al. (2007), using DR5 data. All images were flux calibrated and corrected for Galactic extinction following Schlegel et al. (1998). Five of the 39 BCDs are not covered by the SDSS DR5 and one BCD is excluded because of a nearby other galaxy (VCC-1944). Figure 2.1 shows the distribution of the remaining 33 BCDs within the Virgo cluster. The position of M87 is marked with a black cross.
To avoid contamination by other sources than the galaxy itself, one has to remove these sources very carefully. This was done for the BCD sample by replacing for instance the flux of a star by the median flux of its environment.
For the analysis of the sample of $\mathrm{dEs}, \mathrm{dE}(\mathrm{bc}) \mathrm{s}$ and dIs we follow the method described in Lisker et al. (2007) and Janz and Lisker (2008). The parameters of the additional sample are measured within an elliptical aperture of two Petrosian radii (Petrosian,

[^4]Table 2.2: Galaxies that were initially classified according to their morphology as "BCD" or "uncertain BCD". In the VCC a ":" indicates a weak uncertainty in the morphological classification of the galaxy, whereas a "?" points to a strong uncertainty. Roman numerals are the luminosity class of the galaxy regarding the original catalogue by Sandage and Binggeli (1984).

| Amount | Type |
| :---: | :--- |
| 38 | "BCD", "BCD?", "BCD:","BCD or merger" |
| 10 | "Im / BCD", "Im III / BCD", "Im III / BCD:" |
|  | "Im III / BCD?", "Im III,pec / BCD" |
| 3 | "Spec / BCD", "Spec, N / BCD" " |
| 3 | "Sm III / BCD","SBm III / BCD" |
| 1 | "dS? / BCD?" |
| 1 | "Sd / BCD?" |
| 1 | "dS0 or BCD" |
| 2 | "unknown" $\rightarrow$ "Im / BCD" |
| 1 | "unknown" $\rightarrow$ "BCD?" |

1976) and objects, which do not belong to the galaxy were masked out. The effective radius was measured via:

$$
\begin{equation*}
\mathrm{R}_{\mathrm{eff}}=\mathrm{a}_{\mathrm{hl}} \cdot \sqrt{(b / a)}, \tag{2.1}
\end{equation*}
$$

with the semi-major axis half-light radius $\mathrm{a}_{\mathrm{hl}}$ and the axis radio (b/a).

### 2.3 Surface photometry and decomposition

BCDs consist of a starburst (SB) component embedded within a more extended stellar LSB-component of older stars (see Section 2.1). The luminosity contribution of the starburst amounts, on average, to $\sim 50 \%$ of the $B$ band emission, and in some extreme cases reaches up to $\sim 90 \%$ (Papaderos et al., 1996b; Noeske, 1999; Salzer and Norten, 1999; Amorín et al., 2009). On the other hand, the starburst component is almost negligible, in terms of its fraction by mass, as its $\mathrm{M} / \mathrm{L}$ ratio is several times lower than that of the LSB-component. An adequate separation of its emission via 1D or 2D decomposition is clearly necessary for isolating the emission and studying the structural properties of the host galaxy. Indeed, the composite SBP of a BCD holds prior to decomposition little insight into the photometric structure. For example, Papaderos et al. (1996b) have pointed out that the SBP of a nE BCD can closely


Figure 2.1: Map of the Virgo cluster. Red symbols refer to BCDs, big black dots to galaxies of the class "unknown", small black dots to early-type galaxies, green asterisks to $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$ and the black cross to M87.
resemble the $\mathrm{R}^{1 / 4}$ SBP of a massive elliptical, whereas a typical feature of the SBPs of iE BCDs is an extended plateau, which can be fitted by a Sérsic law with an exponent of $2.5 \lesssim \eta \lesssim 4$, and resembling the bulge of S 0 galaxies. Similarly the effective radius $\mathrm{R}_{\text {eff }}$, a quantity commonly used in studies of high- $z$ galaxies, can vary by up to a factor $\sim 3$, depending on the starburst luminosity fraction (cf. e.g. Papaderos et al., 2006).
For galaxies with a smooth appearance like elliptical galaxies the easiest way to obtain SBPs is to use elliptical apertures and sum up the enclosed flux. Due to the irregular morphology of BCDs, however, such an approach is impractical, on the one hand because the choice of the 'galaxy center' is subjective and wavelength-dependent (thus SBPs and radial colour profiles are not easily reproducible), and, on the other hand, because SBPs derived in this way show, in the case of iE- and iI-BCDs, discontinuous jumps. In some cases, the latter can significantly affect the intensity profile of the LSB-component, thereby biasing studies of the photometric structure of BCDs. 2D axis-symmetric models to BCDs other than those falling into the nE class also yield systematic residuals, unless one carefully masks out and avoids fitting of the SFcomponent, as was done in e.g. Amorín et al. (2009).

Therefore, we applied method iv) of Papaderos et al. (2002), which was also used by Noeske et al. (2006) under the name LAZY. This method, which has as input a set of co-aligned multi-band images of the same pixel scale (SDSS: $0.396 \mathrm{arcsec} / \mathrm{pix}$ ) and point spread function (psf), does not require a choicse for the 'centre' of a galaxy. In our study, an average FWHM of the psf of 4 pix ( 1.584 arcsec ) was applied to smooth the input images with a Gaussian convolution kernel. From these co-aligned images a $\mathrm{S} / \mathrm{N}$ weighted average image (called reference frame) is created, which is used to calculate a mask for each intensity interval $\Delta \mathrm{I}$ within the range $\mathrm{I}_{\text {min }}$ to $\mathrm{I}_{\text {max }}$. Pixels outside the intensity interval of $\mathrm{I}-\Delta \mathrm{I} \leq \mathrm{F} \leq \mathrm{I}$ are set to zero, while pixels within the interval are given full consideration (set to unity). In the next step, each mask is multiplied by all input images (ugriz-band) and the flux is measured within each irregular mask. The corresponding photometric radius $\mathrm{R}^{*}$ is calculated as:

$$
\begin{equation*}
\mathrm{R}^{*}=\left(\frac{\mathrm{A}_{\mathrm{I}}+\mathrm{A}_{\mathrm{I}-\Delta I}}{2 \pi}\right)^{1 / 2} \tag{2.2}
\end{equation*}
$$

In Equation 2.2, $\left(\mathrm{A}_{\mathrm{I}}+\mathrm{A}_{\mathrm{I}-\Delta \mathrm{I}}\right)$ are the areas with intensities above I and $\mathrm{I}-\Delta \mathrm{I}$. In the case of multiple SF regions the $\mathrm{R}^{*}$ derived in this way corresponds to the sum of their area, a concept which translates into a monotonous increase of radius with decreasing intensity threshold.
A validity check of Lazy can be found in Section 2.3.6, where the SBPs of dEs are compared for different methods of profile derivation.

### 2.3.1 LSB-component: exponential model

To distinguish the SB- from the LSB-component we use (g-i)-colour profiles. At smaller $\mathrm{R}^{*}$ the ( g - i -colour profile is due to the superposition of the contribution of the young SB population and the old population of the LSB-component, which results in relatively blue values. At larger R* the contribution of the SB-component vanishes and the old LSB-component dominates, which results in relatively constant red colours. The radius beyond which the colour index levels off to a red, nearly constant value is referred to as transition radius $\mathrm{R}_{\mathrm{tr}}$ (Papaderos et al., 1996b; Cairós et al., 2003; Noeske et al., 2003). We used this characteristic radius to define the minimum R* for fitting an exponential law of the form

$$
\begin{equation*}
\mathrm{I}_{\mathrm{LSB}}\left(\mathrm{R}^{*}\right)=\mathrm{I}_{\mathrm{LSB}, 0} \exp \left[-\left(\mathrm{R}^{*} / \alpha\right)\right], \tag{2.3}
\end{equation*}
$$

to the LSB-component. In units of mag arcsec ${ }^{-2}$ Equation 2.3 reads as

$$
\begin{equation*}
\mu_{\mathrm{LSB}}\left(\mathrm{R}^{*}\right)=\mu_{\mathrm{LSB}, 0}+1.086\left(\mathrm{R}^{*} / \alpha\right) \tag{2.4}
\end{equation*}
$$

where $\mathrm{I}_{\mathrm{LSB}, 0}$ and $\mu_{\mathrm{LSB}, 0}$ are the central intensity and central surface brightness of the LSB-component, and $\alpha$ is its exponential scale length.

### 2.3.2 LSB-component: inner flattening

In some cases the extrapolation of the exponential fit to $\mathrm{R}^{*}=0$ arcsec exceeds the intensity observed at intermediate and small radii. This implies that the exponential law is not applicable in the central part of the LSB-component and must flatten to a central surface brightness that is lower than the extrapolated value $\mu_{\mathrm{LSB}, 0}$ of the fit for $\mathrm{R}^{*} \geq \mathrm{R}_{\mathrm{tr}}$. This kind of perfectly exponential SBPs with a flat core were noticed by Binggeli and Cameron (1991) who have called them type V. Papaderos et al. (1996b) introduced a modified exponential law to approximate such profiles, motivated by two considerations: First, a Sérsic law with a shape parameter $\eta \lesssim 0.5$ cannot fit type V SBPs without producing systematic residuals (see Noeske et al., 2003, for a detailed discussion of this subject). Secondly, as shown by radiation transfer models by Papaderos et al. (1996b), a Sérsic law with $\eta \lesssim 0.5$ implies an extended 'hole' in the intrinsic luminosity density of a spheric-symmetric emitter, if radiation isotropy, and a uniform $\mathrm{M} / \mathrm{L}$ ratio and intrinsic extinction are assumed. As these authors considered the evacuation of the dwarf galaxy centres to be improbable, they introduced an intensity profile for which the intrinsic luminosity density increases monotonously with decreasing radius (for small flattening parameters at least; see below) and has a finite central value. This modified exponential fitting law (modexp) has the form

$$
\begin{equation*}
\mathrm{I}_{\mathrm{LSB}}\left(\mathrm{R}^{*}\right)=\mathrm{I}_{\mathrm{LSB}, 0} \exp \left(-\frac{\mathrm{R}^{*}}{\alpha}\right)\left[1-\mathrm{q} \cdot \exp \left(-\mathrm{P}_{3}\left(\mathrm{R}^{*}\right)\right)\right] \tag{2.5}
\end{equation*}
$$

where

$$
\begin{equation*}
\mathrm{P}_{3}\left(\mathrm{R}^{*}\right)=\left(\frac{\mathrm{R}^{*}}{\mathrm{~b} \alpha}\right)^{3}+\left(\frac{\mathrm{R}^{*}}{\alpha} \cdot \frac{1-\mathrm{q}}{\mathrm{q}}\right), \tag{2.6}
\end{equation*}
$$

with the typical ratio $\mathrm{b} / \mathrm{q}$ being of the order of three (Papaderos et al., 1999 and Papaderos and Östlin, 2011). b is a measure of the radial extent of the central core in units of the exponential scale length $\alpha$ and $q=\frac{\Delta I}{I_{\text {LSB }, 0}}$ describes the attenuation of the modexp fit, as compared to the central intensity $\mathrm{I}_{\mathrm{LSB}, 0}$ predicted by the pure exponential fitting law. A parameter of $(b, q)=(2.40,0.80)$ corresponds to a flattening of the central intensity for a pure exponential LSB-component of $20 \%$ and a core radius of $2.4 \alpha$. Due to the poor knowledge of the structure of the LSB-component in its central part it is not clear at the moment which parameter combination of (b,q) satisfactorily describes the original shape. Therefore, the choice of $(\mathrm{b}, \mathrm{q})$ is not a straightforward task
and has to rely on plausibility considerations (cf. Noeske et al., 2003). In this study a flattening toward smaller radii was only applied when $\mu_{\text {LSB }, 0}$ exceeds the LSB + SB central surface brightness $\mu_{\mathrm{tot}, 0}$, and the spectrum shows clear signs of star formation. For a detailed discussion about the advantage of using flattening formula instead of Sérsic fits we refer to the study of Noeske et al. (2003) and Cairós et al. (2003). By using a parameter of $(b, q)=(2.40,0.80)$ the resulting magnitude of the LSBcomponent is increased by 0.48 mag and $\mu_{0}$ is reduced by $1.747 \mathrm{mag} / \mathrm{arcsec}^{2}$. Since the magnitude difference of 0.48 mag is applied to all filter bands, the colours of the LSB-components are not affected.
The SBP of the galaxy VCC-0641 exemplifies the central flattening of the exponential LSB-component (see Fig. 2.5 and Section 2.4.1).
The mean effective surface brightness $\langle\mu\rangle_{\text {eff }}$ of the LSB-component was calculated by:

$$
\begin{equation*}
\langle\mu\rangle_{\mathrm{eff}, \mathrm{LSB}}=\mathrm{m}_{\mathrm{LSB}}+2.5 \log \left(2 \pi \mathrm{R}_{\mathrm{eff}, \mathrm{LSB}}^{2}\right) . \tag{2.7}
\end{equation*}
$$

The effective radius $\mathrm{R}_{\text {eff,LSB }}$ of the LSB-component was determined by the integration of the exponential or flattened SBP to the radius within which one half of the total flux is enclosed.

### 2.3.3 LSB-component: outer tail

The inspection of the SBPs of our BCD sample reveals that in most cases the exponential slope of the LSB-component shows a slight flattening for large radii, corresponding to very faint surface brightness levels ( $\mu \gtrsim 27 \mathrm{mag} / \operatorname{arcsec} 2$ ). This is illustrated in Fig. 2.2 for one galaxy, and is also clearly visible in Fig. 2.5 for VCC1744 for a radius $\mathrm{R}>15$ arcsec. We note that such a flattening has been reported in the outskirts of some BCDs, e.g. II Zw 71 and Mrk 178, and it contributes for no more than $\sim 3 \%$ of the total luminosity (Papaderos et al., 2001, 2002).
Due to this additional flux contribution from the outermost luminosity component, it is possible that the $\mathrm{R}_{\text {eff }}$ is slightly increased, even though this effect is likely marginal. To check the influence of the tail on the $\mathrm{R}_{\text {eff }}$, the LSB-component was integrated for an exponential slope until $\mathrm{R}_{\text {tail }}$ and then for $\mathrm{R}>\mathrm{R}_{\text {tail }}$ with the observed tail (the $\mathrm{LSB}+$ tailprofile). The radius where the observed SBP of the LSB-component deviates from a pure exponential slope is called tail-radius $\mathrm{R}_{\text {tail }}$. This tail-radius $\mathrm{R}_{\text {tail }}$ is located in Fig. 2.5 for VCC-1744 at a radius $\mathrm{R}>15$ arcsec. If necessary, a flattening of the SBP towards smaller radii was applied too.
In Fig. 2.2 one can see an example of an SBP with a tail.

From the integration of the LSB+tail-profile one can also obtain the Petrosian radius where the $\eta$-function ${ }^{5}$ reaches $\eta=0.2$ (Blanton et al., 2001; Yasuda et al., 2001). For reasons of homogeneity all parameters are derived within two Petrosian radii of the LSB+tail-profile. The tail of the BCDs may be caused by nearby bright stars or galaxies, which increase the local sky. LAZY, just like all surface photometry codes, is very sensitive on the quality of sky subtraction, and objective selection and removal of areas with enhanced sky level adjacent to a galaxy is difficult, and may introduce additional uncertainties. To overcome this problem the parameters of the LSB-components were derived within $\mathrm{R}_{\text {tail }}$ and subsequently corrected for the missing flux. The latter was estimated by integrating the SDSS images in each filter in the interval of $\left[\mathrm{R}_{\text {tail }}: \mathrm{R}_{2 \text { Petrosian }}\right]$ using elliptical apertures and the masking program SExtractor (Bertin and Arnouts, 1996). This analysis resulted in a correction value $\delta \mathrm{X}$ of the magnitude and $\mathrm{R}_{\mathrm{eff}}$, respectively. The final parameter X was corrected by the following equation:

$$
\begin{equation*}
\mathrm{X}_{\text {corrected }}=\mathrm{X}\left(\leq \mathrm{R}_{\text {tail }}\right)+\delta \mathrm{X} . \tag{2.8}
\end{equation*}
$$

Using this correction all photometrical and structural parameters were derived within two Petrosian radii (Petrosian, 1976; Graham and Driver, 2005) with $\eta\left(\mathrm{R}_{\text {Petrosian }}\right)=0.2$ using the combined SBP of the LSB-component. This combined LSB-profile accounts for additional flux at very low surface brightness level. For the diagrams and tables of this work, all value were corrected for the tail with the above described technique.

### 2.3.4 Starburst component

The SBPs of BCDs display at small photometric radii a luminosity excess that is due to emission from the SF regions. However, this central peak in SBPs does not necessarily imply that all SB regions are centrally confined to the BCD host, given that LAZY automatically ascribes the smallest $\mathrm{R}^{*}$ to the brightest region in a galaxy. Besides a central, nearly Gaussian luminosity peak, reflecting the emission of the brightest SB region (if available), the SBPs of BCDs display at small to intermediate radii the plateau feature mentioned in Section 3. The plateau can be approximated (see e.g. Papaderos et al., 1996b) by the Sérsic law (Sérsic, 1968):

$$
\begin{equation*}
\mathrm{I}_{\mathrm{P}}\left(\mathrm{R}^{*}\right)=\mathrm{I}_{\mathrm{P}, 0} \exp \left[-\left(\mathrm{R}^{*} / \beta\right)^{\eta}\right] \tag{2.9}
\end{equation*}
$$

with the scale length $\beta$ and the Sérsic index $\eta$.

[^5]

Figure 2.2: SBP of a BCD. Left panel: Line corresponds to an exponential fit to the LSB-component and the red data points show the outer tail in r-band. Right panel: same, but for the g,r and i-filters. Obvious are the outer tails of the profiles at low surface brightness levels. The horizontal bars indicate the radius at which the profile deviates from a pure exponential fit.

The corresponding surface brightness reads:

$$
\begin{equation*}
\mu_{\mathrm{P}}\left(\mathrm{R}^{*}\right)=\mu_{\mathrm{P}, 0}+1.086\left[\left(\mathrm{R}^{*} / \beta\right)^{\eta}\right] . \tag{2.10}
\end{equation*}
$$

Even in the case of a flatting of the SBP towards smaller photometric radii, a weak starburst is assumed and superposed on the LSB component.

### 2.3.5 Determination of the apparent axis ratio

To derive the apparent axis ratio $\mathrm{b} / \mathrm{a}=1-\epsilon$ - where a and b are the semi-major and semi-minor axis, respectively and $\epsilon$ is the ellipticity of the LSB-component - we used the IRAF task ellipse (Jedrzejewski, 1987). The axis ratio b/a was measured at a radius of one Petrosian radius (Petrosian, 1976) with an elliptical annulus. The

Petrosian radius was measured for the entire BCD without any decomposition into star burst and LSB-component. We choose the one Petrosian radius to avoid disturbances due to the star forming inner regions of the BCDs. The axis ratio was determined by the co-added optical gri-image of the BCD, without taking into account the flattening of the LSB-component. Since the flattening is only important for the inner part of the LSB-component the axis ratio at larger radii is not effected and therefore no special algorithm was applied.

### 2.3.6 Additional tests: comparison of SBPs for different methods

To check the validity of the code LAZY, we compare the results of LAZY with IRAF/ellipse for two dwarf elliptical galaxies. Since the light distribution in dEs is a smooth function, both methods should provide the same results. Figure 2.3 shows the comparison of both methods for two example dEs from the Virgo cluster (VCC0218 and VCC-0916). As one can see, both method indeed produce almost the same SBPs. To obtain the magnitude and the effective radius of the galaxies, the SBP were integrated over the entire range of radius. The results of the integration are shown in Tab. 2.3.

Table 2.3: Comparison between LAZY and IRAF/ellipse.

| VCC | $\mathrm{m}_{\text {ellipse }}$ <br> $[\mathrm{mag}]$ | $\mathrm{m}_{\text {LAZY }}$ <br> $[\mathrm{mag}]$ | $\mathrm{R}_{\text {eff,ellipse }}$ <br> [arcsec] $]$ | $R_{\text {eff,LAZY }}$ <br> [arcsec] |
| :---: | :---: | :---: | :---: | :---: |
| 0218 | 14.13 | 14.02 | 12.80 | 13.48 |
| 0916 | 14.81 | 14.83 | 5.33 | 5.249 |

Additionally, values of the total SBP, where no decomposition was applied, of the two BCDs are displayed in Tab. 2.4. The SBPs were integrated by using three different methods, namely a profile integration using the program python, LAZY and AMOR. For these two example galaxies, one can see that the results are in good agreement. However, LAZY seems to produces slightly larger effective radii.

### 2.4 Results

### 2.4.1 Decomposition examples

In this section we discuss in more detail our analysis methodology on the basis of three illustrative examples. The upper panels of Fig. 2.4 show the co-added gri-SDSS


Figure 2.3: SBPs of two dEs from the early-type sample of Janz and Lisker (2008); Janz and Lisker (2009), derived with LAZY (green) and IRAF/ellipse (red data points), respectively. Left: VCC-0219; Right: VCC-0916.

Table 2.4: Comparison between python-, LAZY- and Amor-integration. "expo" corresponds to an exponential fit to the LSB-component.

| VCC | $\mathrm{m}_{\text {python }}$ <br> $[\mathrm{mag}]$ | $\mathrm{m}_{\text {LAZY }}$ <br> $[\mathrm{mag}]$ | $\mathrm{m}_{\text {Amor }}$ <br> $[\mathrm{mag}]$ | $\mathrm{R}_{\text {eff,python }}$ <br> $[\operatorname{arcsec}]$ | $\mathrm{R}_{\text {efffLAZY }}$ <br> $[\operatorname{arcsec}]$ | $\mathrm{R}_{\text {eff,Amor }}$ <br> $[\operatorname{arcsec}]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0001 | 14.76 | 14.73 | 14.73 | 4.59 | 4.84 | 4.53 |
| 0024 | 14.82 | 14.80 | 14.88 | 4.49 | 4.77 | 4.23 |
| 0001 (expo) | 14.86 | 14.86 | - | 5.04 | 5.03 | - |
| 0024 (expo) | 15.08 | 15.08 | - | 5.30 | 5.28 | - |

images of these three galaxies. VCC-1744 (left) shows a very strong star forming region, which is off-centred. On the other hand, the star forming regions in VCC-0130 are spread over the entire galaxy, resulting in an irregular LSB-component. This is also visible in the contour map of the galaxy (middle right panel). VCC-0641 also shows star formation spread over a large fraction of its surface with one dominant SF knot, which is off-centered. The SDSS spectrum of VCC-1744 shows very strong emission lines, as typical for BCDs. VCC-0130 and VCC-0641 also show emission lines and a rising blue continuum in their spectra, even though these emission lines are not as strong as in the case of VCC-1744. Figure 2.5 shows the results of our decomposition analysis of these three BCDs. The upper panels of Fig. 2.5 show the SBPs in the $r$ band and the lower panels show the colour profiles. The colour profiles for the galaxies do not cover the same range in photometric radius as the SBPs because the $\mathrm{S} / \mathrm{N}$ in the outskirts of the LSB-component become lower, resulting in a large variation in
the colours. VCC-1744 on the left hand side of Fig. 2.5 is an example for a galaxy with an exponential LSB-component. By extrapolating its slope to $\mathrm{R}^{*}=0 \operatorname{arcsec}$ one obtains a central surface brightness of $\mu_{0, \mathrm{LSB}}=21.61 \mathrm{mag} / \mathrm{arcsec}^{2}$. At $\mu=$ $25 \mathrm{mag} / \mathrm{arcsec}^{2}$ the size of the SB- and LSB-component amounts $\mathrm{R}_{\mathrm{SB}_{25}}^{*}=2.4 \operatorname{arcsec}$ and $\mathrm{R}_{\mathrm{LSB}_{25}}^{*}=11.6$ arcsec, respectively. The transition radius - where the colours are getting roughly constant, or the slope of the colour-profile changes - was found to be $R_{\text {trans }} \approx 4$ arcsec. This can be seen in the colour profile in the lower left panel in Fig. 2.5. At smaller radii ( $\mathrm{R}^{*}<4$ arcsec) the colours became bluer, due to the strong star forming region. The SBPs of the galaxy VCC-0130 (middle) and VCC-0641 (left) of Fig. 2.5 show that a pure exponential approximation of the LSBcomponent would overestimate the central surface brightness of the LSB-component ( $\mu_{0, \mathrm{LSB}}$ ). Therefore, we assume a central flattening with a flattening parameter of ( $\mathrm{b}, \mathrm{q}$ ) $=(2.4,0.8)$ and obtain $\mu_{0, \text { LSB,flat }}=22.13 \mathrm{mag} / \operatorname{arcsec}^{2}$ for VCC-0130 and $\mu_{0, \text { LSB,flat }}=$ $21.95 \mathrm{mag} / \mathrm{arcsec}^{2}$ for VCC-0641. Furthermore, looking at the central upper panel of Fig. 2.4, one can see that the optical image of VCC-0130 shows several star forming regions, which further justify the application of the modexp fitting function.
The colour profile of VCC-0130 shows a very flat form with no hint of a strong star formation, but with the informations from the optical image and the spectrum it is evident that the BCD hosts significant SF activity, a fact again motivating the application of modexp. On the other hand, the colour profile of VCC-0641 shows redder colours at larger radii, indicating a change of the stellar population from its SF component towards its LSB-component. Another aspect, which is apparent from the colour profile of VCC-0641, is some ambiguity in the classification of BCDs: the starburst component is offset from the centre indicating an irregular starburst. The LSB-component has an elliptical shape, but with some distortion, thus it is possible to classify it both as iI or iE.

### 2.4.2 Comparison of integral photometric properties with those of the LSBcomponents

The structural and colour properties of the BCDs are illustrated in Fig. 2.6 for the entire galaxies, and in Figs. 2.7 through 2.9 for the LSB-components only. Table 6.2 summarises the results of our analysis. For every single BCD the results for the total BCD and the underlying LSB-component are given (column [5] to [12]). The membership (ms) to the Virgo cluster was adapted from Binggeli et al. (1985) and Binggeli et al. (1993) with an update by T. Lisker (see Lisker et al. (2006a)


Figure 2.4: Co-added gri-SDSS images (upper), contour-maps (middle) and spectra (lower) of VCC-1744 (left panel), VCC-0130 (middle panel) and VCC-0641 (right panel).
and references therein) with recent values from the literature (column [2]). The classification criteria of the LSB-component (column [3]) are given according to Tab. 2.1. Column [4] indicates whether a flattening towards smaller radii was applied (see Section 2.3 for more details). Column [14] shows the total absolute magnitude of the entire BCD when applying the distance given by the GOLDMine data base. The differences between a distance of 16.52 Mpc and the GOLDMine distance are discussed in detail in Section 2.5.2.


Figure 2.5: Upper panels: SBPs (up) and colour profiles (down) of galaxy VCC-1744 (left) and VCC-0130 (right) in the r-band. Lower panel: same for VCC-0641. In the case of VCC-1744 a pure exponential law was assumed to fit the LSB-component, while for VCC-0130 and VCC-0641 a profile flattening towards smaller radii was assumed.


Figure 2.6: The $\mathrm{M}_{\mathrm{r}}-\mathrm{r}_{\mathrm{eff}}-\langle\mu\rangle_{\mathrm{eff}}$-plane of BCDs and ETGs. No decomposition into starburst and LSB-component was applied. Red squares represent the BCDs. Black dots represent the sample of early-type galaxies taken from Janz and Lisker (2008); Janz and Lisker (2009), and green asterisks are dE(bc)s from Lisker et al. (2006a). The $2 \sigma$ deviations of the ETGs is displayed by the vertical bars. Blue squares are possible BCDs, but with an uncertain morphological classification.

Table 2.5: Derived structural parameters of the BCDs and their LSB-components.

| $\begin{aligned} & \hline \text { VCC } \\ & {[1]} \end{aligned}$ | ms <br> [2] | $\begin{gathered} \hline \text { LSB } \\ {[3]} \end{gathered}$ | Flattening <br> [4] | $\begin{gathered} (\mathrm{g}-\mathrm{i})_{\mathrm{tot}} \\ {[5]} \end{gathered}$ | $\begin{gathered} (\mathrm{g}-\mathrm{i})_{\mathrm{LSB}} \\ {[6]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}, \mathrm{r}} \\ {[7]} \end{gathered}$ | $\begin{gathered} \mathrm{M}_{\mathrm{LSB}, \mathrm{r}} \\ {[8]} \end{gathered}$ | $\overline{\mathrm{R}_{\mathrm{efff}, \mathrm{tot}, \mathrm{r}}[\mathrm{kpc}]}$ [9] | $\begin{gathered} \mathrm{R}_{\text {eff,LSB, },[\mathrm{kpc}]} \\ {[10]} \end{gathered}$ | $\begin{gathered} \hline\langle\mu\rangle_{\text {eff.tot,r }} \\ {[11]} \end{gathered}$ | $\begin{gathered} \langle\mu\rangle_{\text {eff,LSB, }, r} \\ {[12]} \\ \hline \end{gathered}$ | $\begin{gathered} \mathrm{D}[\mathrm{Mpc}] \\ {[13]} \end{gathered}$ | $\begin{gathered} \mathrm{M}_{\mathrm{tot}, \mathrm{r}} \\ {[14]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0001 | 2 | nE | 0 | 0.881 | 0.873 | -16.32 | -16.27 | 0.38 | 0.39 | 20.16 | 20.30 | 32.0 | -17.75 |
| 0010 | 1 | iE?? | 0 | 0.735 | 0.754 | -16.09 | -15.87 | 0.37 | 0.42 | 20.36 | 20.84 | 32.0 | -17.52 |
| 0022 | 1 | $n \mathrm{E}$ | 0 | 0.638 | 0.734 | -15.45 | -15.30 | 0.39 | 0.42 | 21.07 | 21.42 | 32.0 | -16.89 |
| 0024* | 2 | $n \mathrm{E}$ | 0 | 0.695 | 0.765 | -16.24 | -16.01 | 0.36 | 0.42 | 20.12 | 20.69 | 32.0 | -17.68 |
| 0074 | 1 | $n \mathrm{E}$ | 1 | 0.736 | 0.765 | -15.23 | -15.23 | 0.59 | 0.57 | 22.22 | 22.12 | 17.0 | -15.30 |
| 0130 | 1 | iI | 1 | 0.417 | 0.400 | -14.65 | -14.73 | 0.36 | 0.35 | 21.73 | 21.60 | 17.0 | -14.72 |
| 0144 | 2 | iE | 0 | 0.149 | 0.249 | -16.51 | -15.97 | 0.25 | 0.34 | 19.08 | 20.26 | 32.0 | -17.94 |
| 0172 | 1 | iI | 0 | 0.489 | 0.743 | -16.82 | -16.30 | 0.84 | 1.07 | 21.38 | 22.42 | 32.0 | -18.25 |
| 0207 | 1 | iI | 0 | 0.271 | 0.215 | -14.57 | -14.19 | 0.21 | 0.26 | 20.67 | 21.50 | 32.0 | -16.01 |
| 0223 | 1 | nE ? | 0 | 0.623 | 0.758 | -15.37 | -15.01 | 0.32 | 0.41 | 20.74 | 21.63 | 32.0 | -16.81 |
| 0274 | 1 | iI | 0 | 0.473 | 0.622 | -14.40 | -14.34 | 0.55 | 0.60 | 22.87 | 23.13 | 32.0 | -15.84 |
| 0324 | 1 | i? | 0 | 0.476 | 0.630 | -17.46 | -17.25 | 0.96 | 1.17 | 21.03 | 21.68 | 17.0 | -17.52 |
| 0334* | 2 | iE | 0 | 0.526 | 0.698 | -16.02 | -15.72 | 0.43 | 0.54 | 20.74 | 21.53 | 17.0 | -16.08 |
| 0340* | 1 | nE | 0 | 0.632 | 0.753 | -17.00 | -16.46 | 0.52 | 0.68 | 20.18 | 21.28 | 32.0 | -18.43 |
| 0410 | 2 | nI | 0 | 0.294 | 0.460 | -14.42 | -14.31 | 0.30 | 0.31 | 21.53 | 21.70 | 17.0 | -14.48 |
| 0428 | 2 | iI | 0 | -0.049 | 0.164 | -14.41 | -14.32 | 0.42 | 0.44 | 22.30 | 22.47 | 17.0 | -14.47 |
| 0429 | 1 | iI/nI | 0 | 0.630 | 0.883 | -15.16 | -14.95 | 0.79 | 0.98 | 22.92 | 23.60 | 17.0 | -15.23 |

continued

| $\begin{aligned} & \hline \text { VCC } \\ & {[1]} \\ & \hline \end{aligned}$ | ms <br> [2] | $\begin{gathered} \hline \text { LSB } \\ {[3]} \end{gathered}$ | Flattening <br> [4] | $\begin{gathered} (\mathrm{g}-\mathrm{i})_{\mathrm{tot}} \\ {[5]} \end{gathered}$ | $\begin{gathered} \hline(\mathrm{g}-\mathrm{i})_{\mathrm{LSB}} \\ {[6]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot,r}} \\ {[7]} \end{gathered}$ | $\mathrm{M}_{\mathrm{LSB}, \mathrm{r}}$ <br> [8] | $\mathrm{R}_{\mathrm{eff}, \text { tot, }[ }[\mathrm{kpc}]$ <br> [9] | $\begin{gathered} \hline \mathrm{R}_{\text {eff, LSB, },}[\mathrm{kpc}] \\ {[10]} \end{gathered}$ | $\overline{\langle\mu}\langle\mu\rangle_{\mathrm{eff}, \mathrm{tot,r}}$ <br> [11] | $\begin{gathered} \langle\mu\rangle_{\text {eff,LSB,r }} \\ {[12]} \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{D}[\mathrm{Mpc}] \\ {[13]} \\ \hline \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{M}_{\mathrm{tot}, \mathrm{r}} \\ {[14]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0459 | 2 | iI | 0 | 0.489 | 0.636 | -16.75 | -16.20 | 0.53 | 0.67 | 20.47 | 21.51 | 17.0 | -16.81 |
| 0513 | 1 | iI | 0 | 0.823 | 0.770 | -16.48 | -16.51 | 0.49 | 0.56 | 20.55 | 20.83 | 17.0 | -16.54 |
| 0562 | 2 | iI | 0 | 0.316 | 0.399 | -15.50 | -15.45 | 0.64 | 0.53 | 22.12 | 21.76 | 17.0 | -15.56 |
| 0641 | 2 | iI | 1 | 0.329 | 0.406 | -15.51 | -15.54 | 0.48 | 0.46 | 21.49 | 21.37 | 23.0 | -16.23 |
| 0802 | 2 | iI,C | 0 | 0.630 | 0.726 | -14.98 | -14.96 | 0.68 | 0.68 | 22.76 | 22.79 | 17.0 | -15.05 |
| 0841* | 2 | iE | 0 | 0.761 | 0.853 | -16.46 | -16.19 | 0.91 | 1.14 | 21.92 | 22.67 | 17.0 | -16.53 |
| 0890* | 1 | iE? | 0 | 0.531 | 0.819 | -15.07 | -14.81 | 0.31 | 0.39 | 20.98 | 21.72 | 23.0 | -15.79 |
| 1141 | 2 | nE | 0 | 0.733 | 0.738 | -15.50 | -15.34 | 0.40 | 0.46 | 21.08 | 21.57 | 23.0 | -16.22 |
| 1313 | 2 | iI | 0 | -0.263 | 0.058 | -14.50 | -13.66 | 0.18 | 0.26 | 20.33 | 22.03 | 17.0 | -14.56 |
| 1411 | 2 | iE | 0 | 0.650 | 0.768 | -16.19 | -16.20 | 0.98 | 1.09 | 22.34 | 22.58 | 17.0 | -16.25 |
| 1437* | 2 | nE | 0 | 0.715 | 0.801 | -16.74 | -16.40 | 0.37 | 0.50 | 19.69 | 20.70 | 17.0 | -16.80 |
| 1459 | 1 | iE? | 0 | 0.762 | 0.820 | -15.22 | -15.21 | 0.65 | 0.72 | 22.44 | 22.67 | 17.0 | -15.28 |
| 1572 | 1 | iE | 0 | 0.769 | 0.914 | -15.79 | -15.57 | 0.91 | 1.06 | 22.58 | 23.14 | 17.0 | -15.86 |
| 1713 | 1 | iI | 1 | 0.599 | 0.750 | -15.71 | -15.70 | 0.79 | 0.89 | 22.37 | 22.64 | 17.0 | -15.77 |
| 1744 | 2 | nI, C | 0 | 0.204 | 0.598 | -14.80 | -14.59 | 0.49 | 0.60 | 22.23 | 22.87 | 17.0 | -14.87 |
| 2015 | 1 | nE ? | 0 | 0.655 | 0.771 | -15.49 | -15.36 | 0.60 | 0.66 | 21.99 | 22.33 | 17.0 | -15.55 |
| 2033* | 2 | iE | 0 | 0.691 | 0.826 | -16.59 | -16.17 | 0.63 | 0.84 | 21.01 | 22.04 | 17.0 | -16.65 |

Derived structural parameters of the BCDs and their LSB-components.

Notes. [1]: VCC-numbers marked with "**" were also discussed in Lisker et al. (2006a); [2]: membership (ms) of the BCDs regarding the VCC; [3]: detailed classification of the LSB-component (see Tab. 2.1); [4] flattening: 1 = flattening was applied; [13]: GOLDMine distance; [14]: magnitude with GOLDMine distance.

### 2.4.3 Colour-magnitude diagrams

Figure 2.7 shows the colour-magnitude diagram (CMD) of the LSB-components (red squares) of our BCD sample. Galaxies with a profile flattening toward smaller radii are indicated with a black cross and the corresponding change in the magnitude is described by a vector. As mentioned in Section 2.3 the profile flattening parameters $(b, q)$ are uncertain, therefore, the vector is to be regarded as an aid to the eye, pointing to the locus of the diagram where the true values are expected to be. The criteria for the profile flattening were only fulfilled for three BCDs (VCC-0074, VCC-0130 and VCC-0641).
Blue squares show the results for galaxies with an uncertain morphological classification (see also Section 2.4.8). Additionally, we plot with black dots the CMD of ETGs within the Virgo cluster (Janz and Lisker, 2009). Horizontal bars indicate the $2 \sigma$ deviations of the early-types within a magnitude bin of $\Delta \mathrm{M}_{\mathrm{r}}=1 \mathrm{mag}$ (vertical length of the bars). The green asterisks correspond to dEs with a blue core in the centre, taken from Lisker et al. (2006a). On average the LSB-components of BCDs are still bluer compared to the ETG population in the Virgo cluster. The LSB-components of the BCDs show a large spread in (g-i)-colours of about 1.2 mag , in contrast to the dEs with a colour range of 0.96 mag. On the other hand, there is one extremely blue BCD (VCC-1313), which can be found at $(\mathrm{g}-\mathrm{i})_{\mathrm{LSB}} \approx 0.03 \mathrm{mag}$. Such an extremely blue $(\mathrm{g}-\mathrm{i})_{\text {LSB }}$-colour can naturally arise from extended nebular emission, as is the case for the XBCDs SBS 0335-052 E (Papaderos et al., 1998) and I Zw 18 (Papaderos et al., 2002). Extended nebular emission could diminish morphological asymmetries, and due to its nearly exponential drop-off for large radii, it can easily be mistaken for a stellar disc in distant, poorly resolved starburst galaxies (Papaderos et al., 2002; Papaderos and Östlin, 2012). Therefore, some caution is in order when the data of VCC-1313 are interpreted.
Looking at Fig. 2.7, one can see that the $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$ are slightly offset from the Virgo dEs on the CMD plot, showing bluer colours with an average ( $\mathrm{g}-\mathrm{i}$ )-colour of $\langle\mathrm{g}-\mathrm{i}\rangle=$ 0.81 mag. It looks as if $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$ are the brighter extension of the LSB-components of BCDs. There are two low luminous dE(bc)s (VCC-0674 and VCC-0901) with $M_{r}>-15$ mag. These two galaxies have a low $S / N$ and therefore their derived properties should be regarded with some caution.
Lisker et al. (2006a) included in their study an additional sample of galaxies that also includes BCDs of our here presented study. These galaxies from Lisker et al. (2006a) are marked in Tab. 6.2 with a "*". Due to the sample selection, all BCDs in the sample
of Lisker et al. (2006a) show regular elliptical isophotes, corresponding to the nE type. To gain a deeper insight on the structure of the BCDs we also classify the shape of the starburst region and of the LSB-component based on the classification schema of Loose and Thuan (1986). These subtypes of the BCDs are also shown in Tab. 6.2. Calculating the average (g-i)-colours of the different subtypes of BCDs one finds the tendency that BCDs with regular LSB-components ( nE or iE ) are redder than irregular shaped LSB-components (see Tab. 2.6). The right hand side of Fig. 2.7 shows the same CMD but additionally the Virgo dwarf irregulars (dIs) ${ }^{6}$ are plotted. No obvious separation between BCDs or dIs can be found, since they almost cover the same region in the CMD. However, the dIs tend to have redder colours at fainter magnitudes, with colours comparable to the ETGs or even redder. The upper right panel of Fig. 2.6 shows for comparison the same CMD, but there the total values of the entire BCDs are used.

Table 2.6: Average values of different dE- and LSB-types.

| Subtype | $\langle\mathrm{g}-\mathrm{i}\rangle$ | $\sigma$ | $\left\langle\mathrm{M}_{\mathrm{r}}\right\rangle$ | $\sigma$ | $\left\langle\mathrm{R}_{\text {eff,r }}\right\rangle[\mathrm{kpc}]$ | $\sigma$ | $\left\langle\left\rangle_{\text {eff,r }}\right\rangle\right.$ | $\sigma$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| dE(bc) | 0.814 | 0.084 | -16.81 | 0.93 | 1.10 | 0.34 | 21.89 | 0.88 |
| BCDs (total) | 0.535 | 0.243 | -15.68 | 0.82 | 0.53 | 0.22 | 21.36 | 0.98 |
| BCDs (LSB) | 0.647 | 0.217 | -15.49 | 0.80 | 0.61 | 0.26 | 21.84 | 0.90 |
| nE (LSB) | 0.760 | 0.051 | -15.64 | 0.55 | 0.55 | 0.14 | 21.56 | 0.79 |
| iE (LSB) | 0.711 | 0.213 | -15.81 | 0.49 | 0.68 | 0.33 | 21.70 | 1.17 |
| iI (LSB) | 0.521 | 0.264 | -15.14 | 0.89 | 0.60 | 0.26 | 22.13 | 0.79 |

Notes. The averages were determined for colours of $(\mathrm{g}-\mathrm{i})>0$ mag to avoid extremely blue BCDs, which could be contaminated by strong nebular emission. $\sigma$ corresponds to the standard deviation of the mean.

### 2.4.4 Sizes of BCDs

Figure 2.8 shows the effective radius $\mathrm{R}_{\text {eff }}$ vs. the mean effective surface brightness $\langle\mu\rangle_{\text {eff,r }}$ in r-band of the LSB-components and ETGs. The meaning of the symbols is the same as in Section 3.3.2. From Fig. 2.8 one can see that the LSB-components of the BCDs occupy a region at the edge of the $2 \sigma$ area of the ETGs. With decreasing $\langle\mu\rangle_{\mathrm{eff}, \mathrm{r}}$ the LSB-components tend to become more compact than the ETGs. When

[^6]

Figure 2.7: Left: CMD of the LSB-components of our BCDs and for comparison the early-type galaxies of the Virgo cluster are plotted, too. Red squares correspond to values of the LSB-components, where black crosses show LSB-components with inner profile flattening. The notation "LSB(corr)" indicates that the parameters of the LSB-components are corrected for the outer tail (see Section 2.3.3). The black vectors display the change of the parameters when an inner profile flatting is applied. Right: same CMD, but additionally plotted are the dIs (cyan triangles).
applying a profile flattening to the LSB-components regarding the criterion described in Section 2.3, one can see that the resulting values are well within the $2 \sigma$ area of the ETGs. Nevertheless, the LSB-components of the BCDs are still very compact compared to all dEs, but still they cover the same region as the most compact dEs. This supports also the early findings of Drinkwater and Hardy (1991) on a small sample of Virgo BCDs (also see Section 2.5.5). Table 2.6 summarises the average values of the different subtypes of dwarf galaxies. Since $n E-B C D s$ have similarities to the $d E(b c) s$ it could have been expected that their sizes are comparable. However, we find that the nE-BCDs are significantly more compact. In summary, even with the subdivision of the BCDs, the LSB-components tend to be more compact than the dEs. Regarding the average of $\langle\mu\rangle_{\text {eff,r }}$ there are no significant differences between the BCD subtypes.
The right hand side of Figs. 2.8 and 2.9 shows additionally the results of the dI galaxies. In both diagrams one finds a clear separation between the LSB-components and dIs


Figure 2.8: Effective radius $\mathrm{R}_{\text {eff }}$ vs. mean effective surface brightness $\langle\mu\rangle_{\text {eff,r }}$. Symbols are the same as in Fig. 2.7.
with a transition region, where both types can be found. Two BCDs (VCC-0274 and VCC-1572) are in the main region of the dIs. A detailed inspection of the gri-image and spectrum of VCC-1572 shows a diffuse galaxy and a spectrum without strong emission lines but a rising blue continuum. Therefore, VCC-1572 could be interpreted as a galaxy in a post-starburst phase, or simply as a misclassified dI. In contrast to VCC-1572, the galaxy VCC-0274 shows in the gri-image a clearly separated starburst region and the spectrum also shows strong emission lines. To summarise, at a given $\mathrm{R}_{\text {eff }}$ the LSB-components of BCDs tend to be brighter and have a higher $\langle\mu\rangle_{\text {eff,r }}$. As in Fig. 2.8, the LSB-components in Fig. 2.9 tend to be more compact at a given magnitude than the ETGs. Comparing the LSB-components and the $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$ one may conclude that $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$ are the linear extension of the LSB-components toward larger $\mathrm{R}_{\mathrm{eff}}$ and brighter magnitude. To illustrate this we overlay linear fits to the LSB-components (red-line) and dE(bc)s (green lines) in Fig. 2.9.
Figure 3.19 shows mean surface brightness $\langle\mu\rangle_{\mathrm{eff}, \mathrm{r}}$ as a function of the absolute magnitude $\mathrm{M}_{\mathrm{r}}$ for the LSB-components of the Virgo cluster on the left hand side and additionally plotted are the dIs on the right panel. The BCDs and the dIs are separated from each other in this diagram, but they share the same location as the dEs. There are seven LSB-components, which are outside the $2 \sigma$ deviation of the dEs.


Figure 2.9: Effective radius $\mathrm{R}_{\text {eff }}$ vs. absolute magnitude $\mathrm{M}_{\mathrm{r}}$ in r-band. Symbols are the same as in Fig. 2.7. Red dashed line shows a fit to the LSB-components and the green lines fits to the $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$ (line 1: fit without the two faint galaxies with $\mathrm{M}_{\mathrm{r}}>-15 \mathrm{mag}$; line 2: fit with the two faint galaxies).

### 2.4.5 Properties of the starburst regions

Figure 2.11 shows the concentration index (CI) of our BCD sample. The concentration index, according to the definition of Papaderos et al. (1996b) (hereafter P96b), describes the relative isophotal area of the SF-component compared to that of the LSBcomponent, both defined at $\mu(B)=25 \mathrm{mag} / \operatorname{arcsec}^{2}$. This definition differs from the commonly used logarithmic ratio of the radii encircling, e.g. $20 \%$ and $80 \%$ of the total luminosity, which primarily measures the (passband dependent) luminosity fraction from the SB. Especially for studies of higher-z starburst galaxies, the CI can be a useful supplement to the CAS quantitative morphology scheme, as it has the advantage of being less sensitive to k corrections and cosmological dimming. CI is defined as

$$
\begin{equation*}
\mathrm{CI}=1-\left(\frac{\mathrm{P}_{25}}{\mathrm{E}_{25}}\right)^{2}, \tag{2.11}
\end{equation*}
$$

where $P_{25}$ and $E_{25}$ denote the isophotal radius of the SB- and LSB-component at an extinction-corrected surface brightness level of $\mu_{B}=25 \mathrm{mag} / \mathrm{arcsec}^{2}$. P96b have used the CI to quantify the morphological variation from nE- towards iE-BCDs, which


Figure 2.10: Absolute magnitude $\mathrm{M}_{\mathrm{r}}$ vs. mean effective surface brightness $\langle\mu\rangle_{\mathrm{eff}, \mathrm{r}}$ in r-band. Symbols are the same as in Fig. 2.7.
show a high and low CI , respectively. They also pointed out that interacting BCDs systematically deviate from the CI vs. $\mathrm{M}_{\mathrm{B}}$ sequence delineated by relatively isolated ones.
The horizontal lines in the figure represent the different regimes of CI. CI $>0.8$ corresponds to a centrally concentrated starburst, as in the case for nE-BCDs, whereas iE-BCDs are typically characterised by a $\mathrm{CI} \leq 0.7$. In BCDs exhibiting a nearly galaxy-wide SF activity (i0 type), hence the LSB-component is almost invisible at the surface brightness cutoff of $\mu_{B}=25 \mathrm{mag} / \mathrm{arcsec}^{2}$, the CI drops to $<0.5$. Over their whole magnitude range, BCDs in the Virgo cluster tend to have a CI larger than $\mathrm{CI}>0.65$, with an average value of $\langle\mathrm{CI}\rangle=0.9(\sigma=0.06)$ without accounting for the inner flattening and $\langle\mathrm{CI}\rangle=0.8(\sigma=0.12)$ when a modexp fit is applied to their LSB-component. In contrast to the sample of field BCDs studied by Papaderos et al. (1996a) we do not find any trend of CI with absolute r-band magnitude, a fact which may provide a hint to the influence of the cluster environment on the spatial distribution of SF activities in dwarf galaxies. Looking at Fig. 2.11, one also realises that BCDs with central flattening tend to have a $\mathrm{CI}<0.8$. This finding is to be expected, given that an extended central core in the LSB-component would imply a larger luminosity fraction and spatial extent for the SF-component. Since the star formation in $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$
is located within the central region of the galaxy, the CI of the $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$ must have a value of $\mathrm{CI}>0.8$, thereby similar to the majority of the Virgo BCDs.


Figure 2.11: Absolute magnitude $\mathrm{M}_{\mathrm{r}}$ vs. concentration index CI in r-band. Symbols are the same as in Fig. 2.7.

### 2.4.6 Apparent axis ratio of the LSB-component

Figure 2.12 shows the number distribution of the axis ratio of the LSB-components in the r-band and additionally the distribution of Virgo dIs, dEs and dE(bc)s. The LSB-components have an average axis ratio of $\langle\mathrm{b} / \mathrm{a}\rangle_{\mathrm{LSB}}=0.64 \pm 0.2$. We found that the average axis ratio of our BCD sample is slightly smaller than that for dEs from Janz and Lisker (2009) (hereafter JL09) with an average value of $\langle\mathrm{b} / \mathrm{a}\rangle_{\mathrm{dE}}=0.70 \pm 0.18$, implying that BCDs are slightly flatter than dEs. This is in agreement with the results for BCDs by Sung et al. (1998), who found $\langle\mathrm{b} / \mathrm{a}\rangle_{\mathrm{BCD}, \mathrm{Sung}}=0.67$, which is even closer to the dE value. Interestingly, there is no difference between the axis ratio of the LSBcomponents and $\mathrm{dE}(\mathrm{bc}) \mathrm{s}\left(\langle\mathrm{b} / \mathrm{a}\rangle_{\mathrm{dE}(\mathrm{bc})}=0.63 \pm 0.2\right)$. In Fig. 2.13 the axis ratio is plotted versus the r -band magnitude of the LSB-components and compared against the dE sample of JL09. One can see that BCDs cover the same area like the dEs.


Figure 2.12: Number distribution of Virgo cluster dwarf galaxies. The dEs also include the $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$, which are additionally plotted in the lower panel. The upper panels show the number distributions of the axial ratios of the LSB-components and dIs.

### 2.4.7 Colour gradients within BCDs

A large gradient in the colour profile indicates a change in the stellar population of the galaxy. Studies by Clemens et al. (2011) and Peletier et al. (2012) on early-type dwarf galaxies in the infrared showed that they do not have constant colour profiles. Therefore, in respect to the possible evolutionary connection between BCDs and dEs, a constant colour profile is not peremptory (especially in case of the $\mathrm{dE}(\mathrm{bc})$ ). However, these studies observed in the IR, which makes a comparison to optical data difficult. In another study by Koleva et al. (2011) about gradients in early-type dwarf galaxies


Figure 2.13: Axial ratio b/a vs. absolute magnitude $\mathrm{M}_{\mathrm{r}}$ in r of the LSB-components and ETGs. Symbols are the same as in Fig. 2.7.
using long-slit spectroscopy, they pointed out that dEs with gradients are may be the progenitors of BCDs with low angular momentum. Table 2.8 shows the colour gradients of the LSB-components. The gradients were measured by the difference of the colours at a radius of $r<1 / 8 R_{\text {eff }}$ and at $r=2 R_{\text {eff }}$. The colour of the LSBcomponent at the inner radius of $r<1 / 8 R_{\text {eff }}$ - which is probably within the seeing of most of the cases - depends strongly on the fitting of the LSB-component. Therefore, the following analysis serves as a first-order approximation.
The average difference of the ( $\mathrm{g}-\mathrm{i}$ )-colours at $\mathrm{r}<1 / 8 \mathrm{R}_{\text {eff }}$ (following Peletier et al., 2012) and at $\mathrm{r}=2 \mathrm{R}_{\mathrm{eff}}$ of all LSB-components amounts for $\langle\Delta(\mathrm{g}-\mathrm{i})\rangle=0.26$ mag. To compare the colour gradients of the LSB-components to the ones of the dEs and $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$, additionally the colour at $\mathrm{r}=1 \mathrm{R}_{\text {eff }}$ was measured for all three classes. The dEs have an almost constant $(\mathrm{g}-\mathrm{i})$-colour profile $\langle\Delta(\mathrm{g}-\mathrm{i})\rangle_{\mathrm{dE}}=0.01 \mathrm{mag}$, while the $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$ (due to the blue colour excess in the centre) show a slightly gradient $\langle\Delta(\mathrm{g}-\mathrm{i})\rangle_{\mathrm{dE}}=0.03$ mag. The LSB-components have an average gradient for $\mathrm{r}=1 \mathrm{R}_{\text {eff }}$ to $\mathrm{r}=2 \mathrm{R}_{\text {eff }}$ of $\langle\Delta(\mathrm{g}-\mathrm{i})\rangle_{\text {LSB }}=0.14 \mathrm{mag}$, therefore, the $(\mathrm{g}-\mathrm{i})$-colour profile clearly differs from the early-type galaxies. However, for simplification the flattening of the SBPs of the BCDs was not taken into account, what would slightly decrease
$\langle\Delta(\mathrm{g}-\mathrm{i})\rangle_{\text {LSB }}$. To obtain an almost flat ( $\mathrm{g}-\mathrm{i}$ )-colour profile for the LSB-components of the BCDs, the composition of the stellar population has to change. The change in the colour profile would also influence the effective radius of the LSB-components in the way that it became larger, bringing them even more into the regime of the dEs.
But how is the effective radius influenced in detail when the colour gradients in the BCDs almost vanish? To investigate the influence on $\mathrm{R}_{\text {eff }}$, the BCD VCC-0274 was used as an example. Since the r-band was used as the reference band in our study, the ( $\mathrm{g}-\mathrm{r}$ )-colours were used instead of ( $\mathrm{g}-\mathrm{i}$ ). The difference of the innermost ( $\mathrm{g}-\mathrm{r}$ )-colour to the colour of the LSB-component at $\mathrm{r}=2 \mathrm{R}_{\text {eff }}$ amounts for $\Delta(\mathrm{g}-\mathrm{i})_{\text {LSB }}=0.178$ mag. Assuming that the LSB-component has constant (g-r)-colours, the innermost colour has to be increased by 0.26 mag and $\mu_{0}$ has to be increased, too. With this information we were able to construct the new SBP and integrate it to obtain the new effective radius. In case of VCC-0274, the effective radius change from $\mathrm{R}_{\text {eff }}=0.599 \mathrm{kpc}$ to $\mathrm{R}_{\text {eff }}=0.602 \mathrm{kpc}$. Thus, the influence on the effective radius due to a change of the colour gradient is minor (see Tab. 2.7).
The (g-i)-colour profiles of the LSB-components are shown in Fig. 2.14 and the averages of the different morphological types are summarised in Tab. 2.7.


Figure 2.14: Radial colour profiles of the BCDs without decomposition (upper panel) and of the LSB-components (lower panel). The colour profiles are normalised to the effective radii of the LSB-components.

Figure 2.15 shows the colour gradients of the Virgo galaxies as a function of their magnitude. The scatter for the dEs became larger with increasing magnitude, but on average it is almost zero. This scatter is reasonable, since with increasing magnitude

Table 2.7: Average (g-i)-colour gradients of the different morphological types.

| Type | $\langle\Delta(\mathrm{g}-\mathrm{i})\rangle$ | Radius range |
| :--- | :---: | :--- |
| BCD | 0.26 | $\mathrm{r} / \mathrm{R}_{\text {eff }}=1 / 8$ to $\mathrm{r} / \mathrm{R}_{\text {eff }}=2$ |
| LSB | 0.14 | $\mathrm{r} / \mathrm{R}_{\text {eff }}=1$ to $\mathrm{r} / \mathrm{R}_{\text {eff }}=2$ |
| dE | 0.01 | $\mathrm{r} / \mathrm{R}_{\text {eff }}=1$ to $\mathrm{r} / \mathrm{R}_{\text {eff }}=2$ |
| $\mathrm{dE}(\mathrm{bc})$ | 0.03 | $\mathrm{r} / \mathrm{R}_{\text {eff }}=1$ to $\mathrm{r} / \mathrm{R}_{\text {eff }}=2$ |

also the $\mathrm{S} / \mathrm{N}$ decreases. The same behaviour like for the dEs can be observed for the LSB-components with much larger scatter. In case of the LSB-components, this scatter cannot be explained by the $\mathrm{S} / \mathrm{N}$, since the scatter is present at all magnitudes. Without accounting for the fainter $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$ with $\mathrm{M}_{\mathrm{r}}>-15 \mathrm{mag}$, the $\mathrm{dE}(\mathrm{bc})$ s show an increase of the colour gradient for fainter magnitudes.


Figure 2.15: (g-i)-colour gradient as a function of the absolute r-band magnitude for the LSB-components (red dots), dEs (black dots) and dE(bc)s (green asterisks).

### 2.4.8 Where are galaxies with uncertain morphological classification?

In the classification scheme of the VCC one also finds 19 galaxies with an uncertain morphological classification like "Im / BCD". In Tab. 2.9 these uncertain types are summarised. Using the above described parameter-space $\left(\mathrm{M}_{\mathrm{r}}-(\mathrm{g}-\mathrm{i})_{\mathrm{LSB}}-\mathrm{R}_{\text {eff }}-\langle\mu\rangle_{\text {eff }}\right)$ of the BCDs, we are able to conclude whether these galaxies truly share the same parameter-space like the BCDs or not. If a galaxy occupies the same region in the
parameter-space, then one can conclude that this galaxy is probably a BCD. In a first step, we compared the derived parameters of irregular (dI) galaxies with that of BCDs. By inspecting the properties of these types in the parameter-space, we found that VCC1374 clearly falls within the same parameter-space like the other BCDs.
For the galaxies VCC-0309 and VCC-2037, we found that most of the parameters are in the same region where irregular galaxies are located. For the rest of the galaxies a clear distinction between BCDs and dIs is not possible, because they are located in the overlap zone between these two dwarf galaxy classes.
There are three galaxies initially classified as "Sm / BCD" and "SBm / BCD", respectively. VCC-1356 is located in the same parameter-space as the BCDs. VCC1725 and VCC-1791 do not show a clear membership or separation from the BCDs. The galaxy VCC-0281 was taken from JL09 and was classified as "dS?/BCD?" with strong morphological uncertainties. It was also described by Lisker et al. (2006a) as a $\mathrm{dE}(\mathrm{bc})$. In the $(\mathrm{g}-\mathrm{i})-\mathrm{M}_{\mathrm{r}}$-plane, VCC-0281 occupies the same region like the BCDs, however, this is not the case with regard to the other two planes $\left(\mathrm{R}_{\text {eff }}-\mathrm{M}_{\mathrm{r}}-\langle\mu\rangle_{\text {eff }}\right)$. Regarding their classification, VCC-0213 ("dS?/BCD?") and VCC-1955 ("Spec/BCD") show features of spiral structure. Indeed, looking at the co-added SDSS gri-images one can find these spiral structures. The spiral structure is clearly visible when applying unsharp masks (see Lisker et al. (2006b) and Lisker et al. (2006a) for detailed description). In the classification scheme by Sung et al. (2002), they reported of "spiral structures with compact off-centred core" for the "Postmerger BCDS". But, to the knowledge of the authors, there is no detailed study about BCDs with spiral structure visible at optical wavelengths. Only HI observations of BCDs by van Zee et al. (2001) reveal rotation in BCDs, but the optical morphology of the used sample has smooth, symmetrical isophotes. In the CMD of Fig. 2.7, galaxy VCC-0213 and VCC-1955 are located at the brighter and redder end of the BCD regime. In Fig. 2.8, VCC-0213 has almost the same parameters like the BCDs. The effective radius of VCC-1955 is too large compared to the other BCDs. In Fig. 2.9, VCC-0213 is again in the regime of the BCDs, albeit it is at the boundary of the parameter-space populated by BCDs. VCC1955 shows too large values for both parameters. Using all these informations, it is possible that VCC-0213 is a BCD with obvious optical spiral structure. Nevertheless, the red core of VCC-0213 contradicts the classification of a BCD. The spiral structure, together with the red core (may be a bulge?), might hint to a spiral galaxy with a larger distance than the one of the GOLDMine data base of $\mathrm{D}=17.0 \mathrm{Mpc}$. To investigate this issue it is worth to analyse the spectrum of the galaxy. The spectrum
of VCC-0213 gives a redshift of $\mathrm{z}=-0.0006$ (blue shifted), with a corresponding heliocentric recessional velocity regarding GOLDMine of $\mathrm{V}_{\text {helio }}=-165 \mathrm{~km} / \mathrm{s}$, which indeed locates VCC-0213 at a distance of 17 Mpc (see Fig. 2.17). The upper left panel of Fig. 2.16 shows the co-added SDSS gri-image of galaxy VCC-0213 and the residual image in the upper right panel. The lower panel of Fig. 2.16 shows the residual fluxconserving unsharp mask image of VCC-0213 (cf. Papaderos and Östlin, 2012), which shows the colours of the spiral structure. Interestingly, the colour distribution of the spiral structures is not homogeneous, but shows a gradient from left to right.


Figure 2.16: Top: co-added SDSS gri-image (left) and zoomed residual image using a two pixel kernel (right) of the galaxy VCC-0213. Bottom: flux-conserved and combined unsharp mask image of VCC-0213.

At first glance the galaxy VCC-0655 (Spec/BCD) seems to be too bright compared to the BCDs. Additionally, its $\mathrm{R}_{\text {eff }}$ is too large. The gri-image displays several SF regions in the central part of the galaxy, which are surrounded by a redder stellar host. This is typical for BCDs and furthermore VCC-0655 has some similarities to the BCD Mrk-
86. The sample of BCDs in Papaderos et al. (2008) also shows a large spread in radius from very compact objects up to object with $\mathrm{R}_{\text {eff }} \gtrsim 1 \mathrm{kpc}$ (see also Section 2.5.4). Thus, it cannot be ruled out that VCC-0655 is a BCD.
The galaxy VCC-0135 (Spec/BCD) covers almost the same region in the $\mathrm{R}_{\mathrm{eff}}-\mathrm{M}_{\mathrm{tot}^{-}}$ $\langle\mu\rangle_{\text {eff }}$ as the other BCDs. However, its ( $\mathrm{g}-\mathrm{i}$ )-colours are too red. The gri-image shows a galaxy with red core, comparable to VCC-0123, but there is no obvious spiral structure visible. The spectrum shows strong emission lines and indeed a detailed inspection of the "core-region" reveals a blue structure, which is located in the region of the SDSS spectroscopic fibre. With this information about VCC-0123, we go along with the former classification, but indicate the new classification with a strong uncertainty regarding the BCD class ("Spec/BCD?").
The analysis of the two unknown galaxies ("?" in Tab. 2.9) shows that the properties of VCC-0429 and VCC-1713 have more similarities to the dIs. This is also supported by the optical images of VCC-0429 and VCC-1713. Therefore, these two galaxies should be classified as a transition type ( $\mathrm{BCD} / \mathrm{dI}$ ).

Table 2.8: Colour gradients of the BCDs.

| VCC | $\Delta(\mathrm{g}-\mathrm{i})$ | $\Delta(\mathrm{g}-\mathrm{r})$ | $\Delta(\mathrm{i}-\mathrm{z})$ | $\Delta(\mathrm{u}-\mathrm{g})$ | $\mathrm{R}_{\text {eff,r }}[\mathrm{kpc}]$ | $\mathrm{R}_{\text {eff,r,nograd }}[\mathrm{kpc}]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0001 | 0.052 | -0.018 | -0.240 | -0.172 | 0.395 | 0.401 |
| 0010 | -0.100 | -0.117 | -0.145 | -0.338 | 0.422 | 0.451 |
| 0022 | 0.300 | 0.132 | -0.103 | -0.584 | 0.423 | 0.464 |
| 0024 | 0.147 | 0.025 | 0.035 | -0.589 | 0.419 | 0.452 |
| 0074 | 0.101 | 0.141 | 0.499 | -0.304 | 0.566 | 0.597 |
| 0130 | 0.093 | 0.042 | -0.133 | 0.309 | 0.355 | 0.351 |
| 0144 | 0.564 | 0.322 | 0.106 | 0.466 | 0.339 | 0.396 |
| 0172 | 0.030 | -0.003 | -0.293 | -0.096 | 1.068 | 1.146 |
| 0207 | 0.903 | 0.305 | 0.658 | -1.177 | 0.263 | 0.293 |
| 0223 | 0.178 | 0.162 | 0.184 | -0.610 | 0.408 | 0.472 |
| 0274 | 0.249 | 0.178 | 0.546 | -1.261 | 0.599 | 0.602 |
| 0324 | 0.112 | 0.036 | 0.271 | -0.006 | 1.171 | 1.227 |
| 0334 | 0.233 | 0.009 | 0.247 | -0.537 | 0.542 | 0.566 |
| 0340 | 0.331 | 0.061 | 0.271 | -0.584 | 0.675 | 0.760 |
| 0410 | 0.306 | 0.292 | 1.424 | -1.195 | 0.306 | 0.344 |
| 0428 | 0.609 | -0.517 | -0.766 | 0.956 | 0.438 | 0.418 |
| 0459 | 0.439 | 0.215 | -0.298 | -0.779 | 0.669 | 0.717 |
| 0513 | 0.012 | -0.084 | -0.119 | 0.463 | 0.563 | 0.553 |
| 0562 | 0.749 | 0.371 | 1.224 | -1.081 | 0.532 | 0.565 |
| 0641 | 0.005 | 0.151 | 0.162 | -0.489 | 0.462 | 0.474 |
| 0802 | 0.364 | 0.188 | -0.369 | -0.467 | 0.683 | 0.762 |
| 0841 | 0.096 | -0.031 | 0.043 | -0.021 | 1.135 | 1.235 |
| 0890 | 0.067 | -0.019 | -0.167 | 0.779 | 0.388 | 0.381 |
| 1141 | 0.246 | 0.054 | 0.143 | -0.275 | 0.462 | 0.463 |
| 1313 | 0.384 | 1.164 | 1.842 | -2.629 | 0.264 | 0.308 |
| 1411 | 0.484 | 0.135 | 0.368 | -1.530 | 1.092 | 1.053 |
| 1437 | 0.170 | 0.051 | -0.185 | -0.103 | 0.504 | 0.560 |
| 1459 | 0.169 | 0.054 | -0.533 | 0.301 | 0.723 | 0.628 |
| 1572 | -0.284 | -0.252 | 1.056 | -1.482 | 1.061 | 1.008 |
| 1744 | 0.496 | 0.231 | -0.199 | -0.056 | 0.597 | 0.579 |
| 2015 | 0.317 | 0.243 | 0.133 | 0.263 | 0.660 | 0.706 |
| 2033 | 0.453 | 0.077 | -0.401 | -0.243 | 0.842 | 0.856 |
|  |  |  |  |  |  |  |

Table 2.9: Uncertain morphological classifications.

| VCC | Type | $\mathrm{M}_{\text {tot, }}$ | $\mathrm{R}_{\text {eff,r }}[\mathrm{kpc}]$ | $\langle\mu\rangle_{\text {eff,r }}$ | $(\mathrm{g}-\mathrm{i})$ |
| :--- | :---: | :---: | :---: | :---: | :---: |
| 0135 | Spec / BCD | -17.34 | 0.67 | 20.36 | 0.98 |
| 0213 | $\mathrm{dS} ? ~ / ~ B C D ? ~$ | -17.50 | 0.61 | 20.01 | 0.80 |
| $0281^{*}$ | $\mathrm{dS0}$ or BCD | -16.37 | 1.13 | 22.50 | 0.77 |
| 0309 | $\mathrm{Im} / \mathrm{BCD}$ | -15.33 | 1.07 | 23.41 | 0.50 |
| 0429 | $?$ | -15.10 | 0.73 | 22.81 | 0.65 |
| 0446 | Im / BCD: | -15.94 | 0.87 | 22.34 | 0.75 |
| 0655 | Spec,N: / BCD | -18.73 | 1.49 | 20.73 | 0.86 |
| $0737^{* *}$ | Sd / BCD? | - | - | - | - |
| $0848^{* *}$ | ImIIIpec / BCD | - | - | - | - |
| 1179 | ImIII / BCD | -16.34 | 0.97 | 22.18 | 0.64 |
| 1356 | SmIII / BCD | -16.34 | 0.54 | 20.89 | 0.48 |
| 1374 | ImIII / BCD | -16.99 | 0.79 | 21.09 | 0.40 |
| 1427 | Im / BCD: | -16.45 | 0.93 | 21.98 | 0.68 |
| 1713 | ? | -15.66 | 0.75 | 22.30 | 0.56 |
| 1725 | SmIII / BCD | -17.33 | 1.56 | 22.22 | 0.54 |
| 1791 | SBmIII / BCD | -17.20 | 1.43 | 22.15 | 0.34 |
| 1804 | ImIII / BCD | -16.04 | 0.92 | 22.37 | 0.82 |
| 1955 | Spec / BCD | -18.35 | 2.33 | 22.07 | 0.85 |
| 1960 | ImIII / BCD? | -13.70 | 0.37 | 22.73 | 0.74 |
| 2007 | ImIII / BCD: | -15.89 | 0.68 | 21.85 | 0.69 |
| 2037 | ImIII / BCD | -16.26 | 1.66 | 23.42 | 0.79 |

Notes. * from JL09; ** not covered by SDSS

### 2.5 Discussion

### 2.5.1 Difference between certain and possible cluster member BCDs

From Tab. 2.10 it is apparent that various photometric parameters of certain Virgo BCDs and cluster-member candidates are identical within $\sim 1 \sigma$. Therefore, in the following discussion no differentiation based on Virgo membership was applied, even though photometric quantities in Tab. 2.10 are listed separately.
To avoid contamination due to extreme outliers, the average values were determined within the colour interval $(\mathrm{g}-\mathrm{i})>0 \mathrm{mag}$ (see Section 2.7 for explanation).

Table 2.10: Comparison of the BCDs regarding their membership to Virgo.

| Parameter | Certain | $\sigma$ | Possible | $\sigma$ |
| :--- | :---: | :---: | :---: | :---: |
| $\left\langle M_{r}\right\rangle[\mathrm{mag}]$ | -15.488 | 0.835 | -15.489 | 0.783 |
| $\langle g-i\rangle[\mathrm{mag}]$ | 0.589 | 0.254 | 0.696 | 0.172 |
| $\left\langle R_{\text {eff, },}\right\rangle[\mathrm{kpc}]$ | 0.569 | 0.251 | 0.650 | 0.275 |
| $\langle\mu\rangle_{\text {eff,r }}\left[\mathrm{mag} / \mathrm{arcsec}^{2}\right]$ | 21.697 | 0.843 | 21.963 | 0.950 |

### 2.5.2 Distance dependence of the results

The GOLDMine data base provides the distances of the BCDs, which can be found in column [13] in Tab. 6.2. Table 2.5 .2 shows the different distances and the resulting errors if one assumes a distance of $\mathrm{d}=16.52 \mathrm{Mpc}(\mathrm{m}-\mathrm{M}=31.09)$ instead of the GOLDMine value. $\alpha\left(\mathrm{R}_{\text {eff }}\right)$ corresponds to the factor, about which $\mathrm{R}_{\text {eff }}[\mathrm{kpc}]$ changes due to different distances. Additionally, the spacial location within the Virgo cluster of the corresponding distances are given (Gavazzi et al., 1999, hereafter G99). The distances of different substructures of Virgo in Tab. 2.5.2 were derived by using the Tully-Fisher relation (Tully and Fisher, 1977) for a sample of spiral galaxies and the Fundamental Plane (Djorgovski and Davis, 1987; Dressler et al., 1987) for the E/S0 in the study of G99. G99 found three separate regions in their Figs. 4 and 5 (velocity $\left(v_{\mathrm{LG}}\right)^{7}$ versus the distance of the galaxies), which in turn define the distances of the substructures. However, since the allocation of the BCDs to the different substructures of the Virgo cluster only depends on the projected location of the galaxies, the GOLDMine distances can not be taken for granted. For instance, the column depth

[^7]of the different substructures in Virgo is not taken into account.
Additionally, it needs to be pointed out that the finite extent of the cluster causes a natural scatter in the distance. As an example, galaxy VCC-1313 has the smallest cluster centric distance to M87, but it also has the bluest (g-i)-colours within the sample. Therefore, VCC-1313 could fall into the Virgo cluster from a larger deprojected cluster centric radius in the fore- or background of the cluster. Such affects are not accounted for by the GOLDMine data base.

Table 2.11: GOLDMine distances of the different Virgo clouds. $\Delta \mathrm{m}$ and $\alpha\left(\mathrm{R}_{\mathrm{eff}}\right)$ indicate the variation of the corresponding parameters if the GOLDMine distance is used instead of a constant distance of 16.52 Mpc .

| Distance <br> $[\mathrm{Mpc}]$ |  | Location in Virgo cluster | $\Delta \mathrm{m}[\mathrm{mag}]$ <br> to d$=16.52 \mathrm{Mpc}$ |
| :---: | :--- | :---: | :---: | | $\alpha\left(\mathrm{R}_{\text {eff }}\right)$ |
| :---: |
| to d $=16.52 \mathrm{Mpc}$ |

Figure 2.17 shows the distance in Mpc of the Virgo BCDs against the velocity given by GOLDMine (see also Fig. 5 in G99). Regarding G99 galaxies with $v_{\text {LG }}>1900 \mathrm{~km} / \mathrm{s}$ belong to the M or W cloud with a distance of 32 Mpc . As one can see, four BCDs (VCC-0022, VCC-0024, VCC-0274 and VCC-0340) have a GOLDMine distance of 32 Mpc , but do not have velocities of $\mathrm{v}_{\mathrm{LG}}>1900 \mathrm{~km} / \mathrm{s}$. Therefore, we do not assume a distance of 32 Mpc , but instead a distance of 16.52 Mpc . These galaxies are displayed by green squares in Fig. 2.18. For all other BCDs the GOLDMine distances are adopted.
Figure 2.18 shows the $\left(\mathrm{M}_{\mathrm{r}}-\mathrm{R}_{\text {eff }}-\left\langle\mu_{\text {eff }}\right\rangle\right)$-plots of the analysed galaxies with the GOLDMine distances. Interestingly, there are only three LSB-components of the BCDs that do not fall into the ETG's $2 \sigma$-plane and these BCDs are classified as "il" and "nI". One of these three BCDs also needs a profile flattening of the LSB-component, bringing it even more into the $2 \sigma$-plane.
In the $\mathrm{M}_{\mathrm{r}}-\mathrm{R}_{\text {eff }}$-diagram, there are two other BCDs outside the $2 \sigma$-deviation of the dEs , where the velocity criterion is not fulfilled. The comparison of the late-type dwarfs ( $\mathrm{BCD}(-\mathrm{LSB}$ ) and dI) and the early-type dwarfs reveals another interesting result. The LSB-components of the BCDs and the dI galaxies form a continuous sequence, which has the same location as the early-type dwarf galaxies. This is consistent with the hypothesis that, once star formation in Virgo BCDs has ceased, their


Figure 2.17: GOLDMine distance vs. the velocity of the Virgo BCDs. The vertical solid line at $\mathrm{v}_{\mathrm{LG}}=1900 \mathrm{~km} / \mathrm{s}$ corresponds to the limit from which galaxies belong to the M and W cloud regarding Gavazzi et al. (1999). The four BCDs in the box are below or slightly above this velocity limit, but have a GOLDMine distance of 32 Mpc .
descendants could evolve into dEs. Additionally, Virgo dIs could form the progenitor population of dEs with the largest effective radii. Clearly, one has to keep in mind that $R_{\text {eff }}$ is not a particularly sensitive discriminator between SBPs of substantially different shapes, consequently the alignment of all dwarf galaxy types along a common sequence on the $\mathrm{M}_{\mathrm{r}}-\mathrm{R}_{\text {eff }}$ diagram is unsurprising. Moreover, it should be pointed out that these findings are critically dependent on the distance of the galaxies. As seen in Figs. 2.8 and 2.9, there are more BCDs outside the $2 \sigma$ deviation when a constant distance of 16.52 Mpc is adopted.

### 2.5.3 Comparison between BCDs and the inner components of dEs

A recent study by Janz et al. (2012) investigated the different components of $\sim 70$ dEs in the Virgo cluster in H -band in the range of $-19 \leq \mathrm{M}_{\mathrm{r}} \leq-16$ mag. As mentioned before, dEs do not form a homogeneous class of objects, but show disc features and inner components, which were modelled with GALFIT 3.0 (Peng et al., 2010) in the study of Janz et al. (2012). Apart from the H-band, also r-band profiles of the dEs were analysed, thus, we are able to compare the inner components of the dEs with our sample. Figure 2.19 shows the results of our study of the LSB-components and the results for the different components of the dEs.
Interestingly, the inner components of dEs and the LSB-components cover a very sim-


Figure 2.18: Same as Figs. 2.8, 2.9 and 3.19, but instead of a constant distance of 16.52 Mpc , the distances given by the GOLDMine data base are used for the LSBcomponents. Green squares correspond to BCDs with a GOLDMine distance of 32 Mpc, but with velocity $\mathrm{v}_{\mathrm{LG}} \lesssim 1900 \mathrm{~km} / \mathrm{s}$. Therefore, a distance of 16.52 Mpc was applied.


Figure 2.19: Same as Fig. 2.9, but additionally the inner components of dEs are shown. Lower panels show the inverted SDSS images of three example dEs, which inner components cover the region of $0.3<\mathrm{R}_{\text {eff }}<0.7 \mathrm{kpc}$ and $-14.4>\mathrm{M}_{\mathrm{r}}>-16.4$ mag. This parameter region corresponds to the bulk of the LSB-components.
ilar region in the $\mathrm{R}_{\text {eff }}-\mathrm{M}_{\mathrm{r}}$-parameter-space. The physical reason for this coincidence is not clear at the moment and it is also not clear if there are any connections between the inner components of dEs and LSB-components of BCDs. One should also keep in mind that the inner components of dEs are accompanied by outer (disc) components. While the outer tails of some BCDs qualify as an equivalent outer component, for most of them these are not as prominent (or even not detected) as in the multicomponent dEs. This reflects again the fact that the LSB-components, even with the outer tails,
only overlap with the more compact $\sim$ half of the dEs. Therefore, if the similarity to the inner components of dEs had an evolutionary meaning, LSB-components would have to form a surrounding disc to look like present day dEs.

### 2.5.4 Comparison with other BCDs

Table 2.12 shows the results of a sample of BCDs from Papaderos et al. (2008) (hereafter P08) from the SDSS. Values of the LSB-component of this study were transformed into an effective radius using the formula of Graham and Driver (2005). All given values of P08 were measured in the SDSS g-band, therefore, $\mathrm{R}_{\text {eff }}$ for the r-filter was calculated with the following approximation:

$$
\mathrm{R}_{\mathrm{eff}, \mathrm{r}}=\frac{1}{0.95} \cdot \mathrm{R}_{\mathrm{eff}, \mathrm{~g}} .
$$

As one can see, the effective radii of the P08 sample are also very compact with an average of $\left\langle\mathrm{R}_{\text {eff,g }}\right\rangle=0.62 \mathrm{kpc}$, but with a large scatter.

Table 2.12: Galaxies from the study of Papaderos et al. (2008).

| Galaxy | $\mathrm{m}_{\text {tot,g }}$ | $\mathrm{M}_{\text {tot,g }}$ | $\mu_{0, \mathrm{~g}}$ | $\mu_{\mathrm{e}, \mathrm{g}}$ | $\langle\mu\rangle_{\mathrm{e}, \mathrm{g}}$ | $\mathrm{R}_{\mathrm{eff}, \mathrm{g}}$ <br> $[\operatorname{arcsec}]$ | D <br> $[\mathrm{Mpc}]$ | $\mathrm{R}_{\mathrm{eff}, \mathrm{g}}$ <br> $[\mathrm{kpc}]$ | $\mathrm{R}_{\mathrm{efff}, \mathrm{r}}$ <br> $[\mathrm{kpc}]$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{J} 0133+1342$ | 17.9 | -15.0 | 22.1 | 34.7 | 23.22 | 4.63 | 37.8 | 0.85 | 0.89 |
| $\mathrm{~J} 1044+0353$ | 17.2 | -16.3 | 19.9 | 35.3 | 21.02 | 2.32 | 51.2 | 0.58 | 0.61 |
| $\mathrm{~J} 1201+0211$ | 17.4 | -13.3 | 22.2 | 32.5 | 23.32 | 6.10 | 14.0 | 0.41 | 0.44 |
| $\mathrm{~J} 1414-0208$ | 18.1 | -13.7 | 21.6 | 33.6 | 22.72 | 3.35 | 23.0 | 0.37 | 0.39 |
| J2230-0006 | 17.0 | -15.0 | 19.8 | 33.8 | 20.92 | 2.43 | 24.9 | 0.29 | 0.31 |
| J2302+0049 | 18.6 | -17.0 | 20.8 | 37.4 | 21.92 | 1.84 | 134.6 | 1.20 | 1.26 |

However, there are also other examples of BCDs, having a very extended LSBcomponent. One example of these BCDs is Markarian 86 (Mrk-86). The SBP of Mrk-86 is shown in Fig. 2.20. The derived effective radius at an assumed distance of 7.0 Mpc amounts for $\mathrm{R}_{\mathrm{eff}, \mathrm{R}}=0.94 \mathrm{kpc}$ (regarding Graham's formula and with values from Papaderos et al., 1996b) in the Johnson-filter system or $\mathrm{R}_{\text {eff,r }}=1.2 \mathrm{kpc}$ (from profile integration) in $r$ of SDSS. Compared to the Virgo BCDs, even these extended BCDs are in the same radius range. Interestingly, Mrk-86 has a very red LSB-colour of ( $\mathrm{g}-\mathrm{i}$ ) $=0.92$ mag. Together with $\mathrm{R}_{\text {eff }, \mathrm{R}}, \mathrm{M}_{\mathrm{tot}, \mathrm{g}}=-17.06 \mathrm{mag}$ and $\langle\mu\rangle_{\text {eff }, \text { LSB }}=20.98 \mathrm{mag} / \mathrm{arcsec}^{2}$, it occupies the same region as the $\mathrm{dE}(\mathrm{bc})$ population of Virgo.


Figure 2.20: The SBP of the BCD Mrk-86, which is not located in the Virgo cluster.

Since the BCDs of our study are located in the region of the Virgo cluster, one may ask whether they are some special kind of BCD and deviate from other (field) BCDs. To investigate this question a large sample of emission-line galaxies in arbitrary environments (unpublished study by H. T. Meyer, R. Kotulla, P. Papaderos, Y. Izotov, N.G. Guseva and K.J. Fricke) was used. The galaxies were spectroscopically selected, therefore, also non-dwarfish galaxies are included in the initial sample. To include only star forming dwarf galaxies, a g-band magnitude cutoff of $\mathrm{M}_{\text {cutoff }}=-18.0 \mathrm{mag}$ was applied and the resulting sample is called the second sample. To all star forming galaxies of this sample the same photometric analysis (derivation of SBPs and decomposition) was applied, thus the derived parameters of both samples can be compared. The CMD of the LSB-components of the Virgo BCDs and the SFDGs of the second sample is shown in Fig. 2.21. The LSB-components of the Virgo BCDs tend to be more reddish compared the other LSB-components of the second sample. This could be a hint to the influence of the denser cluster environment on the Virgo BCDs, making them redder due to e.g. the removal of gas. In the other two diagrams of Fig. 2.21, no obvious differences between these two samples can be found. Nevertheless, the SFDGs of the second sample tend to have larger effective radii than the BCDs of Virgo, which is may caused by selection effects. It also has to be pointed out that the sample of Virgo BCDs is small, making a comparison quite difficult. Therefore, an increase of the sample size of cluster BCDs (e.g. Perseus-, Coma-, Fornax cluster) would be desirable for future investigations.


Figure 2.21: Parameter-space of the LSB-components of our Virgo BCDs and a comparison sample of SFDGs in arbitrary environments. The black squares corresponds to the second sample of SFDGs, while the other galaxy types are colour code as in the previous figures of this chapter.

### 2.5.5 Evolutionary connection between different dwarf galaxies in galaxy clusters

Various observational and theoretical lines of evidence suggest that the evolutionary pathways of late-type dwarfs in a cluster environment have to significantly differ from those in the field. Even in the cluster periphery, late-type dwarfs have undergone interaction with the hot ICM and other nearby galaxies since a few hundreds of Myr, a time span on the order of their dynamical timescale. Starbursts transforming latetype dwarfs plunging into the ICM into BCDs, may have a different origin than starbursts in the general population of field BCDs. Consequently, one should be cautious when comparing structural and integral photometric properties (e.g. colour, burst parameter, effective radius, concentration index) of cluster-BCDs with literature data for normal BCDs. Various mechanisms for starburst ignition shortly in initial stages of the interaction of a dwarf with the ICM can be envisaged, for example triggered gas collapse by the external ICM pressure (as opposed to gas collapse driven by self-gravity), or strong dissipation and gas cooling in large-scale shocks, followed by collective star formation. For such reasons, it is not even for sure that the compact structure of the LSB-component - a typical, if not the most distinctive characteristic of field BCDs - has to be common in systems classified as BCDs in galaxy clusters. For example, external agents (see above) could ignite a BCD-typical starburst in a genuine (that is, relatively diffuse) dI. Alternatively, an externally triggered destabilisation and inflow of gas could eventually lead to an adiabatic contraction of the LSB-component (Papaderos et al., 1996a), and the opposite may happens after complete removal of the dwarf gas halo, as it sinks deeper within the hostile ICM environment. As our understanding of the early interaction of late-type dwarfs in the cluster environment is still poor, it would be speculative to generalise conclusions drawn from our Virgo BCD sample to the BCD population as a whole. Still, the discussion below is instructive in its own right, because it concentrates on a complete and well-selected sample of cluster-BCDs and probable candidates, thereby minimising contamination of the trends found for field BCDs.

The evolutionary scenario of Davies and Phillipps (1988) describes how a dwarf irregular galaxy can ignite a starburst and become a BCD due to in-falling gas from a surrounding reservoir.
After several starburst (BCD) phases and non-starbursting dI-phases, the galaxy evolves to a passive dE galaxy. If the latter evolutionary sequence was possible at present epochs, and specifically in the outskirts of the Virgo cluster, BCDs and dEs
should have a similar structure.
Observational studies of BCDs by Drinkwater and Hardy (1991) ${ }^{8}$ and Papaderos et al. (1996a) showed that BCDs are very compact objects compared to the dEs. They concluded that only the most compacted dEs may be related to the BCDs. An even more extreme conclusion is reached by Drinkwater et al. (1996), who state that "no blue star-forming progenitors of dE galaxies" exist in the Virgo cluster down to an absolute $B$-magnitude of $\mathrm{M}_{B} \approx-14$ mag. At first view, these findings seem to be supported by our photometric measurements for the galaxies' entire light: Fig. 2.6 (top left panel) shows that the majority of BCDs are more compact than even the $2 \sigma$-range of dEs at a given magnitude.

However, if the gas reservoir that is immediately available to a BCD's starburst region is limited, thus restricting the starburst phase to durations of some $10^{8} \mathrm{yrs}$ or less (Thuan, 1991; Thornley et al., 2000), a population of non-starbursting counterparts must exist within the same spatial volume. Assuming a velocity of $1000 \mathrm{~km} / \mathrm{s}$, a BCD could only move some 100 kpc during the starburst phase, i.e. some tenths of the Virgo cluster's virial radius - which means that the outskirts of the Virgo cluster should hold many non-starbursting counterparts to BCDs. How do these look like?
To investigate this, we determined the structural properties of the underlying LSBcomponent of the Virgo BCDs, i.e. without the contribution of the starburst region. The results already show much more overlap with the dEs in the parameter-space of structure and also colour (see Figs. 2.7, 2.8 and 2.9). When we also take into account the different distances of the various parts of the Virgo region, the LSBcomponents turn out to fit remarkably well to the dE population with small effective radii (Fig. 2.18). The left side of Fig. 2.22 shows a collection of such dEs that fall within the locus of the LSB-components: these may look like the future "red and dead" descendants of BCDs. Average dEs are shown on the right side for comparison. As for the few objects that remain more compact than the dEs, none of them actually has a clearly red LSB-component (Fig. 2.23). Judging from their measured values and their images, only VCC-1313 (probably also VCC-0410) seems to follow the literal meaning of the term "blue compact dwarf", while most of the other BCDs rather have the appearance of a dE or dI with additional starburst. The non-detection of a red LSB-component down to low surface brightness levels for some "extreme" BCDs was already reported by Drinkwater and Hardy (1991): two BCDs in their

[^8]

Figure 2.22: Possible descendants of BCDs in the Virgo cluster. The figure in the middle shows the $\mathrm{R}_{\text {eff }}-\mathrm{M}_{\mathrm{r}}$ diagram of the analysed galaxies with the GOLDMine distances, and the blueish box displays the locus of the LSB-components. The three example BCDs on the left (from top to bottom: VCC-0022, VCC1141 and VCC-1623) are selected in the parameter-interval of $0.5<\mathrm{R}_{\mathrm{eff}}<0.9 \mathrm{kpc}$, $-17.0<\mathrm{M}_{\mathrm{r}}<-14.5 \mathrm{mag}$ and $21.0<\langle\mu\rangle_{\mathrm{eff}, \mathrm{r}}<23.0 \mathrm{mag} / \operatorname{arcsec}^{2}$. The three dEs on the right hand side are randomly selected in the same parameter-interval (from top to bottom: VCC-0178, VCC-1308 and VCC-1980). All images have the same scale (x-edge length is 100 arcsec and 7.95 kpc for a distance of 16.52 Mpc ).
study had no surrounding low-surface-brightness envelopes detected to a limit of $\mu(\mathrm{V})=27 \mathrm{mag} / \operatorname{arcsec}^{2}$, which therefore would be classified as i0-BCDs. For all of these, the Next Generation Virgo Cluster Survey (NGVS, Ferrarese et al., 2012) will provide unprecedented insight into their outer structure and older stellar population, if existent.

Are dEs the only morphological type that qualifies as non-starbursting counterparts of BCDs? Figure 2.24 shows the number distribution of galaxies in the outer region of the Virgo cluster that fall in the same parameter-space of magnitude and mean surface brightness as the LSB-components of the BCDs ${ }^{9}$, displayed separately for blue and

[^9]

Figure 2.23: BCDs that are more compact than the dEs and are outside the $2 \sigma$ deviation in Fig. 2.18. The upper panels show the co-added gri-SDSS images and the lower panels show the contour maps of the BCDs. From left to right: VCC-0130, VCC-0410 and VCC-1313.
red ( $g-i>0.7$ ) galaxies. Only for blue galaxies (left panel), late-type spirals and especially irregular galaxies form a population of significant size compared to the dEs. However, the majority of the LSB-components of BCDs is already in the red colour regime (right panel), where the dEs clearly outnumber all other galaxy types.

Vaduvescu et al. (2006) concluded that "BCDs and dIs are similar structurally", which is in conflict with our results (see Figs. 2.8 and 2.9). We measure dIs to be systematically more extended than the LSB-components of BCDs at a given magnitude, hence only a small fraction of the dIs would qualify as non-starbursting BCDs. Vaduvescu et al. fitted structural parameter relations to dIs and claimed consistency of the BCDs with these relations. Their Figs. 4 and 5 show large scatter among their BCD sample, with at least half of the BCDs lying clearly at smaller radii than the dI relations. Among the five BCDs that lie significantly above the dI relation, the three largest ones have been classified as ambiguous types between $\mathrm{Sm}, \mathrm{dI}$ and BCD in the VCC, and appear closer to Sm/dI by visual inspection of SDSS colour


Figure 2.24: Late-type and early-type galaxies in the outer region of Virgo ( $\mathrm{D}_{\mathrm{M} 87}>0.6 \mathrm{Mpc}$ ). In the left panel a cutoff of $\mathrm{M}_{\mathrm{r}}<-18 \mathrm{mag}$ and $\left.\left\langle\mu_{\mathrm{r}}\right\rangle_{\mathrm{eff}, \mathrm{r}}\right\rangle 23 \mathrm{mag} / \mathrm{arsec}^{2}$ was applied to restrict the parameter region to the locus of most of the LSB-components. The same cutoffs were used in the right panel but with an additionally colour cutoff of $(\mathrm{g}-\mathrm{i})>0.7$ mag. The histogram of the dEs also includes the $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$. The morphological types are taken from the VCC.
images. This may provide an explanation for the discrepant findings.
In summary, if the starbursts of the BCDs are relatively short, then many of the dEs in the Virgo outskirts with similar magnitudes must be the non-starbursting counterparts to the BCDs. Unless we are witnessing all BCDs in their very last starburst phase, which seems unlikely, this unavoidably means that a number of dEs must be able to re-ignite a starburst. Interestingly, a recent paper by Hallenbeck et al. (2012) finds gas in Virgo dEs and furthermore studies by Gu et al. (2006) and Lisker et al. (2006a) find dEs with blue cores, indicating recent or ongoing star formation. At first glance, one may speculate that these type are related with each other. However, we find that the $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$ are on average less compact and brighter than the LSB-components of the BCDs (see Figs. 2.8 and 2.9).
To gain a deeper insight into the connections of BCDs and dEs, a large HI survey of both populations is worthwhile and may also change our view from "red and dead" dEs to more complex systems. This "increase" of complexity already started with several studies of dEs with spiral feature, discs or blue cores (e.g. Jerjen et al., 2000; Lisker et al., 2006b; Hallenbeck et al., 2012), and the improvement of observational
techniques (e.g. E-ELT ${ }^{10}$, JWST ${ }^{11}$ ) and new surveys (e.g. NGVS ${ }^{12}$ ) will push this complexity even further. Of course, this is also true for the BCDs, where new observations will probably reveal new substructure in the SB-regions as well as in the underlying LSB-component - e.g. i0-BCDs, where no LSB-component is detected until now.

[^10]
# Chapter 3 The structural properties of latetype galaxies in the Virgo cluster 


#### Abstract

In this chapter, the structural parameters of the late-type galaxy population of the Virgo cluster are analysed. Furthermore, the distribution of the galaxies within the Virgo cluster and the dependence of the morphology according to the distance to M87 is investigated. Finally, the entire photometric and structural informations of the galaxies are used to review and update the initial classifications. To study the possible evolutionary connections between late- and early-type dwarf galaxies, the properties of these two groups are compared. Together with the results of the BCDs and dIs from the previous chapter, the results of this chapter form the foundation for the investigation of the evolution with time of the late-type galaxies.


The results of this study were partly presented in Meyer and Lisker (2011) and throughout several conference contributions.

### 3.1 Introduction

Galaxies in the Virgo cluster, but also in other regions in the Universe, show a large variety in morphology, ranging from almost featureless galaxies to galaxies with impressing spiral structures. Moreover, galaxies show a diversity in colours, from relatively red to blue colours. The former class of galaxies with their smooth shape and red colours are called early-type galaxies (ETGs), while the latter are called latetype galaxies. The terms "early" and "late" do not give any hint to the evolutionary state of the galaxies and are mostly used for historical reasons. For instance, earlytype galaxies are dominated by an old stellar population, while the luminosity of the
late-type galaxies is dominated by young and massive stars (e.g. Sparke and Gallagher, 2000).

ETGs show a large spread in magnitude, ranging from a r-band magnitude of $\mathrm{M}_{\mathrm{r}}=-13$ mag for the most faintest ones of up to $\mathrm{M}_{\mathrm{r}}=-23 \mathrm{mag}$ for the brightest ETGs in the sample of this thesis. This means that the faintest ETG differs from the most luminous one by a factor of $10^{4}$ ! The gigantic ETGs are also the largest galaxies in diameter of the entire morphological sequence, while the dwarfish ellipticals are the most frequently observed galaxy type.
As mentioned above, a first inspection of the optical images of ETGs show galaxies that are characterised by a smooth appearance and relatively red colours. However, studies of ETGs show that there are several subclasses with centrally concentrated star formation (Lisker et al., 2006a), disc features or spiral arms (Jerjen et al., 2000; Lisker et al., 2006b, 2007). These initially believed to be simple systems, are thus much more complicated.
Apart from the ETGs, there is another population of galaxies with bluer colours at a given magnitude. One of the representative of this morphological class shows prominent and striking spiral arms. These galaxies also show a large range in morphology ranging from Sa to Sm galaxies and are called spiral galaxies.
Looking at optical images of spiral galaxies, one realises that Sa galaxies have an inner dense region, which is most prominent in the class of Sa galaxies, and is called the bulge ${ }^{1}$. The stars in the bulge have random motions and the colours are generally red. From Sa to Sc galaxies the bulge became less prominent and vanished for the late-type spirals of type Sd and Sm . To characterise the importance of the bulge, the bulge-todisc (B/D) ratio was introduced, which describes the ratio of the light contribution of the bulge to the one of the disc (Sandage, 1961; Graham and Worley, 2008). Bulge dominated systems like Sa galaxies have B/D ~ 0.3, while later types like Sc have B/D ~ 0.05 (Mo et al., 2010).
The dust in spiral galaxies is found to be located in a layer, which is very thin and has a spatial extend comparable to the stellar population (Xilouris et al., 1999). The extinction due to dust is also responsible to influence the colours of galaxies. Since dust extinction depends on the observed wavelength in way that the influence in the infrared is less than in UV, galaxies appear to be redder in the presence of dust. Therefore, the colours of a dust dominated galaxy do not represent the light distribution of the current

[^11]stellar population. Due to the distribution of the dust in a layer, the inclination angle also plays an important role. Galaxies that are observed face-on are less influenced by dust extinction than edge-on galaxies (Cunow, 1999).
Along the sequence of late-type galaxies the gas mass fraction varies strongly. McGaugh and de Blok (1997) showed that the gas mass fraction is a function of the central surface brightness $\mu_{0}$ and the colour of the disc in a way that galaxies with a high surface brightness and red colours have a lower gas mass fraction. That means, the gas mass fraction of Sa galaxies is lower compared to late spirals. The different phases of the gas are also used as a tracer e.g. for the star formation rate (SFR) and metallicity $(\mathrm{Z})$ and it was found that the SFR and Z are not homogeneous over the entire galaxy. For instance, observations and simulations of Virgo spirals reveal negative metallicity gradients (Magrini et al., 2011; Roediger et al., 2011).
The informations about the SFR of spiral galaxies are obtained from observation in $\mathrm{H} \alpha$, which are available for many spiral galaxies in the Virgo cluster. Koopmann and Kenney (2004) found a reduction of the SFR in the outer part of the discs of Virgo spirals compared to field spirals of the same morphological type. Since the SFRs in the inner parts of the cluster spirals and field spirals are similar, they concluded that the reduction of the SFR is due the truncation of the disc in the dense cluster environment when the galaxies fall into the cluster. Comparing the late-type population of a cluster to the one in the field, it was shown that both populations differ from each other. For instance, the global SFR differs between galaxies in a cluster and galaxies in the field, what points to a transformation of galaxies, which enters the cluster. Indeed, Tully and Shaya (1984) showed by models that spiral galaxies just entered the Virgo cluster within several Gyr. A more detailed review of the differences of field and cluster galaxies is presented by Boselli and Gavazzi (2006).
Looking at the number distribution of luminous objects ${ }^{2}$, spiral galaxies are the most frequently observed morphological type, but not the most luminous compared to the giant ellipticals. In the Virgo cluster, the faintest spiral galaxies of type Sm differ from the brightest Sc by a factor $\sim 600$ in luminosity.
It has to be pointed out that in early times galaxies of type Sd-Sm-Im were initially classified as an irregular class of Irr I and Irr II. Later, de Vaucouleurs (1974) introduced the classes of $\mathrm{Sd}, \mathrm{Sm}$ and $\mathrm{Im}^{3}$. The " m " in Sm and Im is an abbreviation for Magellanic and indicates the similarities of these type to the Magellanic clouds (see

[^12]also Fig. 1.1 of Chapter 1).
Galaxies which do not belong to any of the above mentioned classes are then classified as irregular galaxies.
Looking at all the different morphological types in the Virgo cluster, but also in other clusters, it was realised that the spatial distribution of the galaxies is not homogeneous. Furthermore, studies showed that the fraction of blue (late-type) to red (early-type) galaxies is a function of the distance of the galaxy to the centre of the cluster, which is called the morphology-density relation (Dressler, 1980). Thus, the galaxies will change their appearance when moving through the dense cluster environment. For instance, Conselice et al. (2001) showed that several early-type dwarf galaxies are not an old cluster population, but originated from in-falling field galaxies and therefore (parts) of the dEs are not formed in-situ. But the question is, which morphological type is able to form an early-type dwarf galaxy? To answer this question, it is crucial to investigate the structural parameters of both, early- and late-type galaxies, and after that, the evolution with time of these parameters by the means of models and simulations.

In this chapter we analyse the photometric and structural properties of the Virgo latetype galaxy population. The results of this chapter will serve as input parameters for the evolutionary synthesis models of GALEV in Chapter 4. The GALEV-code models the time evolution of the galaxies, and will be therefore the foundation of our investigation of possible connections between early- and late-type galaxies.

### 3.2 The sample and data reduction

### 3.2.1 Sample selection

Using the Virgo Cluster Catalog (VCC) by Binggeli et al. (1985) we selected all latetype galaxies within the Virgo cluster up to magnitude of $\mathrm{m}_{\mathrm{B}} \leq 18.0 \mathrm{mag}$, where the VCC was found to be complete. This magnitude cutoff hence ensures that there is no bias due to undetected galaxies with a low surface brightness.

Applying a distance modulus of $\mathrm{m}-\mathrm{M}=31.09 \mathrm{mag}$ (according to a distance of $\mathrm{D}=$ 16.5 Mpc; Mei et al., 2007; Blakeslee et al., 2009) this magnitude limit corresponds to an absolute magnitude of $\mathrm{M}_{\mathrm{B}} \leq-13.09$ mag.
To cover the entire morphological sequence from early-type to late-type galaxies within the Virgo cluster, we supplement the results from Janz and Lisker (2008); Janz and Lisker (2009) on ETGs.

To decide whether a galaxy belongs to the Virgo cluster or is located in the background, a heliospheric velocity criterium of $-730 \leq \mathrm{v}_{\text {helio }} \leq 2990 \mathrm{~km} / \mathrm{s}$ was used (see appendix of Weinmann et al., 2011). All galaxies in this velocity range are treated as certain members. However, there are also galaxies with morphologies similar to Virgo galaxies but without any velocity information. These galaxies are handled as possible members. The original VCC was updated by Dr. T. Lisker using new velocities using the NASA/IPAC Extragalactic Database (NED) or SDSS (see Lisker et al., 2006b, and references therein).
Figure 3.1 shows the velocities and the corresponding distances of our sample taken from the GOLDMine data base (Gavazzi et al., 2003) ${ }^{4}$. Vertical lines correspond to the above mentioned velocity limits. There is one galaxy (VCC-0323) with a slightly larger velocity of $\mathrm{v}_{\text {helio }}=3002 \mathrm{~km} / \mathrm{s}$, which is also included in the analysis ${ }^{5}$. Galaxies


Figure 3.1: Velocities of the sample galaxies vs.their assumed distances. The heliocentric recessional velocities $v_{\text {helio }}$ are taken from the GOLDMine data base, as well as the informations about the distances. To compare the velocities with the study of Gavazzi et al. (1999), which used velocities with respect to the local group (LG), we transformed $\mathrm{v}_{\text {helio }}$ into $\mathrm{v}_{\mathrm{LG}}$ via $\mathrm{v}_{\mathrm{LG}}=\mathrm{v}_{\text {helio }}+220 \mathrm{~km} / \mathrm{s}$.
with nearby bright foreground or background objects and mergers were also removed from the sample, since an accurate analysis of the photometric properties would not be possible. There are several galaxies that are not or only partly covered by SDSS.

[^13]These galaxies were removed from the analysis as well.
The complete sample of late-type galaxies consists of 388 galaxies and a summary of the sample with all derived parameters can be found in Tab. 6.2 of the appendix.

### 3.2.2 Remarks on the distances

During our analysis, we realised that some galaxies were listed as Virgo members, but do not fulfil the above mentioned velocity criterium. In the past, the membership was first determined only by morphology without available distance measurements. With new velocity data from SDSS and other observations, the GOLDMine data were updated and the new velocities were added. It turns out that in some cases the GOLDMine data base has used the wrong spectral informations from the SDSS, as illustrated in Fig. 3.2. The left panel of Fig. 3.2 shows VCC-0703 which is classified as dI (GOLDMine: Im ), with a spectral derived distance of 95 Mpc . The cyan square indicates the location of the SDSS fibre, which was used to determine the distance. It is obvious that the spectrum used to derive the redshift, is that of the wrong galaxy and therefore also the distance is not the one of VCC-0703.
Another example of an issue with GOLDMine is VCC-0574 ( $\mathrm{D}_{\text {GOLDMine }}=307 \mathrm{Mpc}$ ), which is shown in the right panel of Fig. 3.2. In this case the situation is even worse. The coordinates of the dI VCC-0574 point to the galaxy with the cyan square, which clearly is not of type dI. Furthermore, there is a very faint galaxy in the north of the galaxy, which has the right morphological properties of a $\mathrm{dI}^{6}$. Moreover, it is also possible that the galaxy with the cyan square is a distant disc galaxy. Thus, it is very uncertain which galaxy actually is VCC-0574.

### 3.2.3 Classification

As a starting point we used the VCC classification of Binggeli et al. (1985). This step is important to obtain a first overview for the location of the different morphological types within the parameter-space of the derived photometric properties. In the course of this section the classification is reviewed and if necessary the galaxies are re-reclassified. One possible reason for a misclassification can be the usage of photometric plates from the Las Campanas survey. Figure 3.4 shows images of VCC-0979 from three different observations. On the left hand side of Fig. 3.4 the original photometric plate of the Las Campanas survey is shown, while the middle and right panels show the gri-

[^14]

Figure 3.2: Wrong identified distances in GOLDMine due to inaccurate determination of the SDSS fibres. Left: VCC-0702; Right: VCC-0574.

SDSS and Hubble Space Telescope (HST) (WFPC2-F606W filter at 5957Å) image, respectively. Using the photometric plate, VCC-0979 was classified as Sa galaxy in the original work of Binggeli et al. (1985). However, the images of SDSS and HST show significantly more details and reveal possible ongoing or recent star formation in the centre of the galaxy, as well as dust lanes. Thus, with upcoming new, deeper and more detailed observations the morphological types of the galaxies will need to be revised once again.
Figure 3.3 shows the number distribution of the different classes of the late-type galaxies in the sample for the classification by Binggeli et al. (1985) and additionally the new classes after the revision of this study. Additionally plotted is the distribution of possible members of the Virgo cluster relative to the total amount of galaxies. As expected dwarf irregular galaxies (dIs) are the main class of late-type galaxies in the Virgo cluster with 95 galaxies. Compared to the dIs, the BCDs are less abundant with a total amount of 37 galaxies.
Within the class of spiral galaxies, the Sc galaxies are by far the most numerous class (83 galaxies), followed by the Sa ( 35 galaxies). The other spiral classes of Sb (18), Sd (21) and Sm (26) are represented in roughly equal numbers. In the class of galaxies with an unknown morphology, almost all galaxies were reclassified as part of this study. For completeness sake, we also include spiral galaxies of the type Sa , despite knowing that Sa galaxies, due to their large bulge, are not able to transform into early-type


Figure 3.3: Number distribution of the different classes of the Virgo cluster. In each class the left bar corresponds to the classification of Binggeli et al. (1985) and the right one to the revision of our study.
dwarf galaxies. For the same reason of completeness, we also include two galaxies (VCC-0654 and VCC-1883) classified as "RSB0", which are S0 galaxies with a ring structure.
The group of star forming dwarf galaxies (BCDs and dIs) was already subject of the previous Chapter 2.

### 3.2.4 Data reduction

To analyse the optical properties of the late-type galaxy population of the Virgo cluster, we use the ugriz photometry with an effective exposure time of 54 s per filter (Gunn et al., 1998; Smith et al., 2002) from the Sloan Digital Sky Survey (SDSS) Data Release Five (DR5, Adelman-McCarthy et al., 2007). Since the SDSS does not point to each single galaxy, but operates in drift scan mode, the galaxies are not always covered by a single CCD image, which makes an alignment of several images necessary. Due to the high precision of the astrometry (Pier et al., 2003) the necessary alignment can


Figure 3.4: VCC-0979 observed with different telescopes. Left: photographic plate from the Las Campanas survey; Middle: gri-SDSS image; Right: HST image (WFPC2-F606W at $5957 \AA$ ).
be done automatically.
The SDSS pipeline provides measurements of the sky brightness, which turn out to be insufficient for our purpose. Therefore, we used the method of Lisker et al. (2007) for accurate sky subtraction. To determine the sky flux of the images, in a first step masks of all objects were created using the software SExtractor (Bertin and Arnouts, 1996). To avoid additional contamination of the sky due to too small masks, the masks were "blown up" through a gaussian kernel using the astronomical software package IRAF $^{7}$ (Tody, 1993). To increase the $\mathrm{S} / \mathrm{N}$, the $\mathrm{g}, \mathrm{r}$ and i images were co-added to a single image and in combination with the "blown up" mask a new mask was created, which was finally applied to the images in different filters to obtain the accurate sky values. As shown by Lisker et al. (2007) the sky is not constant over the entire image. To account for this non-constant sky, the above mentioned mask was applied to the images and the sky was determined within overlapping boxes of the size $201 \times 201$ pixels. This results in a master sky image of each filter that was finally subtracted from the images.
The sky-subtracted SDSS images were finally flux calibrated by using the information from the SDSS, and Galactic extinction was corrected by using the equation given by Schlegel et al. (1998). The zero point offsets of the $u$ and $z$ filter were also corrected (Oke and Gunn, 1983) with:

$$
\begin{align*}
& \mathrm{u}_{\mathrm{AB}}=\mathrm{u}_{\mathrm{SDSS}}-0.04 \mathrm{mag} \\
& \mathrm{z}_{\mathrm{AB}}=\mathrm{z}_{\mathrm{SDSS}}+0.02 \mathrm{mag} . \tag{3.1}
\end{align*}
$$

[^15]All the above mentioned data reduction steps were done automatically with the program virgocam, which was provided by T. Lisker.
The error of the fluxes $\Delta \mathrm{f}$ in each filter $\mathrm{x}(\mathrm{x}=\mathrm{u}, \mathrm{g}, \mathrm{r}, \mathrm{i}, \mathrm{z})$ was calculated by the formula given in Lisker et al. (2008):

$$
\begin{equation*}
\Delta \mathrm{f}_{\mathrm{x}}=\mathrm{f}_{\mathrm{x}} \cdot\left(\left(\frac{\sigma \cdot \sqrt{N_{p i x}}}{\mathrm{f}_{\mathrm{x}}}\right)^{2}+2 \cdot\left(\frac{0.002 \cdot \sigma \cdot \mathrm{~N}_{\mathrm{pix}}}{\mathrm{f}_{\mathrm{x}}}\right)^{2}+\left(10^{0.4 \cdot \Delta \varphi}-1\right)^{2}+\Theta \cdot\left(10^{0.4 \cdot 0.02}-1\right)^{2}\right)^{0.5} \tag{3.2}
\end{equation*}
$$

with

$$
\Theta= \begin{cases}0, & \text { for } \mathrm{g}, \mathrm{r}, \mathrm{i} \text { and } \mathrm{z}  \tag{3.3}\\ 1, & \text { for } \mathrm{u}\end{cases}
$$

and

$$
\Delta \varphi= \begin{cases}0.02, & \text { for } \mathrm{g}, \mathrm{r} \text { and } \mathrm{i}  \tag{3.4}\\ 0.03, & \text { for } \mathrm{u} \text { and } \mathrm{z}\end{cases}
$$

In Equation $3.2 \mathrm{~N}_{\mathrm{pix}}$, is the number of pixels within the aperture, $\mathrm{f}_{\mathrm{x}}$ the measured flux and $\sigma$ the noise level per pixel. $\Delta \varphi$ accounts for the uncertainties of the photometric calibration in the different filters. The last term in the formula accounts for the red leak in the u-band, resulting in an uncertainty of 0.02 mag. Equation 3.2 also includes the uncertainties regarding the determination of the Petrosian radius (see next Section 3.2.5). It was estimated that the error of the Petrosian determination has the same order as the sky uncertainties (Lisker et al., 2008).
With the propagation of uncertainty the errors of the colours were calculated. The following equation shows the error of the (g-i)-colour, which is the main colour index used in this study:

$$
\begin{equation*}
\Delta(\mathrm{g}-\mathrm{i})=-2.5 \cdot \log \left(1-\sqrt{\left(\frac{\Delta f_{g}}{f_{g}}\right)^{2}+\left(\frac{\Delta f_{i}}{f_{i}}\right)^{2}}\right) \tag{3.5}
\end{equation*}
$$

with the measured fluxes $\mathrm{f}_{\mathrm{g}}, \mathrm{f}_{\mathrm{i}}$ and the corresponding errors of Equation 3.2.

### 3.2.5 Data analysis

In the course of our data analysis, it was realised that the centre coordinates, given by the VCC, are inaccurate for some galaxies. Therefore, we determine the centre


Figure 3.5: Reduction of VCC-0030. Left: gri-SDSS image in false colours; Middle: corresponding mask for the image; Right: Combination of image and mask.
coordinates for the entire sample using the IRAF task ellipse (Jedrzejewski, 1987). This was done in several steps starting with an initial rough eye approach of the centre. With these centre coordinates a first determination of the Petrosian semi-major axis $\mathrm{a}_{\mathrm{p}}$ (Petrosian, 1976) for an elliptical aperture (Lotz et al., 2004), using the co-added SDSS gri-image of the galaxy, was performed by the program AMOR of T. Lisker. The derived parameters ( $\mathrm{x}_{0}, \mathrm{y}_{0}, \epsilon_{0}$ and $\mathrm{a}_{\mathrm{p}, 0}$ ) of this run were used as starting conditions for IRAF/ellipse with a fixed $\mathrm{a}_{\mathrm{p}}$. This analysis yielded the new centre coordinates $\left(\mathrm{x}_{\mathrm{c}}, \mathrm{y}_{\mathrm{c}}\right)$, ellipticity $(\epsilon)$ and position angle (p.a.) for each galaxy. Using these parameters, the Petrosian radius was determined again and in the last step the IRAF parameters ( $\mathrm{x}_{\mathrm{c}}, \mathrm{y}_{\mathrm{c}}, \epsilon$ and p.a.) were fixed and the final Petrosian radius was determined within the co-added gri-image. The derived parameters of the gri-image are used for all filters in the same manner and were used to measure the total flux and the half-light semi-major axis $a_{h l}$ within $2 a_{p}$ aperture for all five SDSS filters.
To avoid flux contamination by sources other than the galaxy, the above mentioned masks were used as a first approximation. To analyse the galaxy, the region around the relevant galaxy has to be unmasked, resulting in an image where all objects, except the sky and galaxy, are masked out. Small scale objects on top of the galaxy, which do not belong to the galaxy (e.g. stars), are not masked by the first mask. Therefore, these objects were visually inspected and are masked by hand. The left panel of Fig. 3.5 shows, in false colours, the image of VCC-0030, with the corresponding mask in the middle. The combination of the image and the mask is shown in the right panel of the figure.
With the final set of $\mathrm{x}_{\mathrm{c}}, \mathrm{y}_{\mathrm{c}}, \epsilon$, p.a. and $\mathrm{a}_{\mathrm{h}}$, we are able to perform our analysis within an aperture of $2 \mathrm{a}_{\mathrm{p}}$, yielding all relevant parameters, including the semi-major half light radius $\mathrm{a}_{\mathrm{hl}}$ in units of pixels. To obtain the effective radius $\mathrm{R}_{\text {eff }}$ in arcsec, the semi-major half light radius $\mathrm{a}_{\mathrm{hl}}$ has to be transformed via:

$$
\begin{equation*}
\mathrm{R}_{\mathrm{eff}}=0.396 \frac{\operatorname{arcsec}}{\text { pixel }} \cdot \mathrm{a}_{\mathrm{hl}} \cdot \sqrt{b / a} \tag{3.6}
\end{equation*}
$$

with the axis ratio (b/a) and the SDSS pixel scale of $0.396 \frac{\text { arcsec }}{\text { pixel }}$.
To convert $\mathrm{R}_{\text {eff }}$ from arcsec into parsec the following conversion was used

$$
\begin{equation*}
\mathrm{R}_{\mathrm{eff}}[\mathrm{pc}]=\mathrm{R}_{\mathrm{eff}}[\operatorname{arcsec}] \cdot \tan \left(\frac{1}{3600} \cdot \frac{\pi}{180}\right) \cdot \mathrm{D}[\mathrm{pc}] \tag{3.7}
\end{equation*}
$$

with the distance D of the object in parsec. In case of the Virgo cluster a constant distance of $\mathrm{D}=16.5 \mathrm{Mpc}$ was assumed (Mei et al., 2007).
The mean effective surface brightness $\langle\mu\rangle_{\text {eff }}$ of the galaxies was calculated by

$$
\begin{equation*}
\langle\mu\rangle_{\mathrm{eff}}=\mathrm{m}+2.5 \log \left(2 \pi \mathrm{R}_{\mathrm{eff}}^{2}\right) \tag{3.8}
\end{equation*}
$$

with the apparent magnitude $m$ and $R_{\text {eff }}$ in units of arcsec.
Since the above described parameters differ from filter to filter, it is reasonable to choose a reference band. If not stated otherwise, the r-band was used as the reference filter.

### 3.3 Results

In this section we present the results of the analysis of the late-type galaxy population within the Virgo cluster. In the first part, the spatial distribution of the galaxy within the Virgo cluster is described. In the second part of this section, the photometric properties of the sample are presented and furthermore the results are compared for different approaches for the distances. Finally, the morphological classes of the VCC are reviewed and if necessary the galaxies are reclassified by using all the informations of the following analysis.

### 3.3.1 Distribution within the Virgo cluster

The distribution of the galaxies within the Virgo cluster is shown in Fig. 3.6, which is divided into two panels for reasons of clearness. The position of M87 is marked with a black cross and defines the assumed centre of the Virgo cluster at a position of $\alpha=187.705907 \mathrm{deg}$ and $\delta=12.3911409 \mathrm{deg}(\alpha=12 \mathrm{~h} \mathrm{30min} 49.4 \mathrm{~s}$; $\delta=12 \mathrm{~d} 23 \mathrm{~min} 28.0 \mathrm{~s})$. Additionally, the sample of early-type galaxies from the study of Janz and Lisker (2008) is plotted as small black data points. Particular in the case of the spirals in the right hand side of Fig. 3.6, one can see a clustering of the morphological types into different clouds, which was already reported by the early studies of de Vaucouleurs (1961), Binggeli et al. (1985) and Binggeli et al. (1993). These clouds were also mentioned in the BCD Section 2 (see Tab. 2.5.2). Figure 3.7 shows the original map from Gavazzi et al. (1999) with the subdivision of the cluster into different regions. The coordinates of the galaxies correspond to the equinox of J1950, in contrast to the coordinates of this study with an equinox of J2000.
Comparing the distribution of the Virgo galaxies with the study of Gavazzi et al. (1999), the galaxies can be associated with the different clouds and sub-clusters. The distance of each galaxy within GOLDMine is based on the location within these cluster regions, thus a galaxy in the A-cluster has a distance of $\mathrm{D}=17 \mathrm{Mpc}$, whereas a galaxy in the M-cloud has a distance of $\mathrm{D}=32 \mathrm{Mpc}$. These distances will be from particular interest in the Sections 3.4.2 and 3.4.3, where we also discuss the implications of assigning distances this way.
With the exact centres of the galaxies it is possible to measure the projected distance to M87. The distances of the galaxies to M87 are calculated with

$$
\begin{equation*}
\mathrm{R}_{\mathrm{CC}}[\text { degree }]=\sqrt{\left(\alpha_{\mathrm{M} 87}-\alpha_{\text {galaxy }}\right)^{2}+\left(\delta_{\mathrm{M} 87}-\delta_{\text {galaxy }}\right)^{2}} \tag{3.9}
\end{equation*}
$$



Figure 3.6: The distribution of late-type galaxies in the Virgo cluster. For comparison the early-type galaxies (black dots) are plotted too. Different sub-regions of the cluster are visible and can be compared with Fig. 3.7.

This radius is also called the cluster centric radius $\mathrm{R}_{\mathrm{CC}}$.
Figure 3.8 shows the number distribution of the different types as a function of the distance to M87. $\mathrm{R}_{\mathrm{CC}}$ was normalised to a virial radius of $\mathrm{R}_{\text {Virial }}=1.5 \mathrm{Mpc}$ (McLaughlin, 1999).
As seen in Fig. 3.8, dIs have a relative mean distance to M87 of $\left\langle\frac{\mathrm{R}_{\mathrm{cc}}}{\mathrm{R}_{\text {Virial }}}\right\rangle=0.88 \pm 0.44$ with a small increase at $\frac{\mathrm{R}_{\mathrm{CC}}}{\mathrm{R}_{\mathrm{V} \text { iral }}}=1.5$. Compared to the dIs, the BCDs are further out at $\left\langle\frac{\mathrm{R}_{c c}}{\mathrm{R}_{\text {Virial }}}\right\rangle=0.99 \pm 0.49$, with two peaks in the number distribution at $\frac{\mathrm{R}_{\mathrm{Cc}}}{\mathrm{R}_{\text {Virial }}}=0.6$ and $\frac{\mathrm{R}_{\mathrm{cc}}}{R_{\text {Virial }}}=1.1$.
Due to the limited number of galaxies of the morphological type unknown, Amorphous and S, the statistic is quite poor. However, they all show a similar distribution of the cluster centric radius with an average of $\frac{\mathrm{R}_{\mathrm{CC}}}{\mathrm{R}_{\text {viral }}} \approx 0.75$.
Looking at the spiral galaxies in the right panel of Fig. 3.8, one can see that the most abundant type are the Sc galaxies, which are located at $\left\langle\frac{\mathrm{R}_{\mathrm{cc}}}{\mathrm{R}_{\text {Viral }}}\right\rangle=0.91 \pm 0.41$. The distribution of the Sa galaxies shows two peaks at $\frac{\mathrm{R}_{c c}}{\mathrm{R}_{\text {virial }}} \approx 0.4$ and $\frac{\mathrm{R}_{c c}}{R_{\text {Virial }}} \approx 1.25$, respectively. The Sb and Sc galaxies show a quite broad distribution with the Sb galaxies peaking at $\frac{R_{c c}}{R_{\text {virial }}}=0.75$. The morphological type of the Sm galaxies are also located in a small distance to M87 and show a broad distribution from $0 \leq \frac{R_{C C}}{R_{\text {Virial }}} \lesssim 2$. Table 3.1 shows the average radii from M87 for the different morphological types. It was found that galaxies of type $S$ have the smallest distance to the centre, while Sd


Figure 3.7: Map of the Virgo cluster, taken from Gavazzi et al. (1999), showing the different regions of the cluster. The position of M87 is indicated by a circle with superscript number 87. Large capitals label the region within the Virgo cluster and the size of the symbols corresponds to the distance modulus $\mu_{0}$ of the galaxies.
galaxies have the largest distance, as expected from the morphology-density relation.


Figure 3.8: Number distribution of different galaxy types regarding the cluster centric radius $\mathrm{R}_{\mathrm{CC}}$, which is normalised to a virial radius of $\mathrm{R}_{\text {virial }}=1.5 \mathrm{Mpc}$ and a step size of $0.2 \mathrm{R}_{\text {virial }}$. Note: the scale of the y -axis changed with the different morphological types.

Table 3.1: Average radii $\mathrm{R}_{\mathrm{CC}}$ to M 87 normalised to a virial radius of $\mathrm{R}_{\text {virial }}=1.5 \mathrm{Mpc}$.

| Type | $\left\langle\mathrm{R}_{\mathrm{CC}} / \mathrm{R}_{\text {virial }}\right\rangle$ | $\sigma$ |
| :---: | :---: | :---: |
| Sa | 0.76 | 0.46 |
| Sb | 1.01 | 0.49 |
| Sc | 0.91 | 0.41 |
| Sd | 1.08 | 0.49 |
| Sm | 0.86 | 0.51 |
| S | 0.67 | 0.31 |
| Amorph | 0.76 | 0.19 |
| BCDs | 0.99 | 0.49 |
| dI | 0.88 | 0.42 |

### 3.3.2 Photometric properties

In this section we present the results of our analysis. For a first approach, we will assume a constant distance of $\mathrm{D}=16.5 \mathrm{Mpc}$, and later on discussing the variation of the derived parameters if we instead use the distances from the GOLDMine data base. In Section 3.4.3, galaxies that are conspicuous in the diagrams are discussed in more detail.

Figures 3.9 to 3.19 show the results of late-type galaxies in the colour-magnitude diagram (CMD) and in the $\mathrm{R}_{\text {eff }}-\mathrm{M}_{\mathrm{r}}-\langle\mu\rangle_{\text {eff }}$-plane (together the parameter-space). For a first approach, we use the galaxy types of the VCC and the first revision of T. Lisker in the diagrams. Therefore, the given average values correspond to the revised version of the VCC by T. Lisker. Later on, the galaxy types are reclassified if their location within the $(\mathrm{g}-\mathrm{i})-\mathrm{M}_{\mathrm{r}}-\mathrm{R}_{\text {eff }}-\langle\mu\rangle_{\text {eff }}$ parameter-space does not fit to the bulk of the single galaxy types.
Parameters of BCDs are taken from Chapter 2 and are corrected for the flux of the starburst region. As in the previous plots, we show the results from Janz and Lisker (2008); Janz and Lisker (2009) for a sample of ETGs in the Virgo cluster as black data points. The vertical bars in the diagrams correspond to the $2 \sigma$ deviation of the ETG sample within a magnitude bin of $\Delta \mathbf{M}_{\mathrm{r}}=0.5 \mathrm{mag}$ or $\Delta\langle\mu\rangle_{\mathrm{eff}, \mathrm{r}}=0.5 \mathrm{mag} / \mathrm{arcsec}^{2}$. Green asterisks correspond to ETGs with a blue core (dE(bc)), which are characterised by ongoing or recent star formation(Lisker et al., 2006a).
In Tab. 6.2 of the appendix, the results of the analysis are summarised. The membership (ms) to the Virgo cluster was adapted from Binggeli et al. (1985) and was updated by T. Lisker with recent values from the literature (column [2]).
In the following we discuss the properties of each morphological galaxy type, as well as potential outliers in the diagrams.

## Sa

Galaxies classified as Sa are located in the region of $0.53<(\mathrm{g}-\mathrm{i})<1.23 \mathrm{mag},-21.91$ $<\mathrm{M}_{\mathrm{r}}<-15.11 \mathrm{mag}, 0.61<\mathrm{R}_{\mathrm{eff}}<5.26 \mathrm{kpc}$ and $18.91<\langle\mu\rangle_{\mathrm{eff}, \mathrm{r}}<23.70 \mathrm{mag} / \operatorname{arcsec}^{2}$ in the $(\mathrm{g}-\mathrm{i})-\mathrm{M}_{\mathrm{r}}-\mathrm{R}_{\text {eff }}-\langle\mu\rangle_{\text {eff }}$ parameter-space in Fig. 3.9 to 3.19.
Calculating the mean values of all parameters reveal a mean colour of $\langle(\mathrm{g}-\mathrm{i})\rangle=1.05 \pm$ 0.14 mag , a magnitude of $\left\langle\mathrm{M}_{\mathrm{r}}\right\rangle=-19.39 \pm 1.48 \mathrm{mag}$, an effective radius of $\left\langle\mathrm{R}_{\text {eff }}\right\rangle=2.04$ $\pm 1.07 \mathrm{kpc}$ and $\langle\langle\mu\rangle\rangle_{\mathrm{eff}, \mathrm{r}}=20.43 \pm 0.96 \mathrm{mag} / \operatorname{arcsec}^{2}$ (see Fig. 3.25). All quoted errors for the average values represent the $1 \sigma$ deviations.


Figure 3.9: Colour-magnitude diagrams of late- and early-type galaxies in the Virgo cluster. Black dots show the ETG sample of Janz and Lisker (2008); Janz and Lisker (2009). The vertical bars correspond to the $2 \sigma$ deviation within a magnitude bin of 0.5 mag. For clearness the CMD is divided into two parts. Left panel: CMD with spirals (Sa-Sd); Right panel: same CMD for Sm, SB, Amorphous, S and unknown galaxies. The CMD of the BCDs and the dIs can be found in Chapter 2.

At first glance one realises galaxies in the region between $1.1<(\mathrm{g}-\mathrm{i})<1.5 \mathrm{mag}$ and $-20.0<\mathrm{M}_{\mathrm{r}}<-17.2$ mag on the left hand side of Fig. 3.9 which are slightly redder than the population of ETGs. One may speculate that these galaxies are inclined ones. However, a detailed inspection of the co-added gri-SDSS images of the Sa galaxies in this region showed that there is no general trend of being edge-on. In this case it is worth to investigate the shifts of the magnitude ${ }^{8}$ introduced by different distance approaches, what will be done in Section 3.4.3.
Sa galaxies have an average axis ratio of $\langle(b / a)\rangle=0.56 \pm 0.20$ with a broad distribution from almost circular galaxies with $(\mathrm{b} / \mathrm{a}) \approx 1$ down to $(\mathrm{b} / \mathrm{a}) \approx 0.2$, indicating very elliptical shapes (Fig. 3.10). Assuming that face-on Sa galaxies are almost round with an axis ratio of $(b / a) \approx 1$ and also assuming that they are very thin in the zdirection, it became clear that several Sa galaxies of the sample have to be inclined.

[^16]With increasing inclination angle the path length through the disc increases and in turn increasing the dust reddening. Therefore, also the magnitude and the corresponding colours are influenced. From this we conclude that the intrinsic colour distribution, corrected for inclination, is somewhat bluer than the observed one.


Figure 3.10: Histogram of the axis ratio of Sa galaxies. The mean axis ratio is indicated by the black line and amounts $\langle(b / a)\rangle=0.56 \pm 0.20$.

## Sb

In the region of $0.47<(\mathrm{g}-\mathrm{i})<1.3 \mathrm{mag},-21.40<\mathrm{M}_{\mathrm{r}}<-17.41 \mathrm{mag}, 0.76<\mathrm{R}_{\text {eff }}<$ 5.3 kpc and $18.78<\langle\mu\rangle_{\mathrm{eff,r}}<22.38 \mathrm{mag} / \operatorname{arcsec}^{2}$ of Figs. 3.9 to 3.19 the Sb galaxies are located. Looking at the CMD in Fig. 3.9, one finds Sb galaxies in the region of ( $\mathrm{g}-\mathrm{i}$ ) > 1.0 mag and $\mathrm{M}_{\mathrm{r}}>-20.0 \mathrm{mag}$, which are, from a visual inspection, invariably inclined galaxies. Except for the galaxies in this region, the other Sb galaxies follow an almost linear correlation in the way that with increasing colours the galaxies getting brighter. In the $\mathrm{R}_{\mathrm{eff}}-\langle\mu\rangle_{\mathrm{eff}, \mathrm{r}}$-diagram of Fig. 3.18, we found that Sc galaxies with a larger effective radius are also brighter.
No obvious outliers can be found, except for VCC-0267, which has a slightly lower $\langle\mu\rangle_{\text {eff,r }}$ compared to the other Sb galaxies in Fig. 3.12 and 3.19.
The following average values and $1 \sigma$ deviations were calculated for the Sb galax-
ies: $\langle(\mathrm{g}-\mathrm{i})\rangle=0.98 \pm 0.25 \mathrm{mag},\left\langle\mathrm{M}_{\mathrm{r}}\right\rangle=-19.46 \pm 1.25 \mathrm{mag},\left\langle\mathrm{R}_{\mathrm{eff}}\right\rangle=2.13 \pm 1.24 \mathrm{kpc}$ and $\left\langle\langle\mu\rangle_{\text {eff }, \mathrm{r}}\right\rangle=20.42 \pm 0.88 \mathrm{mag} / \operatorname{arcsec}^{2}$. The average axis ratio was measured to $\langle(\mathrm{b} / \mathrm{a})\rangle=0.55 \pm 0.23$. However, Fig. 3.11 shows a double peak in the histogram at $(b / a)=0.2$ for highly inclined Sb galaxies and at $(b / a)=0.7$. This double-peak distribution may be explained by the inclination angles. A face-on spiral is easily classified according its bulge-disc ratio and pitch-angle of the spiral arms. However, these features became less obvious in case of increasing inclination.


Figure 3.11: Same as Fig. 3.10, but for the Sb galaxies. $\langle(b / a)\rangle=0.55 \pm 0.23$.

## Sc

In Figs. 3.9 to 3.19 the region of $0.39<(\mathrm{g}-\mathrm{i})<1.21 \mathrm{mag},-22.07<\mathrm{M}_{\mathrm{r}}<-15.05$ $\mathrm{mag}, 0.67<\mathrm{R}_{\mathrm{eff}}<7.62 \mathrm{kpc}$ and $19.72<\langle\mu\rangle_{\mathrm{eff}, \mathrm{r}}<23.16 \mathrm{mag} / \operatorname{arcsec}^{2}$ is populated with Sc-type galaxies.
Inclined Sc galaxies can be found in the region of $(\mathrm{g}-\mathrm{i})>1.0 \mathrm{mag}$ and $\mathrm{M}_{\mathrm{r}}>-19.0$ mag. The rest of the Sc galaxies covers a region in the CMD from $0.39<(\mathrm{g}-\mathrm{i})<1.2 \mathrm{mag}$ and $-22.0<\mathrm{M}_{\mathrm{r}}<-16.0 \mathrm{mag}$, with one outlier (VCC-0989) at $(\mathrm{g}-\mathrm{i})=0.56 \mathrm{mag}$ and $M_{r}=-15.05 \mathrm{mag}$, which will be discussed in Section 3.4.3.
Calculating the mean values of all parameters reveal: $\langle(\mathrm{g}-\mathrm{i})\rangle=0.76 \pm 0.20 \mathrm{mag}$, $\left\langle\mathrm{M}_{\mathrm{r}}\right\rangle=-18.51 \pm 1.31 \mathrm{mag},\left\langle\mathrm{R}_{\mathrm{eff}}\right\rangle=1.97 \pm 1.26 \mathrm{kpc}$ and $\left\langle\langle\mu\rangle_{\mathrm{eff}, \mathrm{r}}\right\rangle=21.20 \pm 0.80$


Figure 3.12: Effective radius $\mathrm{R}_{\text {eff }}$ vs. the mean effective surface brightness $\langle\mu\rangle_{\text {eff }, ~}$ in r-band. Left panel: spirals (Sa-Sd); Right panel: same for Sm, SB, Amorphous, S and unknown galaxies.
mag/arcsec ${ }^{2}$.
Like the Sb galaxies, also the Sc 's show double peak in the histogram for the axis ratio (Fig. 3.13). However, the region between the two peaks is small and thus may disappears when different bin sizes and steps were applied. The average axis ratio over the entire range in the histogram of the Sc galaxies is $\langle(b / a)\rangle=0.54 \pm 0.25$.

## Sd

Galaxies classified as Sd galaxies are located in the region of $0.21<(\mathrm{g}-\mathrm{i})<0.89 \mathrm{mag}$, $-18.86<\mathrm{M}_{\mathrm{r}}<-14.48 \mathrm{mag}, 0.62<\mathrm{R}_{\mathrm{eff}}<3.39 \mathrm{kpc}$ and $21.23<\langle\mu\rangle_{\mathrm{eff}, \mathrm{r}}<24.27 \mathrm{mag}$ $/ \operatorname{arcsec}^{2}$.
Compared to the Sc galaxies, the Sd galaxies in Fig. 3.9 have bluer ( $\mathrm{g}-\mathrm{i}$ )-colours with a minimum at $(\mathrm{g}-\mathrm{i})=0.21 \mathrm{mag}$. At a given $(\mathrm{g}-\mathrm{i})$-colour of $(\mathrm{g}-\mathrm{i}) \approx 0.62 \mathrm{mag}$ two galaxies have relatively faint magnitudes $\left(\mathrm{M}_{\mathrm{r}}>-15.1 \mathrm{mag}\right)$ compared to the other Sc galaxies at the same (g-i)-colour. These two galaxies are VCC-0132 and VCC-1605 and are also conspicuous in Figs. 3.18 to 3.19, particular VCC-1605 stands out with a very small $\mathrm{R}_{\text {eff }}$ compared to the other Sc galaxies.


Figure 3.13: Same as Fig. 3.10, but for the Sc galaxies. $\langle(b / a)\rangle=0.54 \pm 0.25$.

The Sd galaxies have an average ( $\mathrm{g}-\mathrm{i}$ )-colour of $\langle(\mathrm{g}-\mathrm{i})\rangle=0.58 \pm 0.17 \mathrm{mag},\left\langle\mathrm{M}_{\mathrm{r}}\right\rangle=$ $-16.88 \pm 1.03 \mathrm{mag},\left\langle\mathrm{R}_{\text {eff }}\right\rangle=1.66 \pm 0.68 \mathrm{kpc}$ and $\left\langle\langle\mu\rangle_{\text {eff }, r}\right\rangle=22.61 \pm 0.66 \mathrm{mag} / \mathrm{arcsec}^{2}$. Figure 3.14 shows the histogram of the distribution of the axis ratio (b/a), with an average of $\langle(\mathrm{b} / \mathrm{a})\rangle=0.59 \pm 0.29$. The axis ratio has a broad distribution with one peak at $(b / a)=0.8$.

## Sm

In the right panels of Figs. 3.9 to 3.19, we show the results of the Sm galaxies. The Sm galaxies cover a region in the parameter-space of $0.24<(\mathrm{g}-\mathrm{i})<0.90 \mathrm{mag},-19.52$ $<\mathrm{M}_{\mathrm{r}}<-14.22 \mathrm{mag}, 0.48<\mathrm{R}_{\text {eff }}<2.62 \mathrm{kpc}$ and $20.14<\langle\mu\rangle_{\text {eff,r }}<24.02 \mathrm{mag} / \operatorname{arcsec}^{2}$. There are three galaxies (VCC-1554, VCC-1575 and VCC-1686) which are conspicuous in the diagrams. However, only VCC-1554 shows obvious differences from the rest of the Sm population, while the other two galaxies are at the "boundaries" of the Sm distribution. The "boundaries" are described by the margins of the bulk of the population in the parameter-space.

Averaging all Sm galaxies we obtain: $\langle(\mathrm{g}-\mathrm{i})\rangle=0.55 \pm 0.14 \mathrm{mag},\left\langle\mathrm{M}_{\mathrm{r}}\right\rangle=-16.42 \pm$ $1.30 \mathrm{mag},\left\langle\mathrm{R}_{\text {eff }}\right\rangle=1.21 \pm 0.64 \mathrm{kpc}$ and $\left\langle\langle\mu\rangle_{\text {effr }}\right\rangle=22.24 \pm 0.90 \mathrm{mag} / \mathrm{arcsec}^{2}$.
The average axis ratio is $\langle(\mathrm{b} / \mathrm{a})\rangle=0.54 \pm 0.14$ and the distribution for the entire Sm


Figure 3.14: Same as Fig. 3.10, but for the Sd galaxies. $\langle(b / a)\rangle=0.59 \pm 0.29$.
sample can be found in Fig. 3.15.


Figure 3.15: Same as Fig. 3.10, but for the Sm galaxies. $\langle(b / a)\rangle=0.54 \pm 0.14$.

## Unknown

Most of the initially "unknown" galaxies were reclassified by T. Lisker as irregular galaxies. To check this reclassification, all "unknowns" from the original VCC are used and were checked again in the parameter-space.
Galaxies classified as "unknown" in the VCC cover almost the same region in the CMD as the dIs, therefore, the revision by T. Lisker seems plausible. However, in the case of VCC-0020 it is not easy to decide which galaxy was classified in the initial VCC. In the image of the VCC-0020 there are three galaxies close together. The most prominent one has blue colours and could be classified as dI regarding its morphology, while the others have too red colours. The coordinates, given by GOLDMine for VCC-0020, point exactly to a galaxy which is relatively red. Due to the insufficient coordinates, we cannot confirm its VCC classification, and therefore VCC-0020 should not be taken into account for the remaining analysis.
After the first revision by T. Lisker, there is only one galaxy left over in the class of the "unknown" galaxies. This galaxy (VCC-0113) has very similar properties in the parameter-space as the BCDs in Figs. 3.9 and 3.18. However, the surface brightness is too low compared to the BCDs and LSB-components (see Figs. 3.12 and 3.19). Furthermore, the optical SDSS image of VCC-0113 shows some possible hints for interaction, noticeable by an extended halo with some kinds of loops or shells.

## S

The class of "S" galaxies has be introduced for galaxies, showing a spiral structure but which cannot be associated with a special subtype (e.g. Sa or Sm ). Looking at the class of " S " galaxies, one realises that they cover a range in magnitude of $\Delta \mathrm{M}_{\mathrm{r}}=3 \mathrm{mag}$ and $\Delta\langle\mu\rangle_{\mathrm{eff}, \mathrm{r}}=4 \mathrm{mag} / \operatorname{arcsec}^{2}$, locating them in the regime of both giants and dwarfs. Figure 3.9 shows two galaxies (VCC-0135 and VCC-0213), which are offset from the rest with a magnitude $\mathrm{M}_{\mathrm{r}}>-16 \mathrm{mag}$ and were already discussed in Section 2.4.8. They will also be discussed in Section 3.4.3. Averaging over the entire parameterspace yields a (g-i)-colour of $\langle(\mathrm{g}-\mathrm{i})\rangle=0.87 \pm 0.14 \mathrm{mag},\left\langle\mathrm{M}_{\mathrm{r}}\right\rangle=-17.26 \pm 1.53 \mathrm{mag}$, $\left\langle\mathrm{R}_{\text {eff }}\right\rangle=1.18 \pm 0.62 \mathrm{kpc}$ and $\left\langle\langle\mu\rangle_{\text {effr, }}\right\rangle=21.36 \pm 1.29 \mathrm{mag} / \mathrm{arcsec}^{2}$.
The distribution of the axis ratio is shown in Fig. 3.16. It has a discrete distribution with several peaks and an average axis ratio of $\langle(b / a)\rangle=0.57 \pm 0.27$. However, this finding should be handled with care, since there are only a few galaxies in this morphological class, resulting a poor statistic.


Figure 3.16: Same as Fig. 3.10, but for the $S$ galaxies. $\langle(b / a)\rangle=0.57 \pm 0.27$.

## Amorphous

First introduced by Sandage and Brucato (1979), the Amorphous galaxies are "a redefinition of the standard Irr II type" ${ }^{9}$. Here the group of Irr I galaxies shows some hints of spiral arms (or more commonly: a organised structure), while the Irr II galaxies do not show any organised structure (Carroll and Ostlie, 2006). Thus, one would expect galaxies with irregularities. However, the gri-images of the Amorphous galaxies show discy galaxies with a wide range in colour. Comparing the classification of the original VCC and the data from GOLDMine also shows differences. Four out of five Amorphous galaxies are classified as "unknown" in GOLDMine and one galaxy (VCC-1675) is classified as "Pec". The classification of GOLDMine may be the better choice, since the amorphous class is somewhat loosely defined.
In Figs. 3.9 to 3.19, the Amorphous galaxies can be found in the region of 0.65 $<(\mathrm{g}-\mathrm{i})<0.90 \mathrm{mag},-19.84<\mathrm{M}_{\mathrm{r}}<-17.34 \mathrm{mag}, 0.82<\mathrm{R}_{\text {eff }}<1.7 \mathrm{kpc}$ and $18.92<\langle\mu\rangle_{\text {eff,r }}<22.10 \mathrm{mag} / \operatorname{arcsec}^{2}$ with average values of: $\langle(\mathrm{g}-\mathrm{i})\rangle=0.78 \pm 0.11$ $\mathrm{mag},\left\langle\mathrm{M}_{\mathrm{r}}\right\rangle=-18.48 \pm 1.05 \mathrm{mag},\left\langle\mathrm{R}_{\text {eff }}\right\rangle=1.21 \pm 0.41 \mathrm{kpc},\left\langle\langle\mu\rangle_{\text {eff }, \mathrm{r}}\right\rangle=20.39 \pm 1.24$ $\mathrm{mag} / \operatorname{arcsec}^{2}$ and $\langle(b / a)\rangle=0.52 \pm 0.07$. The histogram of $(\mathrm{b} / \mathrm{a})$ is shown in Fig. 3.17 and has a very narrow distribution around the average, partly due to the small amount

[^17]of galaxies with this classification.


Figure 3.17: Same as Fig. 3.10, but for the Amorphous galaxies. $\langle(b / a)\rangle=0.52 \pm 0.07$.


Figure 3.18: Effective radius $R_{\text {eff }}$ vs. absolute magnitude $M_{r}$ in r-band. Left panel: spirals (Sa-Sd); Right panel: same for Sm, SB, Amorphous, S and unknown galaxies.


Figure 3.19: Absolute magnitude $\mathrm{M}_{\mathrm{r}}$ vs. the mean effective surface brightness $\langle\mu\rangle_{\mathrm{eff}, \mathrm{r}}$ in r-band. Left panel: spirals (Sa-Sd); Right panel: same for Sm, SB, Amorphous, S and unknown galaxies.

### 3.3.3 Colour-colour diagrams

Since all galaxies were analysed in all SDSS filter bands from u to z , also other colour informations are available. Figures 3.20, 3.21 and 3.22 for example, show the distance independent colour-colour diagram of all galaxies in our sample. Used colours are the (u-g)-colour, which is most sensitive to the age due to the UV flux of the young stellar population and the (i-z)-colour, most sensitive to the metallicity.
In the sample of ETGs, there are some galaxies that are relatively offset from the rest of the ETG population. This pertains to VCC-1039 and VCC-1719 with an (i-z)-colour of 0.84 mag and 1.15 mag , respectively. In the ( $\mathrm{u}-\mathrm{g}$ )-colours, VCC-0091 is offset from the rest with an $(\mathrm{u}-\mathrm{g})$-colour of $(\mathrm{u}-\mathrm{g})=3.25 \mathrm{mag}$. In the figures a cutoff at $(\mathrm{u}-\mathrm{g})=2.5$ mag and (i-z)=0.6 mag was applied, thus the above mentioned galaxies are not visible in the diagrams.
In Fig. 3.20, the Sa to Sd galaxies are displayed and additionally the ETGs are plotted (black dots). Among the Sa galaxies there are three galaxies which are very conspicuous in the diagram with colours of $(\mathrm{u}-\mathrm{g})<1.25 \mathrm{mag}$ and $(\mathrm{i}-\mathrm{z})<0.04 \mathrm{mag}$, corresponding to VCC-0015, VCC-1358 and VCC-1933. The rest of the Sa galaxies have colours of $(\mathrm{u}-\mathrm{g})>1.3 \mathrm{mag}$ and $(\mathrm{i}-\mathrm{z})>0.1 \mathrm{mag}$
One of the Sc galaxy (VCC-0865) has a very red (i-z)-colour of ( $\mathrm{i}-\mathrm{z}$ ) $=0.25 \mathrm{mag}$ compared to the other Sc galaxies. The other Sc galaxies have at given ( $\mathrm{u}-\mathrm{g}$ )-colours redder (i-z)-colour compared to the ETGs, but cover almost the same range in the ( $u$ g )-colours of the ETGs.
In the class of the Sd galaxies, there is one galaxy (VCC-0132) that totally offset from the rest with extreme $(\mathrm{i}-\mathrm{z})$-colours of $(\mathrm{i}-\mathrm{z})=-0.48 \mathrm{mag}$, which is only reached by the dIs.
The Sm galaxies are displayed in Fig. 3.21 and are found at (u-g)-colours of (i-z) $<1.24$ mag, with a large spread in the (i-z)-colours.
The loosely classified S galaxies, have the same coverage as the Sb and Sc galaxies. Furthermore, looking at the distribution of the Sb galaxies in Fig. 3.20, one can find a kind of underrepresentation of Sb galaxies between $1.25<(\mathrm{u}-\mathrm{g})<1.45 \mathrm{mag}$ and (i -z ) $<0.16$ mag, which is filled by three of the $S$ galaxy (Fig. 3.21), while the other S galaxies cover the same region as the Sb galaxies. However, this is only a weak indication to the exact morphological type of the S galaxies, since the same colours do not necessarily mean that the galaxies have the same morphology (structure). One of the S galaxies $($ VCC-0320) can be found at $(\mathrm{u}-\mathrm{g})=-1.38$ mag and $(\mathrm{i}-\mathrm{z})=-0.39$ mag. Among the spiral galaxies only the Sb galaxies have comparable red (i-z)-colours, what
could be therefore again hint to the exact morphological type.
The only "unknown" galaxy (VCC-0113) is located at ( $\mathrm{u}-\mathrm{g}$ ) $=-0.82 \mathrm{mag}$ and $(\mathrm{i}-\mathrm{z})=-$ 0.46 mag , which is only reached by the dI galaxies (cf. Fig. 3.22).

The colour-colour diagram of the BCDs and dIs is shown in Fig. 3.22. In contrast to the analysis of Chapter 2, the BCDs are not divided into the contribution of the starburst and LSB-component, due to the low $\mathrm{S} / \mathrm{N}$ of the z -filter, what would be resulted in large photometric errors. Therefore, we use the total colours of the BCDs, instead of the colours of the LSB-component. The dIs show a large scatter in the colour-colour diagram and together with the BCDs, they tend to have redder (i-z)-colours at a given (u-g)-colour than the bulk of the dEs. However, also the dEs show a larger scatter in the diagram.


Figure 3.20: Colour-colour diagram of the late-type sample. Displayed are the spirals and the ETGs.


Figure 3.21: Same as Fig. 3.20, but for the Sm, S Amorphous and unknown galaxies.


Figure 3.22: Same as Fig. 3.20, but for the BCDs and dIs. For the BCDs the total values are used, since the $\mathrm{S} / \mathrm{N}$ of the LSB-component in z is too low.

### 3.3.4 The luminosity function

The luminosity function (LF) describes the number of galaxies dN per magnitude bin dM. LF can be approximated by an analytical function, given by Schechter (1976):

$$
\begin{equation*}
\mathrm{N}(\mathrm{M}) \mathrm{dM}=\phi^{*} \cdot 10^{-0.4(\alpha+1) \mathrm{M}} \cdot \mathrm{e}^{\left.10^{\left(0.4\left(\mathrm{M}^{*}-\mathrm{M}\right)\right.}\right)} \mathrm{dM} \tag{3.10}
\end{equation*}
$$

with the free parameters of the slope $\alpha$, the normalisation $\phi^{*}$ and the characteristic magnitude $\mathrm{M}^{*}$.
Figure 3.23 shows the LF of the early-type and late-type population of the Virgo cluster in the range of -22.0 to -13.0 mag with a bin width of 0.5 mag . Errors were estimated by

$$
\begin{equation*}
\Delta N(M)=\sqrt{N(M)} \tag{3.11}
\end{equation*}
$$

The black squares correspond to the overall LF of the entire Virgo cluster population. The green asterisks and red triangles indicate the distribution of the late-type and early-type galaxies, respectively. The solid line corresponds to the overall LF using Equation 3.10. The error weighted fit was performed with a non-linear least-square Marquardt-Levenberg algorithm. The slope $\alpha$ was found to be

$$
\begin{equation*}
\alpha=-1.27 \pm 0.02 \tag{3.12}
\end{equation*}
$$

and the characteristic magnitude has a value of

$$
\begin{equation*}
\mathrm{M}_{\mathrm{r}}^{*}=-21.58 \pm 0.25 \mathrm{mag} . \tag{3.13}
\end{equation*}
$$

Additional to the Schechter-function, also a linear fit (straight line in Fig. 3.23) to the data in the interval $-14.4 \leq \mathrm{M}_{\mathrm{r}} \leq 19.0 \mathrm{mag}$ was applied, with a slope of

$$
\begin{equation*}
\alpha_{\operatorname{lin}}=-1.30 \pm 0.10 \tag{3.14}
\end{equation*}
$$

The slopes $\alpha$ and $\alpha_{\text {lin }}$ are in good agreement with Rines and Geller (2008) also using SDSS r-band data and it is also in good agreement with Blanton et al. (2001) for their sample of 11275 galaxies from the SDSS. However, our values of $\alpha$ are slightly lower compared to other studies of the Virgo cluster by Trentham and Hodgkin (2002) and Sabatini et al. (2003), who found a slope of $\alpha \approx-1.6$ in the B-band. The SDSS analysis of Krywult (2009) on Abell clusters with richness classes (Abell et al., 1989) of $\mathrm{R} \geq 2$ (rich cluster) and $\mathrm{R} \leq 1$ (poor cluster) found quite different results. In the g-band they obtained slopes of Schechter-function fits of $\alpha=-0.70$, and even lower
values ( $\alpha \approx 0.55$ ) for the other SDSS filters for the Abell rich clusters.
Trentham and Hodgkin (2002) investigated in their study whether the parameters of the LF depend on the used filter. Therefore, all five filters are used to derive the parameters of the LF and the results of the analysis are shown in Tab. 3.2. The results of $\alpha$ for the different filters show quite small deviations. When taken the errors of $\alpha$ into account, the differences are negligible.
The data of this thesis were partly used in Lieder et al. (2012) to determine the LF of the Virgo core region. A transformation of our r-band data to the V-band magnitude, using the method of Jester et al. (2005), and a linear fit to the Lieder et al. (2012) sample in the interval $-18.0 \leq \mathrm{M}_{\mathrm{V}} \leq-13.0$ mag reveals a slope of $\alpha=-1.50 \pm 0.05$. Similar to the studies of Trentham and Hodgkin (2002) and Sabatini et al. (2003), this slope is steeper than ours.


Figure 3.23: Luminosity function of Virgo galaxies. Plotted are late-type, earlytype galaxies on its own (green and red symbols) and the total LF. The red line corresponds to a Schechter-function and the blue line to a linear fit in the interval $-14.4 \leq \mathrm{M}_{\mathrm{r}} \leq 19.0$ mag.

Figure 3.24 shows the same LF, but subdivided into the contribution of the different morphological types, using a bin size of 0.25 mag . In both panels of Fig. 3.24 the distribution of the ETGs and $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$ is shown with a different y -scale compared to the other panels. The ETGs show a gradually increase of the LF with a turnover at

Table 3.2: Parameters of the LF for different filters.

| Filter | $\alpha$ | $\Delta \alpha$ | $\mathrm{M}^{*}$ | $\Delta \mathrm{M}^{*}$ |
| :---: | :---: | :---: | :---: | :---: |
| u | -1.371 | 0.091 | -20.049 | 0.519 |
| g | -1.337 | 0.037 | -21.593 | 0.299 |
| r | -1.297 | 0.037 | -22.023 | 0.324 |
| i | -1.276 | 0.028 | -22.367 | 0.265 |
| z | -1.295 | 0.027 | -22.863 | 0.374 |

$\mathrm{M}_{\mathrm{r}} \approx-14 \mathrm{mag}$.
The Sd and Sm galaxies have a relative similar distribution, even though the Sm galaxies are less abundant. Among the spirals, the Sc galaxies have the broadest distribution from almost $\mathrm{M}_{\mathrm{r}}=-23 \mathrm{mag}$ down to $\mathrm{M}_{\mathrm{r}}=-15 \mathrm{mag}$. The bulk of the Sc galaxies can be found around a magnitude of $\mathrm{M}_{\mathrm{r}}=-19$ mag. As already mentioned in Section 3.3.2, Sa galaxies show a large range in magnitude, what is also visible in Fig. 3.24 with a similar range like the Sc galaxies. Inspecting the star forming dwarfs in this figure, one realises that the BCDs have a narrower distribution than the dIs. There are dIs with slightly higher magnitudes than the BCDs, but there are also dIs with lower magnitudes. The LFs of the loosely classified " S " galaxies, together with the unknown and Amorphous galaxies, suffer from the poor number of galaxies within these morphological classes. Therefore, any interpretation has to be handled with care. Also obvious in Fig. 3.24 is the overlap between regions of generally classified dwarfs and spiral galaxies. For instance, the spiral types of Sd and Sm share a large overlap region with the BCDs and dIs, which makes a simple magnitude division in dwarf and non-dwarfs questionable.


Figure 3.24: Luminosity functions of the different morphological types. Note: the plots of the ETGs and $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$ have different y -scales than the other plots.

### 3.4 Discussion

### 3.4.1 Comparison of the different types

The following Figs. 3.25 and 3.26 show the averages and the $1 \sigma$ deviation of the derived parameters of the different morphological types. The derived parameters of our galaxies are given for every morphological type of the original VCC (red bars) and of the revision by T. Lisker (red bars). Additionally to the late-type galaxies, the averages of the complete ETG population are given as well. Furthermore, the sample of ETGs is subdivided into the contribution of the $\mathrm{dEs}\left(\mathrm{M}_{\mathrm{r}}<-18 \mathrm{mag}\right)$ and the $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$

The difference between the VCC and the revision are minor for the magnitude $M_{r}$, ( $\mathrm{g}-\mathrm{i}$ )-colours and $\langle\mu\rangle_{\text {eff }, \mathrm{r}}$. However, the differences for $\mathrm{R}_{\text {eff }}$ and the axis ratio ( $\mathrm{b} / \mathrm{a}$ ) are more prominent, especially for the Sb and Sd galaxies.
In the next step, the morphological types are compared with each other. The parameters of the Sa and Sb galaxies do not show a large deviation from each other, while they differ from the other spirals of the type $\mathrm{Sc}, \mathrm{Sd}$ and Sm . The Sd and Sm galaxies also show similarities in $\mathrm{M}_{\mathrm{r}}$ and $(\mathrm{g}-\mathrm{i})$. However, the Sd galaxies have on average a larger $\mathrm{R}_{\text {eff }}$ and a lower $\langle\mu\rangle_{\text {eff,r }}$ compared to the Sm galaxies.
The unknown galaxies of the original VCC show a similar behaviour like the dIs, which again supports the reclassification of most of them to the morphological class of the dIs by T. Lisker. Owing to the fact that the revision only contains one unknown galaxy, no error bars are given in the diagrams.
Since the main aim of this study is to find the possible progenitors of the early-type dwarf galaxies, it is worth to compare the derived parameters of today's late-type galaxies with the ones of the dEs. To obtain a more sophisticated insight into the possible connections, we will also study the photometric evolution of the late-type population by means of evolutionary synthesis models in Chapter 4.
One of the main findings of Chapter 2 was that the LSB-components of the BCDs and the compact dEs cover almost the same region in the $\mathrm{M}_{\mathrm{r}}-\mathrm{R}_{\text {eff }}-\langle\mu\rangle_{\mathrm{eff,r},}$-plane, what can also be seen in Figs. 3.25 and 3.26. It is also obvious that the LSB-components are on average bluer and more compact than the dEs. However, the effective radius in units of kpc of the LSB-components is also influenced by the different distance approaches, what will be discussed in more detail in Section 3.4.2.
Comparing the results of the other morphological types with the ETGs, one can see obvious differences, but also some similarities. However, when comparing the
averages one should also keep in mind the $1 \sigma$ deviation of the results.
As expected, the averaged parameters of the Sa and Sb galaxies significantly differ from the ones of the dEs. On average the magnitude $M_{r}$ and mean surface brightness $\langle\mu\rangle_{\text {eff,r }}$ of the Sa and Sb galaxies are too high and the effective radius $\mathrm{R}_{\text {eff }}$ is too large compared to the ETGs. Looking at Fig. 3.9, it is visible that most of the spiral galaxies at a given (g-i)-colour are brighter than the ETGs. Furthermore, due to the active star formation of the spiral galaxies, expectedly they have bluer colours than the ETGs. In Figs. 3.12 and 3.19 the spirals within a $2 \sigma$ deviation cover the same region as the ETGs. Galaxies with $\mathrm{M}_{\mathrm{r}}<-20 \mathrm{mag}$ are slightly fainter with increasing radius compared to the ETGs in Fig. 3.18.
As already seen in Fig. 3.24, there is no clear separation between the dwarf galaxies and spirals. This motivates to ask what really defines a dwarf galaxy? The definition by a certain magnitude holds overlaps with spirals as well as in surface brightness and effective radius.


Figure 3.25: Distribution of the averaged results of the late-type galaxies in the Virgo cluster. The red bars correspond to the morphological types derived by the initial VCC, while the black bars show the results for the morphological types classified by T. Lisker. Additionally shown are the results of the early-type galaxies. The horizontal red and blue lines indicate the averages of the dEs and $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$, respectively.


Figure 3.26: Same as Fig. 3.25, but for the axis ratio b/a.


Figure 3.27: Normalised distribution of the axis ratio for the ETGs, dEs and dE(bc)s of the sample of Janz and Lisker (2008); Janz and Lisker (2009) with corresponding averages of $\langle(b / a)\rangle_{\mathrm{ETG}}=0.69 \pm 0.19,\langle(b / a)\rangle_{\mathrm{dE}}=0.70 \pm 0.18$ and $\langle(b / a)\rangle_{\mathrm{dE}(b c)}=$ $0.63 \pm 0.2$.

### 3.4.2 Application of the GOLDMine distance

For all the above showed plots a constant distance of 16.5 Mpc was applied. Since the GOLDMine data base also provides the distances according to the different clouds of the Virgo cluster (see Figs. 3.6 and 3.7), it is worth to compare our results with the different distance estimates. It has to be pointed out that these cloud distances based on Tully-Fisher distance measurements (Tully and Fisher, 1977) of a small sample of spiral galaxies (Gavazzi et al., 1999), therefore, the estimated distance of each galaxy within a cloud based on a small number statistic. Furthermore, galaxies of a certain right-ascension and declination may be associated with e.g. the B-cluster with a distance of 23 Mpc , but in reality the galaxy is located in the periphery of the A-cluster with a distance of 17 Mpc . Therefore, also the different distances of the galaxies in the cloud should be handle with care.
It was realised through the analysis that some galaxies have distances which are well behind the Virgo cluster. To minimise the distortion of the averaged values due to background galaxies, galaxies with a distance of $\mathrm{D}>32 \mathrm{Mpc}$ are excluded from the calculation of the averages. Figures 3.28 to 3.31 show the same diagrams as Figs. 3.9, 3.18, 3.12 and 3.19, but now with the corresponding GOLDMine distances ${ }^{10}$. A summary of the shifts due of the distance depended parameters can be found in Fig. 3.32, which shows the averaged values of $\mathrm{R}_{\text {eff }}$ and $\mathrm{M}_{\mathrm{r}}$ derived with the GOLDMine distances. Since only the effective radius in kpc and the apparent magnitude are a function of the distance, the other parameters do not change and the averaged values can be taken from Figs. 3.25 and 3.26. The red bars in the histograms correspond to the morphological type, which have been reclassified by T. Lisker. The black bars show the same morphological types, but with the application of the GOLDMine distances. The green bars display the averages of the morphological types, which are reclassified by this study and additionally the GOLDMine distance was used (see next Section 3.4.3). All spiral galaxies and the dIs become brighter and larger in $\mathrm{R}_{\text {eff }}$, if one uses the GOLDMine distances. The difference between the magnitudes is about $\Delta \mathrm{M}_{\mathrm{r}} \approx 0.5 \mathrm{mag}$ and $\Delta \mathrm{R}_{\text {eff }} \approx 0.4 \mathrm{kpc}$ for the effective radii. This magnitude difference gets smaller within the spiral types from Sa to Sm . In case of the Sm galaxies, the difference is only $\Delta \mathrm{M}_{\mathrm{r}} \approx 0.15 \mathrm{mag}$.

[^18]

Figure 3.28: Colour-magnitude diagrams of late- and early-type galaxies in the Virgo cluster. Black dots show the ETG sample of Janz and Lisker (2008); Janz and Lisker (2009), where the vertical bars correspond to the $2 \sigma$ deviation within a magnitude bin of 0.5 mag. For clearness the CMD is dived into two parts. Left panel: CMD with spirals (Sa-Sd); Right panel: same CMD for Sm, SB, Amorphous, S and unknown galaxies. To calculate the distance of the galaxies values from the GOLDmine data base are used.


Figure 3.29: Same as Fig. 3.28, but for the effective radius $\mathrm{R}_{\text {eff }}$ vs. the mean effective surface brightness $\langle\mu\rangle_{\text {eff,r }}$ in r-band.


Figure 3.30: Same as Fig. 3.28, but for the effective radius $\mathrm{R}_{\text {eff }}$ vs. absolute magnitude $\mathrm{M}_{\mathrm{r}}$ in r-band.


Figure 3.31: Same as Fig. 3.28, but for the absolute magnitude $\mathrm{M}_{\mathrm{r}}$ vs.mean effective surface brightness $\langle\mu\rangle_{\mathrm{effr}, \mathrm{r}}$ in r-band.


Figure 3.32: Histograms of late-types with the GOLDMine distance. The red and black bars correspond to the initial classification by T. Lisker and the values if one uses the GOLDMine distances, respectively. The green bar shows the averages using the new classification of Section 3.4.3 and additionally the GOLDMine distance instead of a constant distance of $\mathrm{D}=16.52 \mathrm{Mpc}$.

### 3.4.3 Reclassification

Using the informations of all derived parameters, we are able to check the initial classification of the VCC and the revision by T. Lisker. The reclassified galaxies with the old types from the VCC and the corresponding new types will be summarised in Tab. 3.3. Galaxies classified as a mixed type of BCDs like "Im/BCD" were already discussed in Chapter 2.4.8, but are also listed.
In the class of Sa galaxies, there are three galaxies which are offset from the rest of this type, having a lower surface brightness, fainter magnitude and bluer (g-i)-colour. VCC-1933 has the lowest surface brightness, bluest (g-i)-colour and faintest magnitude, what would point to a misclassification in the VCC. Indeed, the GOLDMine data base classifies VCC-1933 as a Sm galaxy. However, with our findings in the parameter-space it is also possible to classify VCC-1933 as a dwarf irregular (dI) galaxy. The gri-image from SDSS in the right panel of Fig. 3.33 shows a galaxy with a possible spiral structure, but it could also interpreted as features of interaction. Therefore, the class of GOLDMine is used, but with a strong uncertainty.
The two other galaxies (VCC-0015 and VCC-1358) have properties of the morphological types of the Sc's and Sd's, therefore, they should be classified as Scd (with more weight to Sd). Moreover, the gri-image of VCC-0015 (left panel of Fig. 3.33) does not show a prominent bulge, what would be typical for Sa galaxies, reinforcing a misclassification. However, when applying the distances of the GOLDMine data base the magnitude of VCC-0015 would be shifted to the regime of the other Sa galaxies. Interestingly, the SDSS image of VCC-1358 (middle panel of Fig. 3.33) shows a galaxy without any spiral structure, but with a blue core in the centre of the galaxy. The optical image has similarities to $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$ and indeed does the location of VCC-1358 fit to the one of the $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$ in the entire parameter-space. Even with the distance of GOLDMine, there is no shift in the parameter-space, since it is stated with 17 Mpc . Therefore, VCC-1358 is reclassified as a dE(bc). The VCC classifies VCC-0989 as a "Sc or Im" and the GOLDMine data base as "Sc (dSc)". Using our recent results, we would classify VCC-0898 as a dI galaxy, since $\langle\mu\rangle_{\text {eff,r }}$ at a given $\mathrm{R}_{\text {eff,r }}$ is quite low compared to the other Sc galaxies. Without deep observation, the GOLDMine class of a dSc could not be ruled out, but the optical image in the left panel of Fig. 3.34 does not give any hint of a spiral structure, which would confirm the irregular class "dI". Three galaxies are conspicuous in the Sd class (VCC-0132, VCC-1605 and VCC-1685). All three galaxies are classified in the VCC and GOLDMine as Sd and SBd, respectively. Comparing the results of the galaxies with other types, VCC-0132 and VCC-1605


Figure 3.33: Galaxies initially classified as Sa , but not belonging to the parameterspace of Sa galaxies (from left to right: VCC-0015, VCC-1358 and VCC-1933).


Figure 3.34: Left: Galaxy VCC-0989 was initially classified as $\mathrm{Sc}(\mathrm{dSc}$ ), which could be reclassified as a dI; Middle: The very puzzling galaxy VCC-1664, which could not be associate to any late-type class; Right: VCC-1955, possibly a dE(bc).
share the same parameter-space like dI galaxies. VCC-1685 also has properties in common with the dIs, but also with the Sm galaxies. The optical image of VCC-0132 (left in Fig. 3.35) shows a faint galaxy with no obvious spiral structure, except for a faint loop in the south, supporting the galaxy class of "dI:", with a minor uncertainty. Furthermore, in the colour-colour-diagram in Fig. 3.20, the (i-z)-colours of VCC-0132 are only reached by far by the dIs.
The images of VCC-1605 (middle) and VCC-1685 (right panel of Fig. 3.35) reveal inclined galaxies. Therefore, the classification of the Sd type cannot be ruled out. There are several galaxies classified as " $S$ " or " $S(\mathrm{dS}$ )" in the sample. VCC-0213 and VCC-1955 were already discussed in Chapter 2, concluding that VCC-0213 could be a BCD with a spiral structure regarding its parameter-space, but nevertheless there are uncertainties regarding the classification.


Figure 3.35: Galaxies initially classified as Sd , but not belonging to the parameterspace of Sd galaxies (from left to right: VCC-0132, VCC-1605 and VCC-1685).

VCC-1574 is classified as a " S " in the VCC and as "?" (unknown) in GOLDMine. Cross-checking the redshift with SDSS reveals $\mathrm{z}=0.0756$, bringing it well behind the Virgo cluster. However, in the GOLDMine data base a velocity of $v_{\text {hel }}=639 \mathrm{~km} / \mathrm{s}$ is given. This value based on HI measurements from Gavazzi et al. (2006) (their Table 3). The optical image shows a galaxy with obvious spiral structure, but with the applied distance it is rather small for a typical spiral galaxy. Therefore, we are not able to give any new information regarding the classification, but VCC-1574 is marked with a ":", indicating the uncertainties in the classification.
As VCC-1358, also VCC-1955 covers almost the same parameter-space as the dE(bc)s and it was also included in the additional sample of Lisker et al. (2006a). The SDSS image of VCC-1955 is shown in the right panel of Fig. 3.34. The centre of the galaxy has blue colours and the surrounding disc does not show any feature of spiral structure. Therefore, VCC-1955 was reclassified to a dE(bc).
The parameter-space of VCC-0320, which is classified as a " $\mathrm{S}(\mathrm{dS})$ " in GOLDMine, shows all structural properties of the dI sample at a distance of $\mathrm{D}=16.5 \mathrm{Mpc}$. Regarding the GOLDMine data base, VCC-0320 has a distance of $\mathrm{D}=32 \mathrm{Mpc}$, making an allocation to the dIs still possible, but also shifting it into the region of the late spirals in the $\mathrm{M}_{\mathrm{r}}-\mathrm{R}_{\text {effi-plane. Despite to the other late spirals, VCC-0320 has a remarkable }}$ low $\langle\mu\rangle_{\text {eff,r },}$, which would confirm the classification as a dI.
Another galaxy classified as " $\mathrm{S}(\mathrm{dS}$ )" is VCC-1086. In the optical image (left panel of Fig. 3.36), one can see that it is strongly inclined and with the informations from the structural parameters we are able to conclude that VCC-1086 is not a dwarf spiral, since it has a too bright magnitude and a too high surface brightness compared to the other star forming dwarfs of the sample. However, due to the inclination it is not
possible to determine the exact type.
Within the class of the Sm galaxies, there is one galaxy that is offset from the other Sm galaxies. This offset for VCC-1554 is visible in the entire parameter-space. The visual inspection of the gri-image in the right panel of Fig. 3.34, reveals a galaxy with extended star forming regions, a bar-like structure and an irregular morphology. Nevertheless, it has a very high surface brightness and a magnitude assimilable to early spirals. Thus, in the parameter-space VCC-1554 is not a Sm galaxy, but it is also not an early spiral regarding its optical morphology. Due to the above mentioned problems a reclassification is not easily possible and we just mark the galaxy with a "pec" to indicate the peculiarity of this galaxy. In the morphological class of the


Figure 3.36: The "S (dS)" galaxy VCC-1086.
"unknown", eight galaxies are in the region of the dIs, thus they were reclassified as "dI" by T. Lisker. There is one galaxy among these eight "unknowns", which optical image (middle panel of Fig. 3.37) shows a galaxy with an obvious starburst region. The comparison with the BCDs reveals that the total parameters of VCC-0237 do only fit to the parameters of the LSB-components, therefore, an allocation to the BCDs would not be sensible and the class of "Im/BCD:" was used.
In the parameter-space, VCC-0031 has in common all properties with the BCDs. However, the optical image shows a red core in the centre of the galaxy. To check the parameters of VCC-0031, it was analysed as a BCD. A fit to the outer exponential component of the surface brightness profile would require an inner flatting of the surface brightness profile and the ( $\mathrm{g}-\mathrm{i}$ )-colour profile also does not have the typical slope as the other BCDs. Therefore, VCC-0031 is still classified as "unknown".
There is one formally classified unknown, which shares the same properties as the BCDs (right panel of Fig. 3.37). Therefore, VCC-1411 is reclassified as a BCD and was analysed and included in the BCD sample of Chapter 2.

Four other unknown galaxies are classified as transition types between dI and BCDs, indicated by "dI/BCD" and were already discussed in Chapter 2.
The only unknown galaxy left over after the revision of T. Lisker, corresponds to VCC0113. Using the informations from the (u-g)-(i-z)-diagram, it became clear that only the dIs share the same location in Fig. 3.21. In the parameter-space it also resides in the same region as the dIs. Thus, VCC-0113 can be reclassified as a dI, however, with some uncertainties due to the visible loops in the optical image of Fig. 3.37 (left). Figures 3.32 and 3.26 show the averaged results of each morphological class. The first


Figure 3.37: Galaxies initially classified as unknown (from left to right: VCC-0113, VCC-0237 and VCC-1411).
bar (red) indicates the average from the VCC and the second bar (black) the averaged values using the classes after the reclassification and application of the GOLDMine distances. As seen in the figures, the influence to the parameters is minor but visible.

Table 3.3: The reclassification of the VCC galaxies.

| VCC | VCC - type | GOLDMine | New class |
| :---: | :---: | :---: | :---: |
| 0015 | Sa | Sa | $\mathrm{Scd} ?$ |
| 0020 | unknown | unknown | dI |
| 0030 | unknown | unknown | dI |
| 0085 | unknown | unknown | dI |
| 0113 | unknown | unknown | $\mathrm{dI}:$ |
| 0132 | Sd | Sd | $\mathrm{dI}:$ |
| 0135 | Spec/BCD | $\mathrm{S} / \mathrm{BCD}$ | $\mathrm{Spec} / \mathrm{BCD} ?$ |
| 0213 | $\mathrm{Sd} ? / \mathrm{BCD} ?$ | $\mathrm{~S} / \mathrm{BCD}$ | $\mathrm{S} / \mathrm{BCD}$ |
| 0237 | unknown | unknown | $\mathrm{dI} / \mathrm{BCD}:$ |
|  |  | continued |  |


| VCC | VCC - type | GOLDMine | New class |
| :---: | :---: | :---: | :---: |
| 0247 | unknown | unknown | dI |
| 0309 | Im/BCD | Im/BCD | dI |
| 0320 | S | $\mathrm{Sc}(\mathrm{dS})$ | dI: |
| 0379 | unknown | unknown | dI |
| 0429 | unknown | unknown | dI/BCD |
| 0446 | Im/BCD | Im/BCD | Im/BCD |
| 0488 | unknown | unknown | dI |
| 0596 | Sc ? | $\mathrm{Sc}(\mathrm{dS})$ | Sc |
| 0655 | Spec,N:/BCD | S/BCD | BCD: (Mrk-86 like) |
| 0737 | Sd/BCD? | S/BCD | - |
| 0848 | ImIIIpec/BCD | Im/BCD | - |
| 0989 | Sc | Sc (dS) | dI |
| 1086 | S | S (dS) | S |
| 1179 | ImIII/BCD | Im/BCD | dI/BCD |
| 1237 | unknown | unknown | dI/BCD |
| 1358 | Sa | Sa | $\mathrm{dE}(\mathrm{bc})$ |
| 1374 | ImIII/BCD | Im/BCD | BCD |
| 1411 | unknown | Pec | BCD |
| 1427 | Im/BCD: | Im/BCD | dI/BCD |
| 1515 | unknown | unknown | dI |
| 1554 | Sm | Sm | Sm(pec) |
| 1574 | S | ? | S: |
| 1605 | Sd | Sd | Sd |
| 1623 | unknown | unknown | dI/BCD |
| 1685 | Sd | Sd | Sd |
| 1713 | unknown | unknown | dI/BCD |
| 1804 | ImIII/BCD | $\mathrm{Im} / \mathrm{BCD}$ | dI/BCD |
| 1933 | Sa | Sm | Sm ? |
| 1955 | Spec/BCD | S/BCD | $\mathrm{dE}(\mathrm{bc})$ |
| 1960 | ImIII/BCD? | $\mathrm{Im} / \mathrm{BCD}$ | dI/BCD |
| 2007 | ImIII/BCD: | Im/BCD | dI/BCD |
| 2037 | ImIII/BCD | Im/BCD | dI |

The reclassification of the VCC galaxies.

Table 3.4: Comparison between certain and possible: $\left\langle\mathrm{M}_{\mathrm{r}}\right\rangle$.

| Type | Certain | $\sigma$ | Possible | $\sigma$ |
| :--- | :---: | :---: | :---: | :---: |
| BCDs | -16.29 | 1.25 | -16.39 | 1.11 |
| dI | -15.67 | 0.92 | -15.45 | 1.14 |
| Sa | -20.21 | 0.93 | -19.43 | 1.70 |
| Sb | -19.99 | 1.14 | -19.79 | 1.10 |
| Sc | -18.91 | 1.24 | -18.74 | 1.46 |
| Sd | -17.43 | 1.27 | -16.95 | 0.65 |
| Sm | -16.88 | 1.18 | -15.99 | 1.52 |

### 3.4.4 Difference between certain and possible cluster members

As mention in Section 3.2.1, the membership of the Virgo galaxies is divided into certain and possible members. However, this membership was not taken into account in diagrams of the last sections. The following Tabs. 3.4 to 3.7 show the differences between the parameters of the certain and possible members. Errors are given by an $1 \sigma$ deviation from the average. Morphological types without any possible member (e.g. Amorphous) are not listed.
To calculate the averages of the parameters, the morphological classes after the reclassification are used, as well as the distance informations from GOLDMine.
The BCDs show no significant deviation between the certain and possible members, except for the (g-i)-colours, which are slightly bluer for the certain members.
The spiral galaxies from Sa to Sm show no major differences between certain and possible members. If there are any differences in the parameters, then they are within the $1 \sigma$ deviation of the averages.
These minor differences of the certain and possible members in Tabs. 3.4 to 3.7 retroactively justify the treatment as one group without any differentiation in the diagrams.

### 3.4.5 Comparison with semi analytical models

Semi analytical models (SAMs) are used to simulate the evolution of galaxies, using analytical recipes to dark matter merger trees. The results of our analysis of the late-type galaxies was included in a study of Weinmann et al. (2011) using SAMs from the study of Guo et al. (2011). The SAMs of Guo et al. (2011) based on the models of De Lucia and Blaizot (2007) with some modifications to account for the

Table 3.5: Comparison between certain and possible: $\left\langle\mathrm{R}_{\text {eff }}\right\rangle[\mathrm{kpc}]$.

| Type | Certain | $\sigma$ | Possible | $\sigma$ |
| :--- | :---: | :---: | :---: | :---: |
| BCDs | 0.64 | 0.30 | 0.73 | 0.31 |
| dI | 1.48 | 0.72 | 1.68 | 1.55 |
| Sa | 2.57 | 1.22 | 2.15 | 0.88 |
| Sb | 2.80 | 1.33 | 2.09 | 1.18 |
| Sc | 2.38 | 1.40 | 2.09 | 1.12 |
| Sd | 2.24 | 1.37 | 1.53 | 0.34 |
| Sm | 1.41 | 0.65 | 0.89 | 0.43 |

Table 3.6: Comparison between certain and possible: $\langle(\mathrm{g}-\mathrm{i})\rangle$.

| Type | Certain | $\sigma$ | Possible | $\sigma$ |
| :--- | :---: | :---: | :---: | :---: |
| BCDs | 0.47 | 0.32 | 0.60 | 0.16 |
| dI | 0.63 | 0.17 | 0.66 | 0.22 |
| Sa | 1.06 | 0.10 | 1.07 | 0.19 |
| Sb | 0.95 | 0.28 | 1.05 | 0.21 |
| Sc | 0.75 | 0.20 | 0.78 | 0.18 |
| Sd | 0.58 | 0.19 | 0.59 | 0.15 |
| Sm | 0.56 | 0.14 | 0.56 | 0.24 |

Table 3.7: Comparison between certain and possible: $\langle\langle\mu\rangle\rangle$.

| Type | Certain | $\sigma$ | Possible | $\sigma$ |
| :--- | :---: | :---: | :---: | :---: |
| BCDs | 21.08 | 1.04 | 21.35 | 0.87 |
| dI | 23.52 | 0.89 | 23.67 | 0.86 |
| Sa | 20.18 | 0.66 | 20.62 | 1.14 |
| Sb | 20.59 | 0.81 | 20.11 | 1.01 |
| Sc | 21.25 | 0.77 | 21.14 | 0.90 |
| Sd | 22.55 | 0.67 | 22.49 | 0.34 |
| Sm | 22.19 | 1.01 | 22.16 | 0.77 |

Table 3.8: Comparison between certain and possible: $\langle\mathrm{b} / \mathrm{a}\rangle$.

| Type | Certain | $\sigma$ | Possible | $\sigma$ |
| :--- | :---: | :---: | :---: | :---: |
| BCDs | 0.64 | 0.20 | 0.61 | 0.18 |
| dI | 0.66 | 0.17 | 0.72 | 0.19 |
| Sa | 0.56 | 0.19 | 0.54 | 0.23 |
| Sb | 0.56 | 0.27 | 0.50 | 0.15 |
| Sc | 0.52 | 0.25 | 0.58 | 0.24 |
| Sd | 0.54 | 0.29 | 0.61 | 0.29 |
| Sm | 0.54 | 0.14 | 0.58 | 0.16 |

increased influence of the effects like SNe- and AGN-feedback or environmental effects due to the shallow gravitational potential of the dwarf galaxies. In addition, the SAMs of Guo et al. (2011) use the high resolution Millennium-II simulation (Boylan-Kolchin et al., 2009), which enables the investigation of dwarf galaxies down to a stellar mass of $\sim 10^{7.5} \mathrm{M}_{\odot}$. Furthermore, the disc sizes of the SAMs galaxies are calibrated to fit the observation of 140000 galaxies from the SDSS by Shen et al. (2003), what makes a comparison with our sample possible.

In Fig. 3.38 the black dots corresponds to the results of the SAMs and the coloured points to the results of this study and additionally the results of Janz and Lisker (2008); Janz and Lisker (2009). The overlap of our sample with the SAM predictions fit quite well.

### 3.4.6 Comparison with the literature

A compilation of data from the literature is given by the study of Graham and Worley (2008) (hereafter G08). The medians of the structural parameters of G08 were given for the different morphological types of spirals including irregular galaxies (Irr in G08). In this study the scale length $h$ of the different types of spirals are given and can be transformed (Mo et al., 2010) into $\mathrm{R}_{\text {eff }}$ via:

$$
\begin{equation*}
\mathrm{R}_{\mathrm{eff}}=1.67 \cdot \mathrm{~h} . \tag{3.15}
\end{equation*}
$$

The scale lengths of G08 are given in the B and I band, therefore the effective radii in $g$ and $i$ of our sample are used for comparison. In Tab. 3.9 the results of $\mathrm{R}_{\text {eff }}$ are summarised. Interestingly, there is a huge difference between our results and the ones of G08, which accounts for more than a factor of two. It has to be pointed out that


Figure 3.38: Comparison of our results of the late-type galaxies within the Virgo cluster with the SAMs of Guo et al. (2011) (Courtesy: T. Lisker).

Equation 3.15 is valid for pure disc galaxies, therefore, in presence of a dominant bulge like in Sa galaxies, the effective radii are slightly underestimated. However, also in systems with a low bulge-to-disc ratio (e.g. Sc or Sd ), the differences between our results and the ones of G08 are impressive. This finding is quite puzzling, since our results fit well to the predictions of the SAMs of Section 3.4.5 and therefore also with the calibration sample from the SDSS of Shen et al. (2003).

Table 3.9: Comparison of $\mathrm{R}_{\text {eff }}$ between Graham and Worley (2008) and our study.

| Type | $\mathrm{R}_{\text {eff,B }}[\mathrm{kpc}]$ <br> (Graham) | $\mathrm{R}_{\text {eff,g }}[\mathrm{kpc}]$ <br> (our study) | $\mathrm{R}_{\text {eff, },}[\mathrm{kpc}]$ <br> $($ Graham $)$ | $\mathrm{R}_{\text {eff, } \mathrm{i}}[\mathrm{kpc}]$ <br> (our study) |
| :---: | :---: | :---: | :---: | :---: |
| Sa | 5.01 | 2.53 | 5.73 | 2.43 |
| Sb | 6.15 | 2.66 | 6.43 | 2.49 |
| Sc | 5.63 | 2.26 | 5.68 | 2.21 |
| Sd | 5.54 | 1.99 | 6.11 | 1.99 |
| Sm | 6.21 | 1.25 | 7.06 | 1.32 |
| Irr | 3.44 | 2.22 | 3.32 | 2.31 |

### 3.5 Summary and outlook

In this chapter we analysed the structural properties of the late-type galaxies within the Virgo cluster. Within the parameter-space the results of the photometry were compared within the different morphological classes and outliers were investigated in more detail and were cross-checked with the GOLDMine data base. If possible, these outliers were reclassified and an updated version of the VCC was created. The results of the reclassification are particular interesting for future investigations, since new found $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$ can be the target for further projects to study their possible connection to other dwarf galaxies, like BCDs.
We also found that there is no clear distinction between dwarf galaxies (dEs, dIs, BCDs) and low-luminosity spiral galaxies. This is particular true for the late spirals of the morphological type Sc to Sm.
After the inspection of the structural parameters of the late-type galaxies within the Virgo cluster, the question arises how these parameters will evolve with time and how the descendants of today's late-type galaxies will look like. To answer these questions, we use the evolutionary synthesis models GALEV (Kotulla et al., 2009). Therefore, the results of this section are the foundation of the next section, which analyses the evolution with time of the late-types and investigates their possible connections to the early-type Virgo population.

## Chapter 4 The evolution of late-type galaxies


#### Abstract

In the prior sections the structural parameters of the late-type galaxies of the Virgo cluster were derived. In this section the evolution of these galaxies is investigated and finally compared to the early-type galaxy population. For this purpose the evolutionary synthesis models of GALEV are used, which are also able to account for the influence of the dense cluster environment. The results of this section will shed light on the possible evolutionary connections between late-type and early-type galaxies.


This study will be timely published together with Thorsten Lisker and Ralf Kotulla.

### 4.1 Introduction: Transformation scenarios for turning late-type into early-types

Several studies on the Virgo cluster revealed a number of substructure (e.g. Binggeli et al., 1985 , 1993) and a large variety in morphology, with a dependence of the fraction of blue-to-red galaxies on the distance to the cluster centre, which is expressed by the morphology-density relation (Dressler, 1980). This substructure, as well as the variety in morphology, could also be confirmed by the results of the previous Chapter 3. As expected, our results have shown that late-type galaxies differ from early-type galaxies, but they have similarities too. Especially in the case of the LSB-components of BCDs, there are striking similarities in the structural and photometric properties, making an evolutionary connection between these two morphological types possible. Naturally the question arises, how the other late-types (e.g. spiral galaxies) will evolve with time and how they will look like in several billion years? Thus the question is, whether the descendants of today's late-type galaxies have similar properties as the early-type galaxies, which are observed in the Virgo cluster today?
When studying the evolution of galaxies it is crucial to take into account effects that
influence the behaviour of galaxies. Since the galaxies of our study are located in the Virgo cluster, environmental effects like ram pressure stripping (RPS), starvation and harassment play an important role in the evolution of late-type galaxies.
When moving through the cluster, the inter stellar medium (ISM) of the galaxies interacts with the hot intra cluster medium (ICM). This interaction removes parts or the entire gas of the galaxies, which consequently will influence the star formation (SF) of these galaxies. In case of RPS (Gunn and Gott, 1972) the entire gas reservoir is removed from the galaxies on a very short time scale of a few 100 Myr (Roediger, 2009, and references therein). As a consequence the star formation rate (SFR) will decrease very rapidly on a comparable time scale. As an example, Boselli et al. (2008) assumed a time scale for the decrease of the SFR of $\sim 150$ Myr. Simulations of RPS have been applied to disc galaxies (Abadi et al., 1999) as well as dwarfs (Mayer et al., 2006).

In the starvation scenario only the surrounding gas envelope of a galaxy is removed, but not the gas in the disc. Hence, the SF will not stop until the remaining gas is exhausted (Larson et al., 1980).
The above described processes of RPS and starvation act on the gas reservoirs of the galaxies due to the dense cluster environment. But also the interaction among the galaxies themselves has an influence on the evolution of galaxies. Due to high-speed encounters between galaxies within the cluster, the gas and the stellar population are disturbed. In the case of disc galaxies, the stellar disc may become thickened by the interaction with other galaxies. The scenario is called galaxy harassment (Farouki and Shapiro, 1981; Moore et al., 1996; Mastropietro et al., 2005) and was also applied to simulations of low mass dwarfs (e.g. González-García et al., 2005; Smith et al., 2010).
Apart from the external processes, there are also processes within the galaxies itself, which drive the evolution too. First of all, every galaxy will simply evolve due to the evolution of the stellar population, since the lifetime of stars is set by their initial mass and metallicity (Sparke and Gallagher, 2000). If a galaxy forms plenty of very high mass stars ${ }^{1}$ within a short time span - a so called starburst - then these stars will end up their lives by supernovae ( SNe ) explosions at almost the same time. Due to the simultaneous explosions of SNe in a small region within the galaxy, a galactic wind may be formed, which is able to "blow out" gaseous material from the galaxies

[^19](Dekel and Silk, 1986; Izotov et al., 1996; Heckman et al., 2001). The loss of gaseous material not only affects the chemical composition of the galaxy, but also decreases the gas reservoir, which is needed to fuel the SF.
In summary, all the above described effects remove gas from galaxies, resulting in a decrease of the SF. Over the past decades it became clear that one of the key parameter in the evolution of galaxies is the star formation and therefore the star formation history (SFH). The SFH characterises the evolution of the SFR, which in turn describes the amount of stellar mass formed within a year. The SFR is expressed in units of $\mathrm{M}_{\odot} / \mathrm{yr}$ and strongly depends on the morphological galaxy type. This can be explained by the Kennicutt-Schmidt law, which describes a power law dependence of the gas surface density and the SFR of the galaxy (Schmidt, 1959; Kennicutt, 1998).
To model the evolution of galaxies and therefore the SFH, several programs and codes were developed over the last decades. One of these codes is GALEV (GALaxy EVolution; Kotulla et al., 2009, 2010) and forms the foundation of our here presented study. GALEV models were successfully applied to the entire range of SFHs that characterise the different morphological types of galaxies (e.g. Krüger et al., 1991, 1995; Falkenberg et al., 2009) and are therefore ideally suited to our study.

### 4.2 Stellar population synthesis models

There are two approaches to model the spectrum of a stellar population and therefore also of a galaxy. The first approach uses a linear combination of stellar spectra to fit the observation and is called stellar population synthesis. The interpretation of the results is limited to the actual observation and no predications of the evolution are made.
The second approach - the evolutionary synthesis models - tries to fit the observation by the means of integrated spectra, colours or the spectral energy distribution (SED) of the galaxy and were first investigated by Tinsley (1968). Current examples for these evolutionary synthesis models are the codes by Leitherer et al. (1999) (Starburst99), Bruzual and Charlot (2003) (BC03), Fioc and Rocca-Volmerange (1997) (PEGASE), Maraston (2005) and GALEV, which is used in the here presented study. In contrast to the stellar population synthesis, all evolutionary synthesis models have to assume an initial stellar mass function (IMF) and SFH, and are able to model the evolution of the galaxy via stellar isochrones. The IMF describes the number of stars that are formed in a specific mass interval. Studies showed that stars are not equally formed per mass bin, but instead less massive stars are formed more often than massive stars. There are
different IMF laws in the literature, which differ in the slopes of the IMF and therefore in the number abundance of star in certain mass bins (see Section 4.2.3).
Thus, with the assumption (or knowledge) of e.g. the IMF and SFH of a stellar population, we are able to investigate its evolution over an arbitrary period and are able to make predications of the future of the stellar population.

### 4.2.1 Why GALEV?

GALEV (GALaxy EVolution) is an evolutionary synthesis code, which models the evolution of galaxies over time. It was developed by Prof. Dr. U. Fritze and her group at the University of Göttingen (Germany) and is maintained by the GALEV-team. To run GALEV models one can use a web interface at www.galev.org, which provides standard configurations, for instance star formation histories or Galactic extinction. To apply more sophisticated configurations like modified SFHs (see Section 4.2.2) one can also use a local installation of GALEV.
To calculate the time evolution of galaxies of different types, GALEV models the spectral evolution of a stellar population and the ISM at the same time. The foundation of the evolution of a stellar population in GALEV are the BaSeL libraries of model atmospheres (Lejeune et al., 1997, 1998) and the stellar evolutionary tracks or isochrones from the Padova group (Bertelli et al., 1994).
The advantage of GALEV over other evolutionary synthesis models (Starburst99, BC03 and PEGASE) is the inclusion of the chemical evolution of the gas additionally to the stellar population. The yields of the stellar population are depending on the metallicity, which is taken into account by GALEV models and is called the "chemical consistent treatment". The yields of the stellar population are taken from Woosley and Weaver (1995) and van den Hoek and Groenewegen (1997) for the different stellar mass ranges.
Also included are the contributions of the nebular emission lines and the continuum of the young stellar population, which could contribute up to 50-60 \% to the total flux (Anders and Fritze-v. Alvensleben, 2003) and is therefore from particular interest for our study.
GALEV does not provide the magnitude in different filters of a galaxy. It just provides the SFHs and the spectra of galaxies of different types at a given time. To obtain the magnitudes as a function of time for a given SFH, the subroutine "COCOS" is used. With the subroutine COCOS, also input parameters like the extinction laws, attenuation and filter systems enter into the analysis.

To obtain the associated magnitudes in each filter of the model, the spectra are folded with the requested filter curve. In the concrete case of this work, the filter curves of SDSS are used, but others, like IR or UV filters, can easily be attached. With the subroutine COCOS one finally obtains the magnitudes for each time step in each filter. To emphasise again, no observational data is used at this point and only theoretically based models are used.
At the moment of writing, galaxies are treated by GALEV as one solid unit without any spatial resolution and internal dynamics.

### 4.2.2 Input parameters

To model the evolution of late-type galaxies one has to choose a reasonable SFH. There are several SFH laws in the literature, like the exponential declining SFR which can be used for elliptical galaxies, or the below described exponentialdelayed. The exponential-delayed SFHs are parametrised by a formula given by Gavazzi et al. (2002), which was first introduced by Sandage (1986) using the study of Gallagher et al. (1984):

$$
\begin{equation*}
\operatorname{SFR}(\mathrm{t})=\frac{\mathrm{t}}{\tau^{2}} \cdot \exp \left(-\frac{\mathrm{t}^{2}}{2 \tau^{2}}\right) . \tag{4.1}
\end{equation*}
$$

All of these exponential-delayed SFHs show an increase ${ }^{2}$ of the SFR up to the point $\tau$ and then an exponential decrease with time (see Fig. 4.1). Different values of $\tau$ correspond to various SFHs of the different spectral galaxy types. For instance, an elliptical galaxy will form almost all its stellar mass within a short period of some Gyrs (e.g. $\tau \lesssim 1 \mathrm{Gyr}$ ), whereas spiral galaxies continuously forming stars until now, but with different strength of SF ( $\tau>1 \mathrm{Gyr}$ ). Therefore, the value $\tau$ can be used as indicator for the spectral type of the galaxy. With some limitations, these spectral types can be associated with the morphological type (MacArthur et al., 2004).
Figure 4.1 shows the SFH for different values $\tau$ for a galaxy with an initial gas mass of $\mathrm{M}=1 \cdot 10^{10} \mathrm{M}_{\odot}$.

### 4.2.3 Starbursts and truncations

The SFH of a galaxy is not always a smooth curve as given by Equation 4.1, but could show short periods of intense SF (starbursts). GALEV models are able to account for these starbursts, which is particular important for BCDs (Krüger et al., 1991;

[^20]

Figure 4.1: SFHs for different sets of $\tau$.

Krüger and Fritze-v. Alvensleben, 1994). The starburst was modelled by a gaussian increase and an exponential decrease of the SFR. The strength of the burst is given by the burst parameter b , which is defined as

$$
\begin{equation*}
\mathrm{b}=\frac{\mathrm{M}_{*, \text { formed }}}{\mathrm{M}_{\mathrm{gas}}} \tag{4.2}
\end{equation*}
$$

where $\mathrm{M}_{*, \text { formed }}$ is the total mass of stars formed during the starburst and $\mathrm{M}_{\text {gas }}$ is the available gas mass at the beginning of the burst. For models with a bursts, the burst parameter was set to $b=0.1$. That means $10 \%$ of the available gas is transformed into stars. As shown by several studies, the starburst phenomenon cannot last for more than $10^{7}$ yrs (Thuan, 1991; Thornley et al., 2000), owing to the fact that the gas reservoir in BCDs is limited. Therefore, the starburst in the models is also limited to this time scale.
Since the galaxies are located in the dense environment of the Virgo cluster, there are several mechanisms - ram pressure stripping (Gunn and Gott, 1972), harassment (Moore et al., 1996) and tidal forces (Mayer et al., 2001) - acting on the cluster galaxies and which depend on the cluster centric radii. As commonality all these mechanisms remove gas from the galaxies and therefore reduce or truncate the SFR. Regarding to the lifetime of a galaxy, the removal of gas in a cluster mostly acts on very short time scales. In the GALEV models the truncation of SFR is simulated by a stop of the SF at a given time (Bicker et al., 2002), which is illustrated as a red line in Fig. 4.2. In our
models we assume that the SF stops today at $\mathrm{t}=13.7$ Gyr. The decline of the SFR due to RPS acts on very short time scales, therefore a decline time scale of 150 Myr (Boselli et al., 2008) was chosen.
At this point, GALEV provides the following initial mass functions (IMFs): Salpeter (1955) and Kroupa (2001). In the nearby future also the IMF of Chabrier (2003) will be implemented, which accounts the different components of a galaxy (disc, bulge, globular clusters). The IMF of Salpeter uses a power law with exponent $\alpha=2.35$ for the entire mass range, while the Kroupa IMF uses a broken power-law for the different mass regimes with various slopes ( $\alpha_{1}=1.3$ for $0.1<\mathrm{M}<0.5 \mathrm{M}_{\odot}$ and $\alpha_{2}=2.3$ for $\mathrm{M}>0.5 \mathrm{M}_{\odot}$ ). These two IMF-laws mainly differ at the low mass slopes. For our galaxies the IMF of Salpeter (1955) was chosen.
To correct the flux for reddening and extinction of the interstellar dust, there are two extinction laws used by GALEV: Cardelli et al. (1989) and Calzetti et al. (2000). Since late-type galaxies with active star formation are in the scope of this study, the extinction law by Calzetti et al. (2000) was chosen for the models. The dust physics in GALEV is quite simple compared to other studies (e.g. Möllenhoff et al., 2006; Piovan et al., 2006a,b), which use the two dimensional structure of the galaxies. Since GALEV has no spatial resolution and the galaxies are therefore treated as one dimensional units, no assumption of the geometry of the galaxy is necessary, which in turn justifies the simple dust approach of a constant value per galaxy.
The chemical abundance of each galaxy can be fixed to a constant value or the evolution of the ISM and stellar population is modelled too. This second approach is called "the chemically consistent treatment" (Moeller et al., 1997).
After the determination of the appropriate SFH for each galaxy, all galaxies can evolve with and without any disturbance, respectively. As an illustration for the evolution of a galaxy with the different model parameters, Fig. 4.2 shows the SFH for a galaxy with $\tau=5 \mathrm{Gyr}$ and an initial gas mass of $\mathrm{M}=1 \cdot 10^{10} \mathrm{M}_{\odot}$ for an undisturbed (green curve), a truncated (red line) and a starbursting (blue line) evolution.
The spectral evolution of a galaxy is shown in Fig. 4.3. The figure displays the spectra of the galaxy at different time steps $\mathrm{t}_{1}=1 \mathrm{Gyr}$ and $\mathrm{t}_{2}=13.6 \mathrm{Gyr}$. The red spectrum displays the galaxy after 1 Gyr and the green spectrum after 13.6 Gyr. Obviously, the flux of the galaxy decreases with time and the strong emission lines disappear too. This is intelligible, since the flux in the UV and blue regime of the spectrum is dominated by young and massive stars, which have a very short lifetime compared to the intermediate stars (Green and Jones, 2004). Therefore, after the death of the massive stars, the


Figure 4.2: SFH for an example galaxy with $\tau=5 \mathrm{Gyr}$, an initial gas mass of $\mathrm{M}=$ $1 \cdot 10^{10} \mathrm{M}_{\odot}$ and three SF-modes: undisturbed (green line), truncation (red line) and burst (blue line).
light is dominated by the intermediate and low mass stellar population, which radiates mostly at optical and infrared wavelength.
Up to this point we just work with theoretical models, but the main question is how these models fit to the observations? To answer this question, the models have to be transformed into observables, which was achieved by folding the theoretical spectrum of the galaxy with the desired filter system. In this study the filter curves of the SDSS (but other filter systems are also possible, like the ones of UKIDSS or HST) are used to obtain the corresponding magnitude (here AB magnitudes, Bohlin and Gilliland, 2004) in a certain filter. The filter curves of SDSS and UKIDSS are exemplarly over plotted in Fig. 4.3. The height of the filter curves corresponds to the throughput of the SDSS telescope and illustrates the coverage of the used filter system.

### 4.3 Fitting the observed SEDs with GALEV models

### 4.3.1 GAZELLE

The most important step is the fitting of the observed data to the models. The finding of the best fitting model enables the investigation of the time evolution of the galaxies. This was done with the program "GAZELLE", which uses a slightly


Figure 4.3: Spectra of a model galaxy. Additionally plotted are the filter curves of the SDSS ( $\mathrm{u}, \mathrm{g}, \mathrm{r}, \mathrm{i}$ and z ) and UKIDSS ( $\mathrm{Y}, \mathrm{J}, \mathrm{H}$ and K ).
modified $\chi^{2}$ minimisation algorithm (Kotulla et al., 2009). As input parameters a set of theoretical SFHs from GALEV, magnitudes from the observations with corresponding photometric errors (see Equation 3.2), the preferred dust extinction law and a filter system are used. The $\chi^{2}$ algorithm uses the following formula for the minimisation ${ }^{3}$ :

$$
\begin{equation*}
\chi^{2}=\sum_{\mathrm{i}}^{\mathrm{n}}\left[\frac{\mathrm{~F}_{\mathrm{i}, \mathrm{mod}}-\alpha \cdot \mathrm{F}_{\mathrm{i}, \text { obs }}}{\sqrt{\sigma_{i, \text { mod }}^{2}+\sigma_{i, \text { obs }}^{2}}}\right]^{2} . \tag{4.3}
\end{equation*}
$$

In this equation $\mathrm{F}_{\mathrm{i}, \text { mod }}$ and $\mathrm{F}_{\mathrm{i}, \text { obs }}$ correspond to the fluxes in the different filters i to n of the model and the observations with the uncertainties given by $\sigma_{\mathrm{i}, \mathrm{mod}}^{2}$ and $\sigma_{\mathrm{i}, \mathrm{obs}}^{2}$. The uncertainties in the models are due to the limited knowledge of the input physics, like the IMF and incomplete stellar libraries. The drawbacks of evolutionary synthesis models due to incomplete knowledge of the input physics are reviewed in Conroy et al. (2009), Conroy et al. (2010) and Conroy and Gunn (2010). At the moment a model uncertainty of $\sigma_{\mathrm{i}, \text { mod }}^{2}=0.1$ is assumed, independent of the used filter.
For our analysis one of the most important factors is $\alpha$, which scales the model luminosity to fit the observation. By scaling the luminosity, it also scales the properties like the mass and SFR of the corresponding galaxy.

[^21]The different distributions of the $\chi^{2}$ values are displayed in Fig. 4.4 for a galaxy with a variable age (left) and fixed age of 13.7 Gyr (right). In our study, we assume a constant age of 13.7 Gyr to reduce the number of free parameters, what is also supported by observations of galaxies, where in almost all cases contributions of an old stellar population can be found ${ }^{4}$. Due to the constant age, all $\chi^{2}$ values are on a vertical line on the right hand side of Fig. 4.4. Colour coded are the different GALEV-models of the grid - and therefore also the different values of $\tau$ that are correlated to the spectral galaxy type - which can be found in Tab. 4.1. The model with the lowest $\chi^{2}$ value was chosen as the best fitting model. Looking at Fig. 4.4, one also realises that the best fitting values of $\chi^{2}$ are below unity and not at unity, as one may expects. This is due to the application of the uncertainties of the model $\sigma_{\mathrm{i}, \mathrm{mod}}^{2}$ and the photometry $\sigma_{\mathrm{i}, \mathrm{obs}}^{2}$ in Equation 4.4. The classical $\chi^{2}$ would only uses the uncertainties of the photometry $\sigma_{\mathrm{i}, \text { obs }}^{2}$, leading to $\chi^{2}$ values close to unity. However, GAZELLE additionally uses the uncertainties of the models $\sigma_{i, m o d}^{2}$, which in turn leads to the $\chi^{2}$ values below unity.


Figure 4.4: $\chi^{2}$ distribution of a galaxy with variable galaxy age (left) and a constant age of 13.7 Gyr (right). Colour coded are the model-numbers according to Tab. 4.1.

Figure 4.5 shows the spectral energy distribution (SED) for two galaxies. In both cases the galaxies were fitted with and without dust. The observed r-band magnitudes of the galaxies are the results of Chapter 3 and are displayed via black data points in Fig. 4.5. The horizontal bars indicate the FWHM of the corresponding SDSS filters. Blue circles in the SEDs show the best fits to the observations when dust extinction is not taken into account, while the red circles belong to the models with dust extinction. The difference between the SED with dust extinction and without are not large, but not negligible. The models that include the dust extinction fit the observation sightly

[^22]better. It is also possible that for the same galaxy, different values of $\tau$ are the best fit for the case of extinction and no extinction. It also has to be pointed out that the extinction is not simply set to zero (see Equation 4.4), but instead we re-run the entire fitting algorithm of GAZELLE to find the best fitting model. If one just set the used dust extinction for a model to zero, the differences of the observed and model SEDs became quite large. For instance, in case of an edge-on spiral, the difference is up to 1 magnitude within a filter.


Figure 4.5: SEDs of two galaxies of different types. Horizontal error bars correspond to the width of each SDSS filter, while the vertical ones display the photometric errors of the data. Left: VCC-0024 (BCD); Right: VCC-0025 (Sc).

### 4.3.2 The grid

To model the observed properties of the galaxies, a model grid was created, using the above mentioned input parameters. $\tau$ was chosen in the range of $0.5 \leq \tau \leq 25$ Gyr. The study of Sandage (1986) assumed for the Sm galaxies a SFH that is rising with time (their Fig. 10 b). Thus, in case of the "Sm" model an exponential increasing SFR was assumed with $\tau=-4$ Gyr in Equation 4.1 (see also MacArthur et al., 2004). This $\tau$ was chosen to reach the minimal ( $\mathrm{g}-\mathrm{i}$ )-colours of the Sm galaxies of our sample. Sm-models with $\tau<-4$ Gyr do not significantly differ in (g-i)-colours from the $\tau=-4$ Gyr model. Therefore, we only include the Sm-model with $\tau=-4 \mathrm{Gyr}$.
For all models the chemically consistent treatment of the ISM of the galaxy was applied. Table 4.1 summarised the used models with the corresponding input parameters. The last model in the table was included as a tracer for errors in the fitting of the models to the observations. If no model matches to the observation, the last model in the list is automatically chosen as the best-fit model. Therefore, we
include a model of a massive and metal rich elliptical galaxy, which should not describe the spectral properties of today's late-type galaxies. Thus, if a galaxy is fitted by the elliptical model it pointed to a failure in the modelling.
To model the analysed galaxies, the grid should cover the entire range of the photometric parameters. Figure 4.6 shows the evolution with time of the (g-i)colours of the different models for an undisturbed evolution without any interaction. The black vertical line at 13.7 Gyr corresponds to the (g-i)-colour of range $0.1<(\mathrm{g}-\mathrm{i})<1.2 \mathrm{mag}$ of the observed late-type galaxies, without including the BCDs with their partially extreme blue colours. Except for the very red (g-i)-colours of $(\mathrm{g}-\mathrm{i})>1.1 \mathrm{mag}$, the grid covers the entire colour-space of the late-type galaxies.
Figures 4.7 and 4.8 show the same models, but with an additional burst and truncation, respectively. The burst was not applied to every single model, but only to models with $\tau>5 \mathrm{Gyr}$, which is typical for bursting dwarfs. It is also possible that the BCDs are in different phases of the starburst. For example, one BCDs perhaps just ignites the starburst, while another is at the end of the starburst. Therefore, we applied different starting points for the starburst - from 13.68 Gyr ( 20 Myr ago) up to 13.60 Gyr ( 100 Myr ago). The bursts are necessary to reach the extreme blue colours of the BCDs with their strong starburst. In the extreme case of VCC-1313, the (g-i)-colour of the entire galaxy reaches a value of $(\mathrm{g}-\mathrm{i})=-0.3$ mag. Interestingly, the influence of the starburst on the (g-i)-colour is minor 1 Gyr after the onset of the burst and is almost invisible after the 2 Gyr.
The truncation was used to simulate the interaction between the galaxy and the ICM of the cluster via RPS, resulting in a gas loss and a decrease of the SF. The truncation starts at 13.7 Gyr with a time scale of $\tau_{\text {trunc }}=150 \mathrm{Myr}$ (Boselli et al., 2008). Figure 4.8 shows the time evolution of the (g-i)-colours of the models if one applies a truncation. Before the truncation sets in, the difference between the minimal and maximal (g-i)colour of the models is $\Delta(\mathrm{g}-\mathrm{i}) \approx 0.85$ mag. However, after 2 Gyr the difference of the models became minimal with $\Delta(\mathrm{g}-\mathrm{i}) \approx 0.3$ mag. Thus, the application of a truncation to the time evolution of the galaxies will bring all galaxies into a narrow (g-i)-colour region of $0.8<(\mathrm{g}-\mathrm{i})<1.1 \mathrm{mag}$.


Figure 4.6: (g-i)-colour evolution with time of the different models. The vertical line indicates the position of today $(\mathrm{t}=13.7 \mathrm{Gyr})$.

Table 4.1: Used model grid for GALEV.

| Model-id | Model name | $\tau$ <br> $[\mathrm{Gyr}]$ | Truncation <br> $[\mathrm{Gyr}]$ | Burst <br> $[\mathrm{Gyr}]$ |
| :--- | :---: | :---: | :---: | :---: |
| model_01 | gavazzi_tau50e08_CC | 0.5 | - | - |
| model_01 | gavazzi_tau10e09_CC | 1.0 | - | - |
| model_02 | gavazzi_tau15e09_CC | 1.5 | - | - |
| model_03 | gavazzi_tau20e09_CC | 2.0 | - | - |
| model_04 | gavazzi_tau30e09_CC | 3.0 | - | - |
| model_05 | gavazzi_tau40e09_CC | 4.0 | - | - |
| model_06 | gavazzi_tau50e09_CC | 5.0 | - | - |
| model_07 | gavazzi_tau60e09_CC | 6.0 | - | - |
| model_08 | gavazzi_tau80e09_CC | 8.0 | - | - |
| model_09 | gavazzi_tau10e10_CC | 10.0 | - | - |
| model_10 | gavazzi_tau15e10_CC | 15.0 | - | - |
| model_11 | gavazzi_tau25e10_CC | 25.0 | - | - |
| model_33 | gavazzi_tau50e08_CC_trunc137e09 | 0.5 | 13.7 | - |
| model_22 | gavazzi_tau10e09_CC_trunc137e09 | 1.0 | 13.7 | - |
| model_23 | gavazzi_tau15e09_CC_trunc137e09 | 1.5 | 13.7 | - |

continued

| Model-id | Model name | $\tau$ | Truncation | Burst |
| :--- | :---: | :---: | :---: | :---: |
|  |  | $[\mathrm{Gyr}]$ | $[\mathrm{Gyr}]$ | $[\mathrm{Gyr}]$ |
| model_24 | gavazzi_tau20e09_CC_trunc137e09 | 2.0 | 13.7 | - |
| model_25 | gavazzi_tau30e09_CC_trunc137e09 | 3.0 | 13.7 | - |
| model_26 | gavazzi_tau40e09_CC_trunc137e09 | 4.0 | 13.7 | - |
| model_27 | gavazzi_tau50e09_CC_trunc137e09 | 5.0 | 13.7 | - |
| model_28 | gavazzi_tau60e09_CC_trunc137e09 | 6.0 | 13.7 | - |
| model_29 | gavazzi_tau80e09_CC_trunc137e09 | 8.0 | 13.7 | - |
| model_30 | gavazzi_tau10e10_CC_trunc137e09 | 10.0 | 13.7 | - |
| model_31 | gavazzi_tau15e10_CC_trunc137e09 | 15.0 | 13.7 | - |
| model_32 | gavazzi_tau25e10_CC_trunc137e09 | 25.0 | 13.7 | - |
| model_34 | gavazzi_tau60e09_CC_burst20 | 6.0 | - | 13.68 |
| model_35 | gavazzi_tau60e09_CC_burst50 | 6.0 | - | 13.648 |
| model_36 | gavazzi_tau60e09_CC_burst100 | 6.0 | - | 13.60 |
| model_37 | gavazzi_tau80e09_CC_burst20 | 8.0 | - | 13.68 |
| model_38 | gavazzi_tau80e09_CC_burst50 | 8.0 | - | 13.648 |
| model_39 | gavazzi_tau80e09_CC_burst100 | 8.0 | - | 13.60 |
| model_40 | gavazzi_tau10e10_CC_burst20 | 10.0 | - | 13.68 |
| model_41 | gavazzi_tau10e10_CC_burst50 | 10.0 | - | 13.648 |
| model_42 | gavazzi_tau10e10_CC_burst100 | 10.0 | - | 13.60 |
| model_43 | Sm_40e09 | -4.0 | - | - |
| model_44 | elliptical_fake | - | - | - |

The model grid.


Figure 4.7: Same as Fig. 4.6, but additionally a starburst was assumed in a subset of the models.


Figure 4.8: Same as Fig. 4.6, but with a truncation at $\mathrm{t}=13.7$ Gyr.

### 4.4 Results

### 4.4.1 The output

After the determination of the best fitting SFH and the convolution with the SDSS observation, GALEV gives plenty of output parameters, which are summarised in Tab. 4.2. The offsets $\alpha$ of Equation 4.4, which are applied to the models to fit the observations, are particular important to model the time evolution of the galaxies. The evolution of the magnitude M of the corresponding galaxy can be calculated by:

$$
\begin{equation*}
\mathrm{M}(\mathrm{t})=\mathrm{M}_{\mathrm{mod}}(\mathrm{t})-\alpha+\left[\mathrm{E}(\mathrm{~B}-\mathrm{V})(\mathrm{t}) \cdot\left(\mathrm{SFR}_{\bmod }(\mathrm{t}) / \mathrm{SFR}_{\mathrm{obs}} \cdot 10^{(-0.4 \cdot \alpha)}\right) \cdot \mathrm{A}_{\mathrm{M}}(\mathrm{t})\right] \tag{4.4}
\end{equation*}
$$

with the parameters of Tab. 4.2.
Since GALEV calculates the evolution of the magnitude for all input filters, we are also able to investigate the evolution of the different colours, like the ( g - i )-colours of Chapter 3.

Table 4.2: Output parameters from the GALEV/GAZELLE models.

| Parameter | Description |
| :--- | :--- |
| mod $_{\text {fit }}$ | Best fitting model |
| age | Age of the galaxy (held constant to 13.7 Gyr) |
| $\alpha$ | Applied offset to the model to fit the observations |
| $\mathrm{M}_{\text {tot }}$ | Total mass (gas + stars) of the galaxy in units of $\mathrm{M}_{\odot}$ |
| $\mathrm{M}_{\text {stellar }}$ | Stellar mass of the galaxy in units of $\mathrm{M}_{\odot}$ |
| $\mathrm{M}_{\text {gas }}$ | Gas mass of the galaxy in units of $\mathrm{M}_{\odot}$ |
| $\mathrm{M}_{\text {non-vis }}$ | Non visible mass of the galaxy in units of $\mathrm{M}_{\odot}$ |
| $\mathrm{Z}_{\text {gas }}$ | Gas metallicity |
| $\mathrm{SFR}_{\text {obs }}$ | Best fitting SFR to the observation |
| $\mathrm{SFR}_{\text {mod }}$ | SFR of the model without offsets |
| $\mathrm{M}_{\text {mod }}$ | Magnitude of the model without offsets |
| $\mathrm{A}_{\mathrm{M}}$ | Dust correction |
| $\mathrm{E}(\mathrm{B}-\mathrm{V})$ | Applied extinction |



Figure 4.9: Distribution of the different values of $\tau$ of the models for the Virgo sample. The first bar ("fail") includes all galaxies, where no appropriate model was found by GAZELLE.

### 4.4.2 Number distribution of the galaxies among the model grid

Figure 4.9 shows the number distribution of the different values of $\tau$ used by the GALEV models. In most of the cases a model with $\tau=25 \mathrm{Gyr}$ was chosen (model-11 in Tab. 4.1). The right hand side of Fig. 4.10 shows the distribution of the different morphological galaxy types within the $\tau=25$ Gyr model. In most of the cases this model was used to fit galaxies of the type $\mathrm{Sc}, \mathrm{dI}$ and $\mathrm{dE} / \mathrm{dI}$, a class that is treated as a subclass of the dIs in our analysis and that corresponds to the "dE/Im" class of the VCC. Due to the relatively large value of $\tau(\tau=25 \mathrm{Gyr})$ it is reasonable that most of the dIs are fitted with this model.

Figure 4.11 shows the ( g - i )-colours as a function of $\tau$. As expected, there is a clear trend for $\tau \leq 15 \mathrm{Gyr}$ of being redder with decreasing $\tau$, which shows the connection between $\tau$ and the morphological/spectral type of the galaxies.


Figure 4.10: Distribution of morphological galaxy types with $\tau=2$ Gyr (left panel) and $\tau=25 \mathrm{Gyr}$ (right panel).

Figures 4.12 and 4.13 show the relative number of each morphological type from the VCC as a function of $\tau$. The majority of the late-type galaxies in the sample avoid the region of $6 \leq \tau \leq 15$ Gyr. In the regime of $16 \leq \tau \leq 24 \mathrm{Gyr}$ there is no models in our grid, which therefore explains the gap in this region in Figs. 4.12 and 4.13. In case of the Sa galaxies, almost $50 \%$ of them were fitted with $\tau=3.0 \mathrm{Gyr}$, but also about $5 \%$ were fitted with a quite large $\tau$ of $\tau=25$ Gyr. Over $70 \%$ of the Sd galaxies were fitted with $\tau=25 \mathrm{Gyr}$ and the rest with $\tau \leq 25 \mathrm{Gyr}$.

### 4.4.3 Physical properties of the morphological types

Figure 4.14 shows the average total masses of the different morphological types of the sample. The most massive galaxies in Fig. 4.14 are the early spirals of type Sa to Sc with $\mathrm{M}_{\mathrm{tot}} \approx 7 \cdot 10^{10} \mathrm{M}_{\odot}$. The late spirals ( Sd and Sm ), together with the loosely classified type of $S$ galaxies, have significantly lower averaged total masses in the range of $5.5 \cdot 10^{9}<\mathrm{M}_{\text {tot }}<9.5 \cdot 10^{9} \mathrm{M}_{\odot}$.
The unknown and dI galaxies have the lowest averaged total mass of the entire sample with $\mathrm{M}_{\text {tot }}<2 \cdot 10^{9} \mathrm{M}_{\odot}$.
The BCDs have an averaged total mass of $\mathrm{M}_{\mathrm{tot}}=5.5 \cdot 10^{9} \mathrm{M}_{\odot}$, similar to the Sm galaxies.
Other output parameters, like the gas metallicity, do not scale with the luminosity.


Figure 4.11: (g-i)-colours of the model galaxies vs. the best fitting value of $\tau$.


Figure 4.12: Relative number of the different morphological types from Sa to Sm vs. the best fitting value of $\tau$.


Figure 4.13: Same as Fig. 4.12, but for the Sm, dI, unknown and Amorphous galaxies.

Therefore, all galaxies with the same $\tau$ will also have the same metallicity. Figure 4.15 shows the evolution with time of the gas metallicity for a set of different values of $\tau$ for an undisturbed and truncated evolution, respectively. In the case of an undisturbed evolution, the metallicity of the models with $\tau \geq 5$ Gyr gradually increases, while for the truncation the models with $\tau>5 \mathrm{Gyr}$ and ages of $\mathrm{t}>13.7 \mathrm{Gyr}$ show an almost constant metallicity due to the cease of the star formation.
Figure 4.17 shows the mass-metallicity relation of the sample. The median masses are calculated within each metallicity bin, while metallicities with $Z>0.031$ are summarised into one bin. As one can see, the most massive galaxies also have the highest metallicity and vice versa in case of the low-mass galaxies, as expected from the observations and theory (e.g. Tremonti et al., 2004). Additionally shown in Fig. 4.18 is the stellar mass $\mathrm{M}_{\text {stellar }}$ (left) and the metallicity Z (right) as a function of $\tau$. For $\tau<5$ Gyr the stellar mass remains almost constant, while the metallicity increases with $\tau$. For $\tau>5$ Gyr both parameters decrease, which again reflects the behaviour of the mass-metallicity relation and shows the connection between $\tau$ and the spectral galaxy type.

### 4.4.4 Colour-colour diagram

Figure 4.19 shows the colour-colour diagram of the sample galaxies, combined with the theoretical evolutionary tracks of the model colours. Randomly for some galaxies are also shown the representative error bars of the photometry. The red lines in the diagram correspond to the evolution with time of the colour using the different models without the application of dust reddening. The current data point of the models at 13.7 Gyr is indicated by a big black data point. In the figure, the different values of $\tau$ are colour coded. The green data points correspond to galaxies that are fitted with the "Sm-model".
Data points that are not located in the regions of the bulk of the galaxies have relatively large errors. Since the errors of photometry enter into the fitting routine of GALEV/GAZELLE, the results for these galaxies should be handled with care. A check of the galaxies with large error bars reveals mostly galaxies classified as dI with low S/N.
As one can see, observational data points with ( $u-r$ ) $\gtrsim 1.75$ mag are mostly fitted by models with $\tau \leq 6 \mathrm{Gyr}$ and vice versa for galaxies with ( $\mathrm{u}-\mathrm{r}$ ) $<1.75 \mathrm{mag}$, which are fitted with models with $\tau>6$ Gyr. This behaviour shows again the connection between $\tau$ and the spectral galaxy type.


Figure 4.14: Mass distribution of the different morphological types. For all galaxies a constant distance of 16.5 Mpc was applied.

13.7 Gyr
$\tau=0.5 \mathrm{Gyr}$
$\tau=1.5 \mathrm{Gyr}$
$\tau=5.0 \mathrm{Gyr}$
$\tau=8.0 \mathrm{Gyr}$
$\tau=10.0 \mathrm{Gyr}$

Figure 4.15: Evolution of the gas metallicity of an undisturbed galaxy for different sets of $\tau$. The black horizontal line indicates the solar metallicity and the vertical line shows the position of $t=13.7$ Gyr.


Figure 4.16: Same as Fig. 4.15, but for a truncation.


Figure 4.17: Mass-metallicity relation of the GALEV-models. Data points display the median values and the $1 \sigma$ deviations of the galaxies sample within the corresponding metallicity bin. The horizontal line corresponds to the solar metallicity.


Figure 4.18: $\tau$ vs. the stellar mass $\mathrm{M}_{\text {stellar }}$ and metallicity Z of the GALEV-models. As in Fig. 4.17, the horizontal line corresponds to the solar metallicity.


Figure 4.19: (u-r)-(gi) diagram. Shown are the colours of the observed galaxies and the theoretical predictions of the GALEV models. Red lines correspond to the colour evolution of the different models and the big black data points display the colours of the models at 13.7 Gyr (today), without the application of any offsets. Colour coded are the different values of $\tau$, while the "Sm-model" is displayed by green data points.

### 4.4.5 The time evolution of the CMDs

Figures 4.20 to 4.26 show the CMDs of the sample galaxies at different time steps from 13.7 Gyr to 15.7 Gyr in steps of 0.5 Gyr . Additionally, the sample of ETGs from Janz and Lisker (2009) is shown for comparison (black data points). The upper panels of each morphological type show the undisturbed evolution without any interaction with the cluster, while the lower panels show the evolution with an applied truncation starting at 13.7 Gyr and a truncation time scale of 150 Myr .
In case of the truncation most of the late-type galaxies lie on distinct lines after 2 Gyr - e.g. at $(\mathrm{g}-\mathrm{i})=0.78$ mag and $(\mathrm{g}-\mathrm{i})=0.95 \mathrm{mag}$ in case of the BCDs. This artificial behaviour can be explained by the time evolution of the ( $\mathrm{g}-\mathrm{i}$ )-colours in Fig. 4.8. For instance, the difference of today's colours for a model with $\tau=6 \mathrm{Gyr}$ to a model with $\tau=8 \mathrm{Gyr}$ is $(\mathrm{g}-\mathrm{i})=0.2 \mathrm{mag}$. However, after 2 Gyr the difference between these two models is almost negligible with $(\mathrm{g}-\mathrm{i})=0.05 \mathrm{mag}$. Even though modelled with different values of $\tau$, these galaxies will share the same region in the colour space.
The evolution of the CMDs of the BCDs is shown in Fig. 4.20. The upper panels show the undisturbed evolution of the BCDs. In contrast to the other morphological galaxy types, the BCDs are fitted with models that have additional bursts superposed on the SFH (see again Fig. 4.2). To account for the different phases of the starburst, the onset of the starburst varies from 13.60 Gyr to 13.68 Gyr.
Due to the star formation of today's BCDs, they have bluer (g-i)-colours than the comparison sample of the dEs. Even after 2 Gyr of undisturbed evolution the BCDs have on average bluer (g-i)-colours than the dEs, indicating the ongoing star formation in most of the BCDs. However, for the undisturbed evolution of the BCDs, there are about eight BCDs after 2 Gyr that share the same region in CMD as the dEs.
If the SFR is decreased due to the truncation, which starts today, the resulting CMD after 2 Gyr looks quite different. Two Gyr after the truncation, the colours of the BCDs are in the same region as the dE population. However, the bulk of the evolved BCDs are still slightly bluer than the dEs (see also discussion of Chapter 2).
The dIs in Fig. 4.21 show almost the same behaviour as the BCDs. On average the dIs have bluer (g-i)-colour after 2 Gyr of undisturbed evolution, while in the case of a truncation most of the dIs are located at three distinct lines at ( $\mathrm{g}-\mathrm{i}$ ) $=0.79 \mathrm{mag}$, ( g i) $=0.87 \mathrm{mag}$ and $(\mathrm{g}-\mathrm{i})=0.90 \mathrm{mag}$. These lines are well located within the region of the dEs.
The undisturbed evolving spiral galaxies in Figs. 4.22 to 4.26 show only a weak evolution in the CMDs. This can be explained by the usage of larger values of $\tau$


Figure 4.20: Evolution of the BCDs (red data points) at different time steps of the GALEV simulation. Upper panel shows the undisturbed evolution, while the lower panel displays the evolution with a truncation at 13.7 Gyr . The evolution of the galaxies is shown in time steps of 0.5 Gyr , starting from $\mathrm{t}=13.7 \mathrm{Gyr}$ (today) up to $\mathrm{t}=15.7 \mathrm{Gyr}$. For comparison the results of Janz and Lisker (2009) for Virgo ETGs are shown in black data points.


Figure 4.21: Same as Fig. 4.20, but for the dIs.


Figure 4.22: Same as Fig. 4.20, but for the Sa galaxies.
by GAZELLE. If the SFR is almost constant over the next 2 Gyr , no large variations in the colours can be expected.
The evolution of the spiral galaxies in the case of the truncation is more obvious. It can also be seen in Figs. 4.22 to 4.26 that the changes in the CMDs from 13.7 Gyr to 14.2 Gyr are very prominent, while the evolution from 14.2 Gyr to 15.7 Gyr is minor. This is reasonable, since the truncation time scale is chosen to be very short with $\tau_{\text {trunc }}=150$ Myr. When the star formation stops due to the lack of gas, the massive stars will end up their life on a short time scale of several million years, what in turn will redden the integrated colours of the galaxies due to the dominance of the still remaining intermediate and low-mass stars.
In case of the Sa galaxies, one might have the impression that some of the galaxies became bluer after the truncation starts. The reason for this behaviour can be found in Equation 4.4 and in the colours of the used model grid for the truncation. When the truncation starts, the SFR of the model will decrease and therefore also the third term in Equation 4.4 with the extinction will become less important. Thus, the (g-i)-colour of the model grid only reaches a maximum of $(\mathrm{g}-\mathrm{i})=1.1 \mathrm{mag}$ (see Fig. 4.6) and with that almost all Sa galaxies will gather around this (g-i)-colour in the lower panel of Fig. 4.22.
The brighter truncated Sb galaxies ( $\mathrm{M}_{\mathrm{r}}<-19 \mathrm{mag}$ ) are on average redder than the faint Sb's and are also located in the same region as the ETGs in Fig. 4.23. The opposite is the case for the truncated Sc galaxies in Fig. 4.24. The most brightest Sc's $\left(\mathrm{M}_{\mathrm{r}}<-20.6 \mathrm{mag}\right)$ also have the bluest colours of $(\mathrm{g}-\mathrm{i})=0.78$ mag. Sc galaxies with $(g-i)>0.79$ cover the same region in the CMD as the dEs, while for galaxies with $(\mathrm{g}-\mathrm{i})=0.79$, only the faintest ones $\left(\mathrm{M}_{\mathrm{r}}>-16 \mathrm{mag}\right)$ share the same region.
The lower panel of Fig. 4.25 shows the evolution of the Sd galaxies in the case of a truncation. The Sd galaxies can be found in the region of $0.79<(\mathrm{g}-\mathrm{i})<0.95 \mathrm{mag}$ and $\mathrm{M}_{\mathrm{r}}>-18 \mathrm{mag}$, with one exception at $(\mathrm{g}-\mathrm{i})=1.06 \mathrm{mag}$ and $\mathrm{M}_{\mathrm{r}}=-18.4 \mathrm{mag}$. All Sc galaxies with $(\mathrm{g}-\mathrm{i})>0.9 \mathrm{mag}$ are located in the same regime as the dEs. At a colour of ( $\mathrm{g}-\mathrm{i}$ ) $=0.79 \mathrm{mag}$, only the faintest galaxies ( $\mathrm{M}_{\mathrm{r}}>-17 \mathrm{mag}$ ) match with the dEs. However, all Sc galaxies are well within the $2 \sigma$ deviation of the dEs.
The Sm galaxies with a truncation show the same behaviour as the Sc galaxies. In the CMD, they are clearly in the same location as the dEs after a truncated evolution of 2 Gyr (Fig. 4.26).
The loosely classified S galaxies are shown in Fig. 4.27. The evolution of the undisturbed and truncated S galaxies is less, thus one can conclude that the galaxies
have been modelled with a low gas mass fraction.
Figure 4.28 displays the evolution of the Amorphous and unknown galaxies within the CMD. After the truncation, the two unknown galaxies have the same ( $\mathrm{g}-\mathrm{i}$ )-colour $((\mathrm{g}-\mathrm{i})=0.86 \mathrm{mag})$ and also the same magnitudes with values of $\mathrm{M}_{\mathrm{r}}=14.3 \mathrm{mag}$ and $\mathrm{M}_{\mathrm{r}}=14.6 \mathrm{mag}$ after 2 Gyr . In the class of the Amorphous galaxies, only for two galaxies a significant influence due to the truncation can be observed. Four out of five of the truncated Amorphous galaxies are well located in the same region as the dEs.

Figures 4.29 and 4.32 show the average r-band magnitudes and colours of the different morphological types after an evolution of 2 Gyr. The averages on the left hand side of the figures are for the undisturbed evolving galaxies, without any interaction. The right hand side of the figures show the averages after 2 Gyr with the application of a truncation at $\mathrm{t}_{\text {trunc }}=13.7$ Gyr. Additionally plotted are the averages of today's early-type galaxy population, subdivided into the contribution of the entire population (ETG), dEs, $\mathrm{dE}(\mathrm{bc})$ and dEs with a disc structure (dE(di), Lisker et al., 2006b).
Regarding the magnitudes, the morphological types of Sa and Sb are on average brighter than the entire ETG population for both (undisturbed and truncated) evolutionary scenarios. We also found that the Sa galaxies have too red colours compared to the ETGs, which may seem disconcerting. In this case, one has to take care about the calculation of the averages, since the ETGs do not follow a constant $(\mathrm{g}-\mathrm{i})-\mathrm{M}_{\mathrm{r}}$ relation in the CMD, but show " S "-shape slope (see Janz and Lisker, 2009) in the way that fainter dEs are bluer than the bright ellipticals. Therefore, the results of the Sa galaxies are not in contradiction with other observations.
The Sd and Sm galaxies show a large overlapping region with the ETGs for an undisturbed evolution after 2 Gyr. With an additionally truncation, the r-band magnitude and (g-i)-(u-r)-colours are in an even better agreement with the ETGs. However, the (i-z)-colours, which are most sensitive to the metallicity, show significant differences to the ETGs. The ETGs have lower (i-z)-colours ( $(\mathrm{i}-\mathrm{z}) \approx 0.13 \mathrm{mag})$ compared to the undisturbed Sd and Sm galaxies $((\mathrm{i}-\mathrm{z}) \approx 0.21 \mathrm{mag})$, but these (i-$\mathrm{z})$-colours are within the $1 \sigma$ deviation of the ETGs. In case of a truncation the Sd and $\operatorname{Sm}$ galaxies still have higher (i-z)-colours $((\mathrm{i}-\mathrm{z}) \approx 0.18 \mathrm{mag})$ than the ETGs, but the differences are now less prominent. These differences in the (i-z)-colours are also visible for all other late-type galaxies in both evolutionary scenarios in Fig. 4.32.
After 2 Gyr of undisturbed evolution, the dIs are bluer and slightly fainter than the ETGs. The difference in magnitude become even larger, when a truncation is applied. Thus, the dIs are on average fainter than the ETGs, but with an overlap region regarding


Figure 4.23: Same as Fig. 4.20, but for the Sb galaxies.


Figure 4.24: Same as Fig. 4.20, but for the Sc galaxies.


Figure 4.25: Same as Fig. 4.20, but for the Sd galaxies.


Figure 4.26: Same as Fig. 4.20, but for the Sm galaxies.


Figure 4.27: Same as Fig. 4.20, but for the S galaxies.


Figure 4.28: Same as Fig. 4.20, but for the unknown and Amorphous galaxies.


Figure 4.29: Average values of the r-band magnitudes after 2 Gyr with an undisturbed (left) and truncated (right) evolution, respectively. Error bars indicate the $1 \sigma$ deviation. In case of the ETGs, dEs and dE(bc)s the averages of today's galaxies are used for comparison. The red and blue lines correspond to the averages of the dEs and $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$. Also shown in a pink bar are the averages of dEs in which discs were identified by Lisker et al. (2006b) (so called dE(di)).
the $1 \sigma$ deviations.
In the case of the "unknown", "S" and Amorphous galaxies, any given average suffers from the small number of galaxies within the bin. Additionally, these types are only loosely classified and do not belong to the standard scheme of the morphological types. Therefore, these three morphological types will not be included into the discussion about the possible progenitors of early-type dwarf galaxies.


Figure 4.30: Same as Fig. 4.29, but for the (g-i)-colours.


Figure 4.31: Same as Fig. 4.29, but for the (u-r)-colours.


Figure 4.32: Same as Fig. 4.29, but for the (i-z)-colours.

### 4.4.6 The evolution of the surface brightness profiles

As seen in the last section, the r-band magnitudes and the ( $\mathrm{g}-\mathrm{i}$ )-colours of some latetype galaxies fit very well to the parameters of the dEs, if a truncation is assumed. However, the matching of these two parameters does not imply that late-type and early-type are structurally similar. Therefore, the investigation of the evolution of the effective radius $\mathrm{R}_{\text {eff }}$ and the mean surface brightness $\langle\mu\rangle_{\text {eff }}$ is of particular interest. Since GALEV has no spacial resolution, it is not possible to simply derive the evolution of $\mathrm{R}_{\text {eff }}$ and $\langle\mu\rangle_{\text {eff }}$. Therefore, we divided every galaxy into four regions with $0.5,1.0$, 1.5 and 2.0 half light semi-major axis $\mathrm{a}_{\mathrm{h}, \mathrm{r}}$ in r -band (see Fig. 4.33). The shape of the galaxy ( $a_{\text {hl, }}$, ellipticity $\epsilon$ and position angle p.a.) was determined by using the method described in Chapter 3.2.4.
The different fluxes of the galaxy were measured within the area $A_{i}$ with the limiting $\operatorname{radii} R_{i-1}$ and $R_{i}(i=1 \ldots n)$. Using the corresponding magnitude of the area $A_{i}$ and the radius

$$
\begin{equation*}
\mathrm{R}_{\mathrm{i}}^{\prime}=\frac{\left(\mathrm{R}_{\mathrm{i}}+\mathrm{R}_{\mathrm{i}-1}\right)}{2} \tag{4.5}
\end{equation*}
$$

a surface brightness profile (SBP) was created and for simplicity this SBP was fitted by an exponential law with Sérsic-index of $\eta=1$ :

$$
\begin{equation*}
\mu(\mathrm{R})=\mu_{0}+1.086 \cdot\left(\frac{\mathrm{R}}{\alpha}\right)^{1 / \eta} \tag{4.6}
\end{equation*}
$$

with the central surface brightness $\mu_{0}$ and the scale length $\alpha$. This exponential fit was used to obtain a first-order approximation of the change of $\mathrm{R}_{\text {eff }}$ and $\langle\mu\rangle_{\text {eff }}$ due to the different evolutionary scenarios.
The surface brightness of the area $A_{i}$ at the radius $\mathrm{R}_{\mathrm{i}}^{\prime}$ was calculated with the following equation:

$$
\begin{equation*}
\mu_{\mathrm{i}}\left(\mathrm{R}_{\mathrm{i}}^{\prime}\right)=\mathrm{m}_{\mathrm{i}}+2.5 \cdot \log \left[\pi \cdot\left(\mathrm{R}_{\mathrm{i}}^{2}-\mathrm{R}_{\mathrm{i}-1}^{2}\right)\right], \tag{4.7}
\end{equation*}
$$

with the magnitude $m_{i}$ of the area $A_{i}$.
The magnitudes of each area $A_{i}$ in the five SDSS filters are then used as input parameters for a new GAZELLE run to simulate the evolution of the magnitudes within the different areas $\mathrm{A}_{\mathrm{i}}$, using the model grid of Tab. 4.1. As a result we obtain the evolution of the SBP at the different radii $\mathrm{R}_{\mathrm{i}}^{\prime}$. The new profile was again fitted by an exponential law with $\eta=1$, resulting in new values of the central surface brightness $\mu_{0}$ and the scale length $\alpha$ of the linear fit at different times, ranging from 13.7 Gyr (today,
$\mathrm{t}=0 \mathrm{Gyr}$ ) to $15.7 \mathrm{Gyr}(\mathrm{t}=2 \mathrm{Gyr})$. For all galaxies we simulate the undisturbed evolution and the evolution with a truncation.
Using $\mu_{0}$ and a Sérsic-index of $\eta=1$, we are able to derive $\langle\mu\rangle_{\text {eff,r }}$ at different times, with several equations given in Graham and Driver (2005):

$$
\begin{equation*}
\langle\mu\rangle_{\mathrm{eff}, \mathrm{r}}=\mu_{\mathrm{e}}-0.699 \tag{4.8}
\end{equation*}
$$

with effective surface brightness ${ }^{5} \mu_{\mathrm{e}}$ at $\mathrm{R}_{\text {eff }}$, defined as:

$$
\begin{equation*}
\mu_{\mathrm{e}}=\mu_{0}+8.327 . \tag{4.9}
\end{equation*}
$$

To calculate $\mathrm{R}_{\text {eff }}$, we used the scale length $\alpha$ of the SBP and the formula

$$
\begin{equation*}
\mathrm{R}_{\mathrm{eff}}=1.678 \cdot \alpha, \tag{4.10}
\end{equation*}
$$

what finally leads to the total magnitude $m_{\text {tot }}$ of the entire galaxy:

$$
\begin{equation*}
\mathrm{m}_{\mathrm{tot}}=\langle\mu\rangle_{\mathrm{eff,r}}-2.5 \cdot \log \left(2 \pi \cdot \mathrm{R}_{\mathrm{eff}}^{2} \cdot \sqrt{b / a}\right) . \tag{4.11}
\end{equation*}
$$

The axis ratio $\mathrm{b} / \mathrm{a}$ was used to account for the elliptical shape of the galaxy and was determined in the course of Chapter 3.
Figures 4.34 and 4.35 show the evolution of the SBPs of two galaxies (VCC-0004 and VCC-0017) within several time steps and evolutionary scenarios. The red data points in the figures display the surface brightness at different radii $\mathrm{R}_{\mathrm{i}}^{\prime}$ of the galaxies, which was measured from the observations. The blue lines correspond to linear fits to the SBPs for an undisturbed evolution after 2.0 Gyr. In contrast to this, the black lines show the SBPs after 2.0 Gyr with an applied truncation. The influence to the SBP of an undisturbed galaxies is minor, resulting in a minor change in $\mathrm{R}_{\text {eff }}$. The increase of the SBP can be explained by the ongoing SF, which in case of the Sm-model will increase over the next 2.0 Gyr.
In case of the truncation the change of the SBP is considerable. The central surface brightness changes from $\mu_{0 ; 0004}=22.0 \mathrm{mag} / \operatorname{arsec}^{2}\left(\mu_{0 ; 0017}=22.1 \mathrm{mag} / \mathrm{arsec}^{2}\right)$ for the present SBP to $\mu_{0 ; 0004}=23.0 \mathrm{mag} / \operatorname{arsec}^{2}\left(\mu_{0 ; 0017}=23.1 \mathrm{mag} / \mathrm{arsec}^{2}\right)$ for the SBP in 2.0 Gyr with a truncation. Therefore, $\mathrm{R}_{\text {eff }}$ changes from $\mathrm{R}_{\text {eff; }}$;004 $=0.82 \mathrm{kpc}$ $\left(\mathrm{R}_{\text {eff; } 0017}=1.38 \mathrm{kpc}\right)$ to $\mathrm{R}_{\text {eff; } 0004}=0.87 \mathrm{kpc}\left(\mathrm{R}_{\text {eff; }} 0017=1.54 \mathrm{kpc}\right)$. A summary of the results of each galaxy can be found in Tab. 6.6 of the appendix.

[^23]Table 4.3 shows the median variation of $\mathrm{R}_{\mathrm{eff}}$ in kpc for an undisturbed evolution and a truncation, respectively. For each morphological type, the median variation of $\mathrm{R}_{\text {eff }}$ was calculated. Also displayed are the $1 \sigma$ deviations from the medians. It has to be pointed out that the deviation are large and therefore any interpretation has to be treated with caution.
After 2.0 Gyr of an undisturbed evolution, the BCDs show a slightly decrease of $\mathrm{R}_{\text {eff }}$ of $3.9 \%$. If we apply additionally a truncation of the SFR , the BCDs became larger and therefore $\mathrm{R}_{\text {eff }}$ increases by $12 \%$.
In case of the dIs, the effective radius remains almost constant for an undisturbed evolution after a period of 2.0 Gyr . Using the truncation scenario, $\mathrm{R}_{\text {eff }}$ increases by about $19.4 \%$. However, the $1 \sigma$ deviation is large with $13.4 \%$.
In contrast to the star forming dwarf, the spiral galaxies show a different behaviour in the evolution of the effective radius. The undisturbed spirals of type $\mathrm{Sa}, \mathrm{Sc}$ and Sd show almost no change in $R_{\text {eff }}$, while the Sb galaxies increase by $\sim 6 \%$. In case of the undisturbed Sm and " S " spirals, the change in $\mathrm{R}_{\text {eff }}$ is negative, meaning that $\mathrm{R}_{\mathrm{eff}}$ decreases by $-7.1 \%$ and $-6.1 \%$, respectively.
If we apply a truncation to the spiral galaxies, the types from Sa to Sb show an opposite behaviour than the BCDs and dIs, in a way that $\mathrm{R}_{\text {eff }}$ is getting smaller by $-3.7 \%$ and $-7.9 \%$, respectively. In contrast to this, the change of $\mathrm{R}_{\text {eff }}$ for the $\mathrm{Sc}, \mathrm{Sd}, \mathrm{Sm}$ and S galaxies is positive and in the order of a few percent ( $<10 \%$ ). This behaviour is reasonable for the early spiral galaxies of type Sa and Sb , due to the existence of the prominent red bulge in the centre of the galaxy and the surrounding blue disc. Therefore, GAZELLE modelled the bulge with a small $\tau$, which has almost no gas and an old stellar population, while the disc is modelled with larger values of $\tau$. As a result, the colours of the gas abundant discs are more affected by the truncation than the inner bulge regions, which in turn explains the decrease in $\mathrm{R}_{\text {eff }}$ of the early spiral galaxies. The Sm galaxies show a similar evolution of $\mathrm{R}_{\text {eff }}$ as the star forming dwarfs, especially in case of the BCDs. The radius of the undisturbed Sm galaxies decreases by $-7.1 \%$, but increases by $5.2 \%$ for the truncated evolution. This result for the evolution of $\mathrm{R}_{\text {eff }}$ for the Sm galaxies is expected, since they have similar properties as the star forming dwarf galaxies.
The "unknown" galaxies show a slightly change in $\mathrm{R}_{\text {eff }}$ in case of an undisturbed evolution, but show the largest change of $\mathrm{R}_{\text {eff }}$ of the entire morphological sequence, with $28.9 \%$. However, it should be kept in mind that these results are based on only a small number of galaxies, which could strongly influence the statistic.

The effective radius of the Amorphous galaxies is almost constant for an undisturbed evolution and is only noticeable in the fourth post decimal position. In case of a truncation, $\mathrm{R}_{\text {eff }}$ increases by $13.5 \%$. But as already mentioned for the S and unknown galaxies, these findings only have a low number statistic.
In Tab. 4.4, the median variations of $\langle\mu\rangle_{\text {eff,r }}$ are shown in the same manner as for Tab. 4.3. In this table a negative change in $\langle\mu\rangle_{\mathrm{eff}, \mathrm{r}}$ implies an increase of the surface brightness (SB) and vice versa for positive changes.
The BCDs show a small decrease of the surface brightness for the undisturbed evolution, while in the case of a truncation the SB decreases by about $0.9 \mathrm{mag} / \mathrm{arcsec}^{2}$. The dIs show an opposite behaviour, with an increase of the SB for the undisturbed evolution by about $-0.6 \%$ and a decrease by $3.5 \%$ in case of a truncation.
The SB for the spirals of type Sa and Sb decreases for an undisturbed evolution, while the SB for $\mathrm{Sc}, \mathrm{Sd}, \mathrm{Sm}, \mathrm{S}$ and unknown galaxies increases, and vice versa for the truncation. Overall, the differences between the SB of an undisturbed and truncated evolution are minor - in all of the case below $1 \mathrm{mag} / \mathrm{arcsec}^{2}$ - for all morphological types.
Figure 4.36 compares the results of the SBP analysis of the different morphological types and in addition the results of today's dEs, $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$, ETGs and LSB-components (see Chapter 2) are shown. To obtain the change of the effective radius and mean effective surface brightness of today's late-type galaxies due to the undisturbed and truncated evolution, we use the information of Tabs. 4.3 and 4.4, and applied the percentage terms of the changes of the effective radius ( $\Delta \mathrm{R}_{\text {eff }}$ ) and mean effective surface brightness $\left(\Delta\langle\mu\rangle_{\text {effr, }}\right)$. For instance, the dIs within the Virgo cluster have an average observed effective radius of $\mathrm{R}_{\text {eff }}=1.6 \mathrm{kpc}$ and change their radius after 2 Gyr due to a truncation by about $19 \%$, which in turn results in a new effective radius of $\mathrm{R}_{\text {eff }, 2 \text { Gyr,trunc }}=1.9 \mathrm{kpc}$.


Figure 4.33: Schematic division of the SBP.


Figure 4.34: Evolution with time of a SBP of galaxy VCC-0004.


Figure 4.35: Evolution with time of a SBP of galaxy VCC-0017.


Figure 4.36: Effective radius $\mathrm{R}_{\text {eff }}$ and mean effective surface brightness $\langle\mu\rangle_{\mathrm{eff}, \mathrm{r}}$ at different evolutionary stages of the morphological types. Today's values with the $1 \sigma$ deviations are displayed by a black bar. The red and blue bars correspond to the same morphological types but after an undisturbed evolution after 2 Gyr (red bar) and after 2 Gyr with an additional truncation (black). Furthermore, today's results of the ETGs, $\mathrm{dEs}, \mathrm{dE}(\mathrm{bc}) \mathrm{s}$ and LSB-components are displayed, but without any time evolution.

Table 4.3: Change of $\mathrm{R}_{\text {eff }}$ after 2 Gyr for undisturbed evolution and evolution with truncation.

| Type | Number | $\Delta \mathrm{R}_{\text {eff,2Gyr }}[\mathrm{kpc}]$ | $\sigma[\%]$ | $\Delta \mathrm{R}_{\text {eff,2Gyr,trunc }}[\mathrm{kpc}]$ | $\sigma[\%]$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| BCDs | 33 | $-0.021(-3.9 \%)$ | 5.046 | $0.064(12.0 \%)$ | 6.061 |
| Im | 87 | $0.005(0.4 \%)$ | 11.095 | $0.260(19.4 \%)$ | 13.394 |
| Sa | 35 | $0.008(0.3 \%)$ | 19.846 | $-0.167(-6.0 \%)$ | 17.507 |
| Sb | 17 | $0.125(5.8 \%)$ | 21.617 | $-0.206(-9.7 \%)$ | 19.842 |
| Sc | 79 | $0.029(1.3 \%)$ | 19.456 | $0.043(1.9 \%)$ | 16.970 |
| Sd | 21 | $0.000(0.0 \%)$ | 12.180 | $0.204(9.6 \%)$ | 11.380 |
| Sm | 26 | $-0.116(-7.1 \%)$ | 10.255 | $0.085(5.2 \%)$ | 10.898 |
| S | 5 | $-0.087(-6.1 \%)$ | 10.872 | $0.054(3.8 \%)$ | 10.245 |
| unknown | 3 | $0.046(8.6 \%)$ | 3.993 | $0.153(28.9 \%)$ | 6.924 |
| amorph | 5 | $0.000(0.0 \%)$ | 7.247 | $0.247(13.7 \%)$ | 7.033 |

Table 4.4: Change of $\langle\mu\rangle_{\text {eff }}$ after 2 Gyr for undisturbed evolution and evolution with truncation.

| Type | Number | $\Delta\langle\mu\rangle_{\text {eff } 2 \text { 2Gyr }}$ | $\sigma$ | $\Delta\langle\mu\rangle_{\text {eff,2Gyr,trunc }}$ | $\sigma$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| BCDs | 33 | $0.122(0.6 \%)$ | 1.342 | $0.915(4.3 \%)$ | 1.316 |
| Im | 87 | $-0.151(-0.6 \%)$ | 0.891 | $0.830(3.5 \%)$ | 0.950 |
| Sa | 35 | $0.294(1.4 \%)$ | 0.720 | $-0.033(-0.2 \%)$ | 1.042 |
| Sb | 17 | $0.143(0.7 \%)$ | 0.728 | $-0.305(-1.5 \%)$ | 1.078 |
| Sc | 79 | $-0.018(-0.1 \%)$ | 0.762 | $0.230(1.1 \%)$ | 1.021 |
| Sd | 21 | $-0.085(-0.4 \%)$ | 0.504 | $0.615(2.7 \%)$ | 0.740 |
| Sm | 26 | $-0.290(-1.3 \%)$ | 0.912 | $0.603(2.7 \%)$ | 0.986 |
| S | 5 | $-0.087(-0.4 \%)$ | 0.831 | $0.302(1.4 \%)$ | 0.841 |
| unknown | 3 | $-0.119(-0.5 \%)$ | 1.733 | $0.631(2.7 \%)$ | 2.130 |
| amorph | 5 | $0.175(0.9 \%)$ | 1.245 | $0.188(0.9 \%)$ | 1.201 |

### 4.5 Discussion

### 4.5.1 Classification by models

In Chapter 3 the galaxies were reclassified based on their morphology and distribution in the parameter-space. In Fig. 4.12, there are four Sa galaxies that are modelled with $\tau=25 \mathrm{Gyr}$. This value of $\tau$ is relatively large for Sa galaxies, since they commonly formed their stellar content within the first Gyrs (Sandage, 1986; Gavazzi et al., 2002). Interestingly, when using the number distribution of the morphological types based on the original morphological types of the VCC, then also misclassified ${ }^{6}$ galaxies are included, which in turn should be noticeable in the results of the SEDs. Apart from the check of the classification with the results of the photometry, we are now able to check the classification on the basis of the SED fitting routine.
In the course of Chapter 3, three out of four conspicuous Sa galaxies were reclassified as $\mathrm{Sc}, \mathrm{Sm}$ and $\mathrm{dE}(\mathrm{bc})$, respectively. This finding is also supported by the analysis of the SED fitting. For example, the SED of VCC-0015 on the left hand side of Fig. 4.37 which was initially classified as Sa galaxy by the VCC and GOLDMine - shows large differences from the SEDs of other Sa galaxies. From its morphology and position in the parameter-space, VCC-0015 was reclassified as Sc galaxy and indeed the SED supports this reclassification to a later spectral type.
Only one Sa galaxy (VCC-0979) is still puzzling. The SED on the right panel of Fig. 4.37 shows that the $r$, $i$ and $z$ filters have similar properties as the other SEDs of Sa galaxies in the same wavelength range. However, the $u$ and $g$ filters do not fit to the SEDs of this spectral sequence of the Sa galaxies. The inspection of the gri-SDSS image in Fig. 4.38 reveals a disc galaxy with apparent blue regions in the core. The blue colours of the regions point to an ongoing or recent star formation, which therefore would explain the enhanced magnitudes in the $u$ and $g$ filter. From this perspective, it also points to a misclassification of VCC-0979, but the exact type is still unknown.

### 4.5.2 Models vs. observation

Looking at the results of our analysis in the CMDs of Figs. 4.20 to 4.28, one realises - especially in the case of the truncation - that the models are too simple. When modelling the galaxies with a truncation event, the galaxies fall on discrete lines in the

[^24]

Figure 4.37: Comparison of the SEDs for Sa galaxies, which were modelled with different values of $\tau$. The differences in the SEDs indicate a misclassification in the case of VCC-0015 (left) and VCC-0979 (right).


Figure 4.38: gri-SDSS image of VCC-0979, which shows blue region in the centre of the galaxy.

CMDs, which therefore gives an artificial impression.
As e.g. shown by McQuinn et al. (2010a) and McQuinn et al. (2010b), the SFHs of galaxies do not have to be smooth curves like the ones of our study. Also the assumption that the truncation time scale is constant for all different morphological types might be too simple. For instance, the ram pressure stripping event, which is responsible for the truncation, is not only able to significantly decrease the SFR of a galaxy, but it is also able to induce new star formation (Gavazzi et al., 1995; Fujita and Nagashima, 1999). Thus, the decreasing SFR will be a superposition a truncation and star formation, which may increase the truncation time scale to values with $\tau_{\text {trunc }}>150$ Myr. However, a larger $\tau_{\text {trunc }}$ will result in bluer ( g - i -colours for the galaxies after 2 Gyr , which in turn also influences the comparison with dEs.
Another simplification of the models concerns the gas content of the galaxies after the truncation. In case of a truncation, GALEV just decreases the SFR but does not "physically" create a gas-poor galaxy. Therefore, the behaviour of the gas metallicity of GALEV and the reality are different. The metallicity of the remaining gas of a real galaxy would increase due to the enrichment by the yields of SNe after the truncation event, which is not the case for the GALEV models.
The evolution of BCDs was modelled by a SFH that has one single strong starburst, superposed on a "normal" SFH given by Equation 4.1. However, it is commonly believed that the SFH of BCDs shows oscillation in the SFR, from episodes with short and strong bursts to long quiescence phases (Thuan, 1991). Our models for the BCDs do not account for the bursty evolution of the SFR, but it was also shown that the influence of a strong burst is less and almost not visible in the ( g - i -colours after several Gyr (see Fig. 4.7). Thus, the induced errors by this simplification of the SFH of BCDs may have a minor effect on the overall evolution of the photometric parameters of the BCDs.

Another issue one has to consider is the mass dependence of mechanisms, which remove the gas from the late-type galaxies. Following Mo et al. (2010), the binding force, which prevents the gas to be removed from a disc galaxy, is proportional to the mean surface density of the ISM and the mean mass density of the disc. If two galaxies move through the cluster with the same velocity ${ }^{7}$, the one with the lower density is more influenced by the stripping event. In case of a starvation the galaxy will proceed its SF over a longer time scale due to the still existing gas reservoir in the disc, which

[^25]again increased $\tau_{\text {trunc }}$. Therefore, also the results of simulations of spiral galaxies in the Virgo cluster (van Gorkom, 2004; Fumagalli and Gavazzi, 2008; Vollmer et al., 2012) should be included in the treatment of the environmental effects, which act on the late-type population. However, an increase of $\tau$ will only affect the parameters in the arbitrarily chosen time span of 2 Gyr. If we take much larger time spans of about 4 Gyr or even longer, then the environmental effects will also remove the remaining gas. Thus, we do not have to ask if the derived parameter-space is reached by different sets of $\tau_{\text {trunc }}$, but when the parameter-space is reached.
Also related to the mass of the galaxies is the completeness limit of the VCC and the applied magnitude cut-off ( $\mathrm{M}_{\mathrm{B}} \leq-13.09 \mathrm{mag}$ ). Especially in the case of the dI one has to account for this when computing average model masses. Due to the magnitude cut-off, low surface brightness dIs are not included in our study, which could influence the derived average model masses.
The evolution of the SBPs of the late-type sample shows that an undisturbed evolving galaxy will only slightly change its effective radius within the next 2.0 Gyr. However, in the case of a truncation, which in turn reduces the SFR, the effects become more prominent for the star forming dwarf galaxies of the type BCD and dI. For these star forming dwarfs the effective radius does indeed change within 2.0 Gyr , but only on the order of a few percent. However, these calculations are only a first-order approximation, since all SBPs are approximated by an exponential law with a Sérsicindex of $\eta=1$, which definitely does not describe the real shape of all galaxies in the sample. Thus, a possible improvement in the fitting procedure of the SBPs would be the inclusion of the Sérsic-index as a free parameter in the fitting procedure of the SBPs.
In agreement with our study, Boselli et al. (2008) also stated that the effective radius is almost constant at larger wavelength when the galaxies evolve with time. To increase our understanding of the evolution of the effective radius, but also of the other parameters, detailed simulations are need for the entire morphological sequence of late-type galaxies. Such studies on the effect of the environment on spiral galaxies in the Virgo cluster already exist (e.g. Vollmer, 2009), but further studies, including the low mass galaxies like dIs and BCDs, are desirable in the future.

### 4.5.3 The possible progenitors of dEs

Using all the derived results of our analysis, we are finally able to investigate the possible progenitors of today's dEs by comparing the structural properties of the
different morphological types. In Chapter 2 it was shown that the parameters of the BCDs - when the contribution of the starburst region is subtracted and only the old stellar population of the LSB-component is taken into account - fit very well to the population of the dEs. Furthermore, we found that the LSB-components of BCDs are on average more compact and fainter than early-type dwarf galaxies with a blue core. In comparison to the LSB-components, the dIs in Virgo are less compact and have a lower surface brightness, which is noticeable by a clear separation between these two morphological types in the $\mathrm{M}_{\mathrm{r}}-\mathrm{R}_{\text {eff }}-\langle\mu\rangle_{\mathrm{eff}, \mathrm{r}}$ parameter-space. Also in terms of the spatial distribution within the Virgo cluster of the BCDs and dEs, there is a probability that these two types are related to each other.
But what about the other morphological types in the Virgo cluster and how will they look like in the future?
The time evolution of the late-type galaxies by the means of GALEV/GAZELLE models (Figs. 4.20 to 4.28 and Fig. 4.36) showed that the structural parameters $\left(\mathrm{R}_{\mathrm{eff}}\right.$, ( g - i ) and $\langle\mu\rangle_{\text {eff }}$ ) of the late-types partly fit to the ones of the early-type (dwarf) galaxies for an undisturbed evolution of two Gyr. Since the late-type galaxies are located within the cluster environment, there are several forces acting on theses galaxies, which may transform them. Therefore, a truncation of the SFR was applied, which simulates the removal of gas due to forces like ram pressure stripping. The results of the GAZELLE runs showed that the structural parameters of late-type galaxies become even more similar to the ones of the ETGs, but this finding depends strongly on the morphological type. The early spirals of type Sa and Sb are in many respects not comparable to the dEs , since they are too bright in the r-band magnitude, have too large effective radii and too high surface brightness. As expected, we can conclude that the descendants of Sa and Sb galaxies will not share the same parameter-space as today's dEs.
Looking at the spirals of type Sc and Sd, we found more overlap with dEs in Virgo. The r-band magnitude and (g-i)-colours of the Sc and Sd galaxies are between the ones of the dEs and the $\mathrm{dE}(\mathrm{bc}) \mathrm{s}$. Compared to the dEs , the Sc and Sd galaxies are rather large in radii and have a relatively high surface brightness. However, the overlap within the $1 \sigma$ deviations of the Sc and Sd to the ones of the dEs and $\mathrm{dE}(\mathrm{bc})$ s showed that they are not completely different according to our parameter-space.
Among the spirals, the Sm galaxies show the best overlap with the dEs in the entire parameter-space, even though the radii are slightly larger.
Using the metallicity sensitive (i-z)-colours, we found that there are differences between the ETGs and the late-type galaxies after two Gyr (with and without truncation),
which may points to different chemical abundances (see Fig. 4.32). However, looking in more detail at the early-types in Fig. 4.32, one realises that the Sd and Sm galaxies are within the $1 \sigma$ deviations of the (i-z)-colours of the ETGs.
The parameter-space of today's dI galaxies is in agreement with the dEs, but the evolution with time and the influence of the cluster environment change this behaviour. After a truncated evolution of two Gyr, the (g-i)-colours of the dIs are in the same region as the dEs. However, the dIs become fainter than the dEs and also the surface brightness is getting too low, even when accounting for the $1 \sigma$ deviations of both populations. Therefore, only the faintest dEs with $\langle\mu\rangle_{\mathrm{eff,r}} \approx 26 \mathrm{mag} / \operatorname{arcsec}^{2}$ have the possibility to be related to the dIs. These faint dIs might be related to the dSph galaxies, which are treated as dEs in our analysis. However, Grebel et al. (2003) showed that the mean stellar metallicity of dSph galaxies is higher than that of the dIs (see also e.g. Richer and McCall, 1995; Richer et al., 1998, for other metallicity estimates for dSphs and diffuse ellipticals), pointing to differences in the SFH of these two types. The low surface brightness field dIs of the sample of Thuan (1985) also show "mutually exclusive metallicity ranges" ${ }^{8}$ compared to the Virgo dEs. The metallicity sensitive (i-z)-colours of Fig. 4.32 show that the Virgo dEs and dIs do not differ extremely and that the (i-z)-colours of the dIs are within the $1 \sigma$ deviations of the dEs. Therefore, we are not able to give such a strong statement as Thuan (1985) regarding the different metallicities. The study of Bothun et al. (1986) also supports our findings, since their Virgo dIs have to fade by $\mathrm{m}_{\mathrm{B}}=1.5 \mathrm{mag}$ to have the same photometric properties as the dEs. Consequently, this would lead to very low surface brightnesses, which are not comparable with the Virgo dEs. Hence, our results and the above mentioned studies indicate that there are no evolutionary connections between dIs and dEs (+dSph) in the Virgo cluster, which is in contrast to other early studies (e.g. Lin and Faber, 1983; Kormendy, 1985).
In summary, we found that there are some morphological types among the late-type galaxies that share a similar region in the parameter-space as the dEs after a truncated evolution of two Gyr within the Virgo cluster. The most promising types are the late spirals of the type Sd and Sm , and the LSB-components of the BCDs, when star formation has ceased. Apart from the photometric and structural properties of the late-type and early-type galaxies of our study, one also has to consider the intrinsic kinematics of the galaxies. Galaxies with an ordered motion of their stars and/or gas will not lose their angular moment on short time scales. Therefore, the progenitors

[^26]must have similar kinematical properties as today's ETGs. In the last years, studies of dEs in the Virgo cluster (e.g. Geha et al., 2003; de Rijcke et al., 2005; Chilingarian, 2009; Toloba et al., 2011) found that parts of the dE population indeed have kinematics that point to rotational flattened or rational supported dEs (see Lisker, 2012, for a recent review). Interestingly, Toloba et al. (2011) found the tendency of two kinematically different populations of dEs in Virgo. The ones at large cluster centric radius are rotational supported and have younger ages, while dEs in the core of Virgo are pressure supported and older. Furthermore, other studies found dEs in the Virgo cluster that are rotational supported and contain HI gas (e.g. Pedraz et al., 2002; Conselice et al., 2003; van Zee et al., 2004) and together with the results of Lisker et al. (2006b), who found dEs with disc features, it points to a possible transformation of gas abundant late-type galaxies into dEs, since the dEs still have the imprints of the late-type galaxies. Harassment-simulations of model late-type dwarf disc galaxies in Virgo by Smith et al. (2010) showed that the harassed discs are able to form spiral structure, which was also confirmed by Lisker et al. (2007) in Virgo dEs from the SDSS. However, Smith et al. (2010) also pointed out that the initial model galaxies were smooth and featureless and therefore not easily comparable with typical late-type galaxies of our sample. As seen in Chapter 2, there is one BCD (VCC-0213) in the Virgo sample that shows spiral features and furthermore, these spiral arms show a colour gradient from left to right ${ }^{9}$, which could be interpreted as a sign of interaction with the dense cluster environment. Another important and interesting question is whether the late-type galaxies will lose their entire gas reservoir when moving through the Virgo cluster. The GALEV/GAZELLE settings used here do not account for gas loss due to external forces - which can be achieved by the inclusion of gas outflows - therefore we consult other studies. The study by Chung et al. (2007) on $\sim 50$ VLA HI Virgo spirals showed that seven of them, which are located at intermediate cluster centric radii, have extended HI tails, indicating the influence of the cluster environment (especially RPS). Furthermore, a subsequent study by (Chung et al., 2009) showed that galaxies at a small distance from the core of the Virgo cluster have HI discs that are smaller than the stellar disc. N-body simulations by Vollmer et al. (2001) also showed that galaxies will lose gas when moving through the core region of Virgo, however, the star formation will not completely stop and furthermore, parts of the gas, which is removed by RPS, can be re-accreted by the galaxies and may amount for up to $10 \%$.
Therefore, the above described effects of harassment and RPS are insufficient to create

[^27]the "red and dead" gas-poor dE. Interestingly, a recent study by Hallenbeck et al. (2012) (see also Conselice et al., 2003) found gas-bearing Virgo dEs, which have HI and stellar masses similar to BCDs and dI. Thus, we may also have to change our view of a "red and dead" gas-poor dE and have to focus on the structural properties of these HI dEs. But, from our analysis of Chapter 2 we also saw that dEs with a blue core do not share the same properties in the $\mathrm{M}_{\mathrm{r}}-\mathrm{R}_{\text {eff }}$ parameter-space as the LSB-components of the Virgo BCDs.
If the late-type galaxies of the type $\mathrm{BCD}, \mathrm{Sd}$ and Sm will completely (which is highly unlikely) evolve into a galaxy type that is comparable to the dE population of Virgo, it also interesting to ask whether the number distribution of these "new dEs" is in the same order as the one of today's dEs. Only accounting for galaxies in the outer part of Virgo with a distance to M87 of $\mathrm{D}_{\mathrm{M} 87}>0.6 \mathrm{Mpc}$ (see Fig. 2.24 of Chapter 2 and Fig. 3.8 of Chapter 3), we found that both populations are almost equally numerous (late-types $=25$ and $\mathrm{dEs}=27$ ). If we also take into account parts of the Sc galaxies and dIs then the late-types outnumber the dEs. However, if the number distribution of the entire Virgo cluster (see Fig. 3.3 of Chapter 3) is used then the late-types ( $\mathrm{Sc}+\mathrm{Sd}+\mathrm{Sm}+\mathrm{BCDs}+\mathrm{dI}$ ) are severely outnumbered by the dEs (amount of dEs $\sim 400$ ). Therefore, a formation of the entire dE population from in-falling late-type galaxies is not possible. Conselice et al. (2001) found that parts of the dEs in Virgo are not formed in-situ, but originate from in-falling late-types, which therefore supports our findings from a different point of view.

### 4.6 Outlook

### 4.6.1 GALEV in the zoom

At the time of writing the galaxies of the sample are treated as 1-zone objects without any change with increasing galaxy radius. However, it was shown that galaxies do change their appearance with radius. Examples are the spiral arms in late-type galaxies, but also dEs with a blue cores.
As a new feature, which is in the testing phase, GALEV/GAZELLE is now able to analyse each galaxy pixel by pixel. The procedure is the same as for the entire galaxy, but now for each single pixel or pixel area (e.g. $4 \times 4$ pixels) a best fitting model is determined by using the same model grid as for the entire sample (see Tab. 4.1). From each morphological class a median galaxy was chosen and analysed by the pixel-pixel analysis. This was done by using the results from Chapter 3.4.3 to minimise the scatter due to a misclassification. Table 4.5 summarises the galaxies used for this analysis.

Table 4.5: Subsample of different morphological types.

| VCC number | Morphological <br> type |
| :--- | :--- |
| VCC-0030 | dI |
| VCC-0048 | Sd |
| VCC-0641 | BCD |
| VCC-0679 | Sm |
| VCC-0938 | Sc |
| VCC-1188 | Amorphous |
| VCC-1290 | Sb |

As a result of our analysis we obtain spatially resolved maps of stellar mass and SFR density. Some of the preliminary results are shown in Fig. 4.39 for a small subsample of our galaxies. To create these maps we broke down each of the fits files into resolution elements of $5 \times 5$ pixels and fitted the resulting SEDs with GALEV simple stellar populations models, yielding the light-weighted stellar population age, stellar mass and extinction for each element. By comparing the stellar mass maps of the dEs with ones of e.g. dIs, we are able to conclude if both morphological classes have the same stellar mass distribution. In this case it is obvious that the dEs have much higher surface mass densities as the dIs. It is also possible to conclude how the star formation
rate of a dI has to change to fit the mass regime of the dEs.


Figure 4.39: Stellar mass surface density maps of a sample of Virgo galaxies.

Another example is shown in Fig. 4.40 for the Sb galaxy VCC-1290. The left panel displays the age map and the right panel the dust reddening map. Obvious features in these maps are the younger ages of the spiral arms and the old stellar population of the bulge. Furthermore, the bulge region is less influenced by dust reddening than the surrounded disc. Figure 4.41 shows the age maps of a BCD and Sm galaxy. Interestingly, these two galaxies have a quite similar age distribution, making further and more detailed investigations worthwhile.
In a next step the results of the pixel-by-pixel analysis will be extended to the entire sample. Furthermore, the impact of the truncation event on the stellar mass and SF density can be investigated in more detail.
To develop further, it might be also possible in the future to refine the truncation and RPS/starvation physics to account for the radius dependence of these processes on disc galaxies.

### 4.6.2 New observation

For the pixel-by-pixel analysis of the galaxy using GALEV/GAZELLE, the resolution of the detector and the quality of the observational conditions (seeing) at the observatory are one of the crucial limitations. Since the resolution of the SDSS is low ( 0.396 arsec/pix) and the median seeing is about 1.6 arcsec, new observations of Virgo galaxies were performed by us (H.T. Meyer and R. Kotulla), using the 3.5 m WIYN


Figure 4.40: Pixel-by-pixel analysis of VCC-1290. Left: age map; Right: dust reddening map.


Figure 4.41: Pixel-by-pixel analysis of the BCD VCC-1290 (left) and the Sm galaxy VCC-0679 (right). Shown are the best fitting ages of the galaxies.
telescope ${ }^{10}$ at the Kitt Peak National Observatory (KPNO, USA). The targets of the observing run are summarised in Tab 4.6. In the first run the galaxies were observed for 2 nights using the WIYN Mini-Mosaic Imager (MIMO), which has a resolution

[^28]of $0.14 \mathrm{arsec} / \mathrm{pixel}$ and a field of view of $9.6 \times 9.6$ arcmin. The used filters for the first observing run were the Harris UBR filters (central wavelenghts: $3640 \AA$; $4304 \AA$ A ; 6394 A) ${ }^{11}$.

The second run was performed with the WIYN High Resolution Infrared Camera (WHIRC) and was scheduled for one night. WHIRC has a resolution of $0.1 \mathrm{arsec} / \mathrm{pixel}$ and a field of view of $3.3 \times 3.3 \mathrm{arcmin}$. In this run the galaxies have been observed in the infrared, using the J filter $(1.250 \mu \mathrm{~m})$, H filter $(1.651 \mu \mathrm{~m})$, and $\mathrm{K}_{\mathrm{s}}$ filter (2.168 $\mu \mathrm{m})^{12}$. Owing to the fact that the observation was performed in H and with the superb observing conditions, an average seeing of 1.0 arcsec and below would be possible. However, due to variations in the weather conditions, this seeing was unfortunately not achieved for the entire sample of our observation.
The combination of the resolutions of the instruments and the superb observational conditions at Kitt Peak allows a detailed study of a set of Virgo galaxies, which will be published in an upcoming paper together with R. Kotulla. At the time of writing the data of these observing runs are partly reduced and are applied to first tests with the GALEV/GAZELLE-code.

Table 4.6: Subsample of different morphological types that was observed with the WIYN telescope. Shown are the used filters and the exposure times in each filter.

| VCC number | Observed filter | Exposure time in UBRJHK [sec] |
| :---: | :---: | :---: |
| VCC-0001 | U,B,R | $1200 / 600 / 600 /-/-/-$ |
| VCC-0087 | H | $-/-/-/-/ 2000 /-$ |
| VCC-0226 | U,B,R,H,K | $1200 / 600 / 600 /-/ 2000 / 4320$ |
| VCC-0241 | J,K | $-/-/-/ 80 /-/ 2640$ |
| VCC-0849 | U,B,R,K | $1200 / 600 / 600 /-/-/ 2720$ |
| VCC-1488 | U,B,R,J,K | $1200 / 600 / 600 / 2560 /-/ 2320$ |
| VCC-1507 | U,B,R,J,K | $1200 / 600 / 600 / 480 /-/ 2560$ |
| VCC-1566 | U,B,R,J,K | $1200 / 600 / 600 / 660 /-/ 2560$ |
| VCC-1789 | U,B,R,J,K | $1200 / 600 / 600 / 60 /-/ 2560$ |
| NGC-4565 | J,K | $-/-/-/ 3255 /-/ 80$ |

[^29]
## Chapter 5 Conclusions and outlook

### 5.1 The possible progenitors of early-type dwarf galaxies

In Chapter 2 the structural properties of the Blue Compact Dwarf Galaxies (BCDs) in the Virgo cluster were analysed by a special algorithm. This algorithm accounts for possible irregularities and multiple star forming regions of a BCDs and enables us to investigate the underlying old stellar population through the creation of surface brightness brightness profiles (SBPs). The SBPs were decomposed into the contribution of the starburst and old stellar component, which is also called low surface brightness (LSB-) component or host galaxy. With this decomposition algorithm we are in the excellent position to analyse only the parameters of the underlying LSB-component, which contributes mainly to the galaxy mass. It turns out that the parameters of the LSB-components of the BCDs and the ones of compact Dwarf Ellipticals (dEs) are in good agreement. Regarding the used parameter-space $\left(\mathrm{M}_{\mathrm{r}}-(\mathrm{g}-\mathrm{i})-\mathrm{R}_{\mathrm{eff}}-\langle\mu\rangle_{\mathrm{eff}}\right)$ of this thesis, we concluded that BCDs are the possible progenitors of compact dEs in the Virgo cluster.
Apart from the BCDs, also all other late-type galaxies of the Virgo cluster were analysed in the course of Chapter 3. The parameters of the different morphological types of the Virgo Cluster Catalog (VCC) were reviewed and if necessary reclassified. As expected, the parameters of the late-type galaxies show significant differences to the dEs. However, the morphological class of the Irregular Galaxies (dI) and the late spirals of type Sd and Sm show some similarity to the dEs according to the parameterspace.
These differences between late- and early-type galaxies are not surprisingly, but the most interesting question is how the late-type galaxies will look like in several billion years? To answer this question, evolutionary synthesis models of GALEV and GAZELLE are used, which model the evolution of galaxies based on a set of input parameters, like a star formation history and our observed magnitudes of Chapter 3.

Since the late-type galaxies are members of the Virgo cluster, it is crucial to take into account the influence of the dense cluster environment, which is able to remove the gas from a galaxy and therefore halting the star formation on very short time-scales (a so called truncation).
In Chapter 4 of this thesis, we found that also the Sd and Sm galaxies, after two Gyr of a truncated evolution, are located in the same parameter-space as the dEs of the Virgo cluster. The dIs, which cover almost the same parameter-space of today's dEs, will evolve into very faint galaxies with a very low surface brightness. Therefore, these future dIs will not fall into the locus of today's dEs and may share the structural properties of dSph of the Local Group.
The findings of Chapter 4 point to evolutionary connections between the late spirals ( Sd and Sm ) and the dEs.
In summary, we found that future BCDs and late spirals ( Sd and $\mathrm{Sm} \mathrm{)}$, formation has ceased, will share the same parameter-space like Virgo dEs. Therefore, the results of this thesis supports the idea that these galaxy types will be transformed by the environment of the Virgo cluster, resulting in objects quite similar to today's early-type dwarf galaxies.

### 5.2 Outlook

As seen in the course of this thesis, there are several galaxies within the VCC, which are not or only partly covered by the SDSS. With the new data releases of the SDSS, it will be possible to increase the coverage of this sample of Virgo galaxies. However, due to the overlap with other galaxies (and stars) in the fore- or background, we will never reached an $100 \%$ coverage of the Virgo galaxies within a certain completeness limit.
As always in observational astronomy, it would be a great pleasure to have much more deeper observations of the Virgo galaxies. This will be achieved by the Next Generation Virgo Cluster Survey (NGVS; Ferrarese et al., 2012). The results of the NGVS will be in particular interest with respect to the low-luminosity dwarf irregular galaxies, with their low $\mathrm{S} / \mathrm{N}$. Furthermore, the NGVS will also shed more light on the faint underlying LSB-components of BCDs and may reveal new features in their structure.
In respect to the BCDs, it would be also very interesting to investigate in more details the differences of structure of BCDs in clusters and in the field. Also the comparison of

BCDs in different clusters, especially clusters at different stages of virialisation, would be of great interest. Of course, this is true for all kind of late-type galaxies.
To check the results of our GALEV/GAZELLE analysis, it would also be useful to compare our results with other evolutionary synthesis models and codes (Hansson, 2012; Hansson et al., 2012). Also the inclusion of spatially resolved maps, e.g. of the stellar mass density and star formation, and the comparison of the different morphological types is of major interest. A first step to these spatially resolved maps was already taken in Chapter 4 and we are excited to see the results of our entire Virgo sample.

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## Bibliography

Abadi, M. G., Moore, B., and Bower, R. G.: 1999, MNRAS 308, 947
Abell, G. O., Corwin, Jr., H. G., and Olowin, R. P.: 1989, ApJS 70, 1
Adelman-McCarthy, J. K., Agüeros, M. A., Allam, S. S., Anderson, K. S. J., Anderson, S. F., Annis, J., Bahcall, N. A., Bailer-Jones, C. A. L., and et al.: 2007, ApJS 172, 634

Aloisi, A., Clementini, G., Tosi, M., Annibali, F., Contreras, R., Fiorentino, G., Mack, J., Marconi, M., Musella, I., Saha, A., Sirianni, M., and van der Marel, R. P.: 2007, ApJ 667, L151

Amorín, R., Alfonso, J., Aguerri, J. A. L., Muñoz-Tuñón, C., and Cairós, L. M.: 2009, AGA 501, 75

Anders, P. and Fritze-v. Alvensleben, U.: 2003, A\&A 401, 1063
Athanassoula, E.: 2005, MNRAS 358, 1477
Bekki, K.: 2008, MNRAS 388, L10
Bertelli, G., Bressan, A., Chiosi, C., Fagotto, F., and Nasi, E.: 1994, A\&AS 106, 275
Bertin, E. and Arnouts, S.: 1996, AEGAS 117, 393
Bicker, J., Fritze-v. Alvensleben, U., and Fricke, K. J.: 2002, A\&A 387, 412
Binggeli, B. and Cameron, L. M.: 1991, AEA 252, 27
Binggeli, B., Popescu, C. C., and Tammann, G. A.: 1993, AEAAS 98, 275
Binggeli, B., Sandage, A., and Tammann, G. A.: 1985, AJ 90, 1681
Binggeli, B., Tammann, G. A., and Sandage, A.: 1987, AJ 94, 251

Binggeli, B., Tarenghi, M., and Sandage, A.: 1990, A\&A 228, 42
Binney, J. and Tremaine, S.: 2008, Galactic Dynamics: Second Edition, Princeton University Press

Blakeslee, J. P., Jordán, A., Mei, S., Côté, P., Ferrarese, L., Infante, L., Peng, E. W., Tonry, J. L., and West, M. J.: 2009, ApJ 694, 556

Blanton, M. R., Dalcanton, J., Eisenstein, D., Loveday, J., Strauss, M. A., SubbaRao, M., Weinberg, D. H., Anderson, Jr., J. E., and et al.: 2001, AJ 121, 2358

Bohlin, R. C. and Gilliland, R. L.: 2004, AJ 128, 3053
Boselli, A., Boissier, S., Cortese, L., and Gavazzi, G.: 2008, ApJ 674, 742
Boselli, A. and Gavazzi, G.: 2006, PASP 118, 517
Bothun, G. D., Mould, J. R., Caldwell, N., and MacGillivray, H. T.: 1986, AJ 92, 1007
Boylan-Kolchin, M., Springel, V., White, S. D. M., Jenkins, A., and Lemson, G.: 2009, MNRAS 398, 1150

Bruzual, G. and Charlot, S.: 2003, MNRAS 344, 1000
Cairós, L. M., Caon, N., Papaderos, P., Noeske, K., Vílchez, J. M., García Lorenzo, B., and Muñoz-Tuñón, C.: 2003, ApJ 593, 312

Cairós, L. M., Caon, N., Vílchez, J. M., González-Pérez, J. N., and Muñoz-Tuñón, C.: 2001a, ApJS 136, 393

Cairós, L. M., Vílchez, J. M., González Pérez, J. N., Iglesias-Páramo, J., and Caon, N.: 2001b, ApJS 133, 321

Calzetti, D., Armus, L., Bohlin, R. C., Kinney, A. L., Koornneef, J., and StorchiBergmann, T.: 2000, ApJ 533, 682

Cardelli, J. A., Clayton, G. C., and Mathis, J. S.: 1989, ApJ 345, 245
Carroll, B. W. and Ostlie, D. A.: 2006, An introduction to modern astrophysics and cosmology

Chabrier, G.: 2003, PASP 115, 763

Chilingarian, I. V.: 2009, MNRAS 394, 1229
Chung, A., van Gorkom, J. H., Kenney, J. D. P., Crowl, H., and Vollmer, B.: 2009, AJ 138, 1741

Chung, A., van Gorkom, J. H., Kenney, J. D. P., and Vollmer, B.: 2007, ApJ 659, L115
Clemens, M. S., Panuzzo, P., Rampazzo, R., Vega, O., and Bressan, A.: 2011, MNRAS 412, 2063

Conroy, C. and Gunn, J. E.: 2010, ApJ 712, 833
Conroy, C., Gunn, J. E., and White, M.: 2009, ApJ 699, 486
Conroy, C., White, M., and Gunn, J. E.: 2010, ApJ 708, 58
Conselice, C. J., Gallagher, III, J. S., and Wyse, R. F. G.: 2001, ApJ 559, 791
Conselice, C. J., O’Neil, K., Gallagher, J. S., and Wyse, R. F. G.: 2003, ApJ 591, 167
Côté, P., Blakeslee, J. P., Ferrarese, L., Jordán, A., Mei, S., Merritt, D., Milosavljević, M., Peng, E. W., Tonry, J. L., and West, M. J.: 2004, ApJS 153, 223

Cunow, B.: 1999, Astrophysics and Space Science Supplement 269, 621
Davies, J. I. and Phillipps, S.: 1988, MNRAS 233, 553
De Lucia, G. and Blaizot, J.: 2007, MNRAS 375, 2
de Rijcke, S., Michielsen, D., Dejonghe, H., Zeilinger, W. W., and Hau, G. K. T.: 2005, A\&A 438, 491
de Vaucouleurs, G.: 1961, ApJS 6, 213
de Vaucouleurs, G.: 1974, in J. R. Shakeshaft (ed.), The Formation and Dynamics of Galaxies, Vol. 58 of IAU Symposium, p. 335

Dekel, A. and Silk, J.: 1986, ApJ 303, 39
Dellenbusch, K. E., Gallagher, III, J. S., Knezek, P. M., and Noble, A. G.: 2008, AJ 135, 326

Djorgovski, S. and Davis, M.: 1987, ApJ 313, 59

Dressler, A.: 1980, ApJ 236, 351

Dressler, A., Lynden-Bell, D., Burstein, D., Davies, R. L., Faber, S. M., Terlevich, R., and Wegner, G.: 1987, ApJ 313, 42

Drinkwater, M. and Hardy, E.: 1991, AJ 101, 94

Drinkwater, M. J., Currie, M. J., Young, C. K., Hardy, E., and Yearsley, J. M.: 1996, MNRAS 279, 595

Elmegreen, B. G. and Elmegreen, D. M.: 2010, ApJ 722, 1895
Falkenberg, M. A., Kotulla, R., and Fritze, U.: 2009, MNRAS 397, 1940

Farouki, R. and Shapiro, S. L.: 1981, ApJ 243, 32

Ferguson, H. C. and Binggeli, B.: 1994, $A \mathcal{E} A R$ 6, 67

Ferrara, A. and Tolstoy, E.: 2000, MNRAS 313, 291
Ferrarese, L., Côté, P., Cuillandre, J.-C., Gwyn, S. D. J., Peng, E. W., MacArthur, L. A., Duc, P.-A., Boselli, A., and et al.: 2012, ApJS 200, 4

Fioc, M. and Rocca-Volmerange, B.: 1997, AEGA 326, 950
Fritz, J. and Hevics Collaboration: 2011, in W. Wang, J. Lu, Z. Luo, Z. Yang, H. Hua, and Z. Chen (eds.), Galaxy Evolution: Infrared to Millimeter Wavelength Perspective, Vol. 446 of Astronomical Society of the Pacific Conference Series, p. 77

Fujita, Y. and Nagashima, M.: 1999, ApJ 516, 619
Fumagalli, M. and Gavazzi, G.: 2008, AEGA 490, 571
Gallagher, III, J. S., Hunter, D. A., and Tutukov, A. V.: 1984, ApJ 284, 544

Gardner, J. P., Mather, J. C., Clampin, M., Doyon, R., Greenhouse, M. A., Hammel, H. B., Hutchings, J. B., Jakobsen, P., Lilly, S. J., Long, K. S., Lunine, J. I., McCaughrean, M. J., Mountain, M., Nella, J., Rieke, G. H., Rieke, M. J., Rix, H.W., Smith, E. P., Sonneborn, G., Stiavelli, M., Stockman, H. S., Windhorst, R. A., and Wright, G. S.: 2006, Space Sci. Rev. 123, 485

Gavazzi, G., Bonfanti, C., Sanvito, G., Boselli, A., and Scodeggio, M.: 2002, ApJ 576, 135

Gavazzi, G., Boselli, A., Donati, A., Franzetti, P., and Scodeggio, M.: 2003, A\&A 400, 451

Gavazzi, G., Boselli, A., Mayer, L., Iglesias-Paramo, J., Vílchez, J. M., and Carrasco, L.: 2001, ApJ 563, L23

Gavazzi, G., Boselli, A., Scodeggio, M., Pierini, D., and Belsole, E.: 1999, MNRAS 304, 595

Gavazzi, G., Contursi, A., Carrasco, L., Boselli, A., Kennicutt, R., Scodeggio, M., and Jaffe, W.: 1995, A\&A 304, 325

Gavazzi, G., O’Neil, K., Boselli, A., and van Driel, W.: 2006, A\&GA 449, 929
Geha, M., Guhathakurta, P., and van der Marel, R. P.: 2003, AJ 126, 1794
Gil de Paz, A. and Madore, B. F.: 2005, ApJS 156, 345

González-García, A. C., Aguerri, J. A. L., and Balcells, M.: 2005, AEAA 444, 803
Goto, T., Yamauchi, C., Fujita, Y., Okamura, S., Sekiguchi, M., Smail, I., Bernardi, M., and Gomez, P. L.: 2003, MNRAS 346, 601

Graham, A. W. and Driver, S. P.: 2005, Publications of the Astron. Soc. of Australia 22, 118

Graham, A. W. and Worley, C. C.: 2008, MNRAS 388, 1708
Grebel, E. K., Gallagher, III, J. S., and Harbeck, D.: 2003, AJ 125, 1926
Green, S. F. and Jones, M. H.: 2004, An Introduction to the Sun and Stars

Gu, Q., Zhao, Y., Shi, L., Peng, Z., and Luo, X.: 2006, AJ 131, 806
Gunn, J. E., Carr, M., Rockosi, C., Sekiguchi, M., Berry, K., Elms, B., de Haas, E., Ivezić, Ž., and et al.: 1998, AJ 116, 3040

Gunn, J. E. and Gott, III, J. R.: 1972, ApJ 176, 1
Gunn, J. E., Siegmund, W. A., Mannery, E. J., Owen, R. E., Hull, C. L., Leger, R. F., Carey, L. N., Knapp, G. R., and et al.: 2006, AJ 131, 2332

Guo, Q., White, S., Boylan-Kolchin, M., De Lucia, G., Kauffmann, G., Lemson, G., Li, C., Springel, V., and Weinmann, S.: 2011, MNRAS 413, 101

Guseva, N. G., Papaderos, P., Izotov, Y. I., Green, R. F., Fricke, K. J., Thuan, T. X., and Noeske, K. G.: 2003, A\&A 407, 105

Hallenbeck, G., Papastergis, E., Huang, S., Haynes, M. P., Giovanelli, R., Boselli, A., Boissier, S., Heinis, S., Cortese, L., and Fabello, S.: 2012, ArXiv e-prints

Hansson, K. A.: 2012, Ph.D. thesis, Heidelberg
Hansson, K. S. A., Lisker, T., and Grebel, E. K.: 2012, ArXiv e-prints
Heckman, T. M., Sembach, K. R., Meurer, G. R., Strickland, D. K., Martin, C. L., Calzetti, D., and Leitherer, C.: 2001, ApJ 554, 1021

Hensler, G.: 2012, Morphological Mutations of Dwarf Galaxies, p. 75
Hubble, E. P.: 1922, ApJ 56, 162
Hubble, E. P.: 1926, ApJ 64, 321
Hubble, E. P.: 1929, ApJ 69, 103
Hunt, L. K., Thuan, T. X., and Izotov, Y. I.: 2003, ApJ 588, 281
Izotov, Y. I., Chaffee, F. H., Foltz, C. B., Green, R. F., Guseva, N. G., and Thuan, T. X.: 1999a, ApJ 527, 757

Izotov, Y. I., Dyak, A. B., Chaffee, F. H., Foltz, C. B., Kniazev, A. Y., and Lipovetsky, V. A.: 1996, ApJ 458, 524

Izotov, Y. I., Papaderos, P., Thuan, T. X., Fricke, K. J., Foltz, C. B., and Guseva, N. G.: 1999b, ArXiv Astrophysics e-prints

Janz, J., Laurikainen, E., Lisker, T., Salo, H., Peletier, R. F., Niemi, S.-M., den Brok, M., Toloba, E., Falcón-Barroso, J., Boselli, A., and Hensler, G.: 2012, ApJ 745, L24

Janz, J. and Lisker, T.: 2008, ApJ 689, L25
Janz, J. and Lisker, T.: 2009, ApJ 696, L102
Jedrzejewski, R. I.: 1987, MNRAS 226, 747

Jerjen, H.: 2012, Dwarf Elliptical Galaxies: United and Divided, p. 133
Jerjen, H., Kalnajs, A., and Binggeli, B.: 2000, AEA 358, 845
Jester, S., Schneider, D. P., Richards, G. T., Green, R. F., Schmidt, M., Hall, P. B., Strauss, M. A., Vanden Berk, D. E., Stoughton, C., Gunn, J. E., Brinkmann, J., Kent, S. M., Smith, J. A., Tucker, D. L., and Yanny, B.: 2005, AJ 130, 873

Jones, M. H. and Lambourne, R. J. A.: 2004, An Introduction to Galaxies and Cosmology

Kenney, J. D. P., van Gorkom, J. H., and Vollmer, B.: 2004, $A J$ 127, 3361
Kennicutt, Jr., R. C.: 1998, ApJ 498, 541
Kissler-Patig, M., Küpcü Yoldaş, A., and Liske, J.: 2009, The Messenger 138, 11
Kniazev, A. Y., Pustilnik, S. A., Grebel, E. K., Lee, H., and Pramskij, A. G.: 2004, ApJS 153, 429

Koleva, M., Prugniel, P., de Rijcke, S., and Zeilinger, W. W.: 2011, MNRAS 417, 1643
Koopmann, R. A. and Kenney, J. D. P.: 2004, ApJ 613, 851
Kormendy, J.: 1985, ApJ 295, 73
Kotulla, R., Anders, P., Weilbacher, P., and Fritze, U.: 2010, in G. Bruzual \& S. Charlot (ed.), IAU Symposium, Vol. 262 of IAU Symposium, pp 366-367

Kotulla, R., Fritze, U., Weilbacher, P., and Anders, P.: 2009, MNRAS 396, 462
Kroupa, P.: 2001, MNRAS 322, 231
Krüger, H. and Fritze-v. Alvensleben, U.: 1994, A\&AA 284, 793
Krüger, H., Fritze-v. Alvensleben, U., and Loose, H.-H.: 1995, A\&A 303, 41
Krüger, H., Fritze-von Alvensleben, U., Loose, H.-H., and Fricke, K. J.: 1991, A\&A 242, 343

Krywult, J.: 2009, Astronomische Nachrichten 330, 946
Kunth, D. and Östlin, G.: 2000, A\&AR 10, 1

Larson, R. B., Tinsley, B. M., and Caldwell, C. N.: 1980, ApJ 237, 692
Leitherer, C., Schaerer, D., Goldader, J. D., Delgado, R. M. G., Robert, C., Kune, D. F., de Mello, D. F., Devost, D., and Heckman, T. M.: 1999, ApJS 123, 3

Lejeune, T., Cuisinier, F., and Buser, R.: 1997, $A \mathcal{E} A S$ 125, 229
Lejeune, T., Cuisinier, F., and Buser, R.: 1998, $A \mathcal{E} A S$ 130, 65
Lelli, F., Verheijen, M., Fraternali, F., and Sancisi, R.: 2012, AEA 537, A72
Lieder, S., Lisker, T., Hilker, M., Misgeld, I., and Durrell, P.: 2012, A\&A 538, A69
Lin, D. N. C. and Faber, S. M.: 1983, ApJ 266, L21
Lisker, T.: 2012, Astronomische Nachrichten 333, 405
Lisker, T., Glatt, K., Westera, P., and Grebel, E. K.: 2006a, AJ 132, 2432
Lisker, T., Grebel, E. K., and Binggeli, B.: 2006b, AJ 132, 497
Lisker, T., Grebel, E. K., and Binggeli, B.: 2008, AJ 135, 380
Lisker, T., Grebel, E. K., Binggeli, B., and Glatt, K.: 2007, ApJ 660, 1186
Loose, H.: 1985, in IUE Proposal, pp 2259-+
Loose, H.-H. and Thuan, T. X.: 1985, in D. Kunth, T. X. Thuan, and J. Tran Thanh van (eds.), Star-Forming Dwarf Galaxies and Related Objects, pp 73-+

Loose, H.-H. and Thuan, T. X.: 1986, ApJ 309, 59
Lotz, J. M., Primack, J., and Madau, P.: 2004, AJ 128, 163
MacArthur, L. A., Courteau, S., Bell, E., and Holtzman, J. A.: 2004, ApJS 152, 175
Magrini, L., Bianchi, S., Corbelli, E., Cortese, L., Hunt, L., Smith, M., Vlahakis, C., Davies, J., Bendo, G. J., Baes, M., Boselli, A., Clemens, M., Casasola, V., de Looze, I., Fritz, J., Giovanardi, C., Grossi, M., Hughes, T., Madden, S., Pappalardo, C., Pohlen, M., di Serego Alighieri, S., and Verstappen, J.: 2011, A\&A 535, A13

Maraston, C.: 2005, MNRAS 362, 799
Marcolini, A., Brighenti, F., and D'Ercole, A.: 2003, MNRAS 345, 1329

Marquart, T., Fathi, K., Östlin, G., Bergvall, N., Cumming, R. J., and Amram, P.: 2007, AGA 474, L9

Mas-Hesse, J. M. and Kunth, D.: 1999, A\&A 349, 765
Mastropietro, C., Moore, B., Mayer, L., Debattista, V. P., Piffaretti, R., and Stadel, J.: 2005, MNRAS 364, 607

Mayer, L., Governato, F., Colpi, M., Moore, B., Quinn, T., Wadsley, J., Stadel, J., and Lake, G.: 2001, ApJ 547, L123

Mayer, L., Mastropietro, C., Wadsley, J., Stadel, J., and Moore, B.: 2006, MNRAS 369, 1021

McGaugh, S. S. and de Blok, W. J. G.: 1997, ApJ 481, 689
McLaughlin, D. E.: 1999, ApJ 512, L9
McQuinn, K. B. W., Skillman, E. D., Cannon, J. M., Dalcanton, J., Dolphin, A., Hidalgo-Rodríguez, S., Holtzman, J., Stark, D., Weisz, D., and Williams, B.: 2010a, ApJ 721, 297

McQuinn, K. B. W., Skillman, E. D., Cannon, J. M., Dalcanton, J., Dolphin, A., Hidalgo-Rodríguez, S., Holtzman, J., Stark, D., Weisz, D., and Williams, B.: 2010b, ApJ 724, 49

Mei, S., Blakeslee, J. P., Côté, P., Tonry, J. L., West, M. J., Ferrarese, L., Jordán, A., Peng, E. W., Anthony, A., and Merritt, D.: 2007, ApJ 655, 144

Meyer, H. T. and Lisker, T.: 2011, in M. Koleva, P. Prugniel, \& I. Vauglin (ed.), EAS Publications Series, Vol. 48 of EAS Publications Series, pp 205-206

Misgeld, I. and Hilker, M.: 2011, MNRAS 414, 3699
Mo, H., van den Bosch, F. C., and White, S.: 2010, Galaxy Formation and Evolution
Moeller, C. S., Fritze-v. Alvensleben, U., and Fricke, K. J.: 1997, AEAA 317, 676
Möllenhoff, C., Popescu, C. C., and Tuffs, R. J.: 2006, AEAA 456, 941
Moore, B., Katz, N., Lake, G., Dressler, A., and Oemler, A.: 1996, Nature 379, 613
Noeske, K. G.: 1999, Master's thesis, Göttingen

Noeske, K. G., Guseva, N. G., Fricke, K. J., Izotov, Y. I., Papaderos, P., and Thuan, T. X.: 2000, A\&A 361, 33

Noeske, K. G., Koo, D. C., Phillips, A. C., Willmer, C. N. A., Melbourne, J., Gil de Paz, A., and Papaderos, P.: 2006, ApJ 640, L143

Noeske, K. G., Papaderos, P., Cairós, L. M., and Fricke, K. J.: 2003, AEGA 410, 481
Oemler, Jr., A.: 1974, ApJ 194, 1
Oke, J. B. and Gunn, J. E.: 1983, ApJ 266, 713
Östlin, G., Amram, P., Masegosa, J., Bergvall, N., and Boulesteix, J.: 1999, A\&AS 137, 419

Papaderos, P. and Östlin, G.: 2011, ArXiv e-prints
Papaderos, P., Fricke, K. J., Thuan, T. X., Izotov, Y. I., and Nicklas, H.: 1999, A\&A 352, L57

Papaderos, P., Guseva, N. G., Izotov, Y. I., and Fricke, K. J.: 2008, AEAA 491, 113
Papaderos, P., Guseva, N. G., Izotov, Y. I., Noeske, K. G., Thuan, T. X., and Fricke, K. J.: 2006, AEA 457, 45

Papaderos, P., Izotov, Y. I., Fricke, K. J., Thuan, T. X., and Guseva, N. G.: 1998, A\&A 338, 43

Papaderos, P., Izotov, Y. I., Thuan, T. X., Noeske, K. G., Fricke, K. J., Guseva, N. G., and Green, R. F.: 2002, AEGA 393, 461

Papaderos, P., Loose, H.-H., Fricke, K. J., and Thuan, T. X.: 1996a, AEAA 314, 59
Papaderos, P., Loose, H.-H., Thuan, T. X., and Fricke, K. J.: 1996b, AEAAS 120, 207
Papaderos, P., Noeske, K. G., Cairós, L. M., Vílchez, J. M., and Fricke, K. J.: 2001, in K. S. de Boer, R.-J. Dettmar, and U. Klein (eds.), Dwarf galaxies and their environment, p. 283

Papaderos, P. and Östlin, G.: 2012, A\&A 537, A126
Pedraz, S., Gorgas, J., Cardiel, N., Sánchez-Blázquez, P., and Guzmán, R.: 2002, MNRAS 332, L59

Peletier, R. F., Kutdemir, E., van der Wolk, G., Falcón-Barroso, J., Bacon, R., Bureau, M., Cappellari, M., Davies, R. L., de Zeeuw, P. T., Emsellem, E., Krajnović, D., Kuntschner, H., McDermid, R. M., Sarzi, M., Scott, N., Shapiro, K. L., van den Bosch, R. C. E., and van de Ven, G.: 2012, MNRAS 419, 2031

Peng, C. Y., Ho, L. C., Impey, C. D., and Rix, H.-W.: 2010, AJ 139, 2097
Petrosian, V.: 1976, ApJ 209, L1
Phillipps, S., Parker, Q. A., Schwartzenberg, J. M., and Jones, J. B.: 1998, ApJ 493, L59

Pier, J. R., Munn, J. A., Hindsley, R. B., Hennessy, G. S., Kent, S. M., Lupton, R. H., and Ivezić, Ž.: 2003, AJ 125, 1559

Piovan, L., Tantalo, R., and Chiosi, C.: 2006a, MNRAS 366, 923
Piovan, L., Tantalo, R., and Chiosi, C.: 2006b, MNRAS 370, 1454
Postman, M. and Geller, M. J.: 1984, ApJ 281, 95
Quilis, V., Moore, B., and Bower, R.: 2000, Science 288, 1617
Recchi, S. and Hensler, G.: 2006, AGA 445, L39
Richer, M., McCall, M. L., and Stasińska, G.: 1998, A\&A 340, 67
Richer, M. G. and McCall, M. L.: 1995, ApJ 445, 642
Rines, K. and Geller, M. J.: 2008, $A J$ 135, 1837
Roediger, E.: 2009, Astronomische Nachrichten 330, 888
Roediger, J. C., Courteau, S., MacArthur, L. A., and McDonald, M.: 2011, MNRAS 416, 1996

Rubin, V. C. and Ford, Jr., W. K.: 1970, ApJ 159, 379
Sabatini, S., Davies, J., Scaramella, R., Smith, R., Baes, M., Linder, S. M., Roberts, S., and Testa, V.: 2003, MNRAS 341, 981

Salpeter, E. E.: 1955, ApJ 121, 161

Salzer, J. J., MacAlpine, G. M., and Boroson, T. A.: 1989, ApJS 70, 479
Salzer, J. J. and Norten, S. A.: 1999, The Low Surface Brightness Universe, IAU Col. 171170

Sandage, A.: 1961, The Hubble atlas of galaxies
Sandage, A.: 1986, AEAA 161, 89
Sandage, A. and Binggeli, B.: 1984, AJ 89, 919
Sandage, A., Binggeli, B., and Tammann, G. A.: 1985, AJ 90, 395
Sandage, A. and Brucato, R.: 1979, AJ 84, 472
Schechter, P.: 1976, ApJ 203, 297
Schlegel, D. J., Finkbeiner, D. P., and Davis, M.: 1998, ApJ 500, 525
Schmidt, M.: 1959, ApJ 129, 243
Searle, L. and Sargent, W. L. W.: 1972, ApJ 173, 25
Sérsic, J. L.: 1968, Atlas de galaxias australes, Cordoba, Argentina: Observatorio Astronomico

Shen, S., Mo, H. J., White, S. D. M., Blanton, M. R., Kauffmann, G., Voges, W., Brinkmann, J., and Csabai, I.: 2003, MNRAS 343, 978

Silich, S. A. and Tenorio-Tagle, G.: 1998, MNRAS 299, 249
Smith, J. A., Tucker, D. L., Kent, S., Richmond, M. W., Fukugita, M., Ichikawa, T., Ichikawa, S.-i., Jorgensen, A. M., Uomoto, A., Gunn, J. E., Hamabe, M., Watanabe, M., Tolea, A., Henden, A., Annis, J., Pier, J. R., McKay, T. A., Brinkmann, J., Chen, B., Holtzman, J., Shimasaku, K., and York, D. G.: 2002, AJ 123, 2121

Smith, R., Davies, J. I., and Nelson, A. H.: 2010, MNRAS 405, 1723
Sparke, L. S. and Gallagher, III, J. S.: 2000, Galaxies in the universe : an introduction, Cambridge University Press

Staveley-Smith, L., Davies, R. D., and Kinman, T. D.: 1992, MNRAS 258, 334

Sung, E.-C., Chun, M.-S., Freeman, K. C., and Chaboyer, B.: 2002, in G. S. Da Costa, E. M. Sadler, \& H. Jerjen (ed.), The Dynamics, Structure \& History of Galaxies: A Workshop in Honour of Professor Ken Freeman, Vol. 273 of Astronomical Society of the Pacific Conference Series, p. 341

Sung, E.-C., Han, C., Ryden, B. S., Chun, M.-S., and Kim, H.-I.: 1998, ApJ 499, 140
Tajiri, Y. Y. and Kamaya, H.: 2002, AEAA 389, 367
Telles, E., Melnick, J., and Terlevich, R.: 1997, MNRAS 288, 78
Thornley, M. D., Schreiber, N. M. F., Lutz, D., Genzel, R., Spoon, H. W. W., Kunze, D., and Sternberg, A.: 2000, ApJ 539, 641

Thuan, T. X.: 1985, ApJ 299, 881
Thuan, T. X.: 1991, in C. Leitherer, N. Walborn, T. Heckman, and C. Norman (eds.), Observations and Models of Blue Compact Dwarf Galaxies, pp 183-+

Thuan, T. X. and Martin, G. E.: 1981, ApJ 247, 823
Tinsley, B. M.: 1968, ApJ 151, 547
Tody, D.: 1993, in R. J. Hanisch, R. J. V. Brissenden, and J. Barnes (eds.), Astronomical Data Analysis Software and Systems II, Vol. 52 of Astronomical Society of the Pacific Conference Series, p. 173

Toloba, E., Boselli, A., Cenarro, A. J., Peletier, R. F., Gorgas, J., Gil de Paz, A., and Muñoz-Mateos, J. C.: 2011, AEA 526, A114

Tonnesen, S., Bryan, G. L., and van Gorkom, J. H.: 2007, ApJ 671, 1434
Tremonti, C. A., Heckman, T. M., Kauffmann, G., Brinchmann, J., Charlot, S., White, S. D. M., Seibert, M., Peng, E. W., Schlegel, D. J., Uomoto, A., Fukugita, M., and Brinkmann, J.: 2004, ApJ 613, 898

Trentham, N. and Hodgkin, S.: 2002, MNRAS 333, 423
Tully, R. B. and Fisher, J. R.: 1977, A\&A 54, 661
Tully, R. B. and Shaya, E. J.: 1984, ApJ 281, 31
Vaduvescu, O., Richer, M. G., and McCall, M. L.: 2006, AJ 131, 1318
van den Hoek, L. B. and Groenewegen, M. A. T.: 1997, AGAS 123, 305
van Gorkom, J. H.: 2004, Clusters of Galaxies: Probes of Cosmological Structure and Galaxy Evolution p. 305
van Zee, L., Salzer, J. J., and Skillman, E. D.: 2001, AJ 122, 121
van Zee, L., Skillman, E. D., and Haynes, M. P.: 2004, AJ 128, 121
van Zee, L., Westpfahl, D., Haynes, M. P., and Salzer, J. J.: 1998, AJ 115, 1000
Vigroux, L., Souviron, J., and Vader, J. P.: 1984, AEA 139, L9
Vilchez, J. M.: 1995, AJ 110, 1090
Vollmer, B.: 2009, AGA 502, 427
Vollmer, B., Cayatte, V., Balkowski, C., and Duschl, W. J.: 2001, ApJ 561, 708
Vollmer, B., Wong, . O. I., Braine, J., Chung, A., and Kenney, J. D. P.: 2012, ArXiv e-prints

Weinmann, S. M., Lisker, T., Guo, Q., Meyer, H. T., and Janz, J.: 2011, ArXiv e-prints
Whitmore, B. C., Gilmore, D. M., and Jones, C.: 1993, ApJ 407, 489
Willman, B., Masjedi, M., Hogg, D. W., Dalcanton, J. J., Martinez-Delgado, D., Blanton, M., West, A. A., Dotter, A., and Chaboyer, B.: 2006, ArXiv Astrophysics e-prints

Woosley, S. E. and Weaver, T. A.: 1995, ApJS 101, 181
Xilouris, E. M., Byun, Y. I., Kylafis, N. D., Paleologou, E. V., and Papamastorakis, J.: 1999, A\&A 344, 868

Yasuda, N., Fukugita, M., Narayanan, V. K., Lupton, R. H., Strateva, I., Strauss, M. A., Ivezić, Ž., Kim, R. S. J., Hogg, D. W., Weinberg, D. H., Shimasaku, K., Loveday, J., Annis, J., Bahcall, N. A., Blanton, M., Brinkmann, J., Brunner, R. J., Connolly, A. J., Csabai, I., Doi, M., Hamabe, M., Ichikawa, S.-I., Ichikawa, T., Johnston, D. E., Knapp, G. R., Kunszt, P. Z., Lamb, D. Q., McKay, T. A., Munn, J. A., Nichol, R. C., Okamura, S., Schneider, D. P., Szokoly, G. P., Vogeley, M. S., Watanabe, M., and York, D. G.: 2001, AJ 122, 1104

York, D. G., Adelman, J., Anderson, Jr., J. E., Anderson, S. F., Annis, J., Bahcall, N. A., Bakken, J. A., Barkhouser, R., and SDSS Collaboration: 2000, AJ 120, 1579 Zwicky, F.: 1937, ApJ 86, 217

Zwicky, F.: 1965, ApJ 142, 1293

## Chapter 6 Appendix

### 6.1 Abbreviations

In Tab. 6.1 the abbreviations of the here presented study are summarised.
Table 6.1: Used abbreviations of this study.

| Abbreviation | Description |
| :---: | :--- |
| BCD | Blue compact dwarf |
| CI | Concentration index |
| CMD | Colour-magnitude diagram |
| dE | Dwarf elliptical |
| $\mathrm{dE}(\mathrm{bc})$ | Dwarf elliptical with blue core |
| dI | Dwarf irregular |
| dSph | Dwarf spheroidal galaxies |
| E | Elliptical galaxy |
| E-ELT European Extremely Large Telescope |  |
| ETG | Early-type galaxy |
| ETDG | Early-type dwarf galaxy |
| EW | Equivalent width |
| FWHM | Full width at half maximum |
| GALEV | GALaxy EVolution code |
| G08 | Graham and Worley (2008) |
| G99 | Gavazzi et al. (1999) |
| ICM | Intra cluster medium |
| IMF | Initial mass function |
| ISM | Inter stellar medium |
| JL09 | Janz and Lisker (2009) |

continued

| Abbreviation | Description |
| :---: | :--- |
| JWST | James Webb Space Telescope |
| LF | Luminosity function |
| LG | Local Group |
| LSB | Low surface brightness |
| NED | NASA/IPAC Extragalactic Database |
| modexp | Modified exponential fitting law |
| NGVS | Next Generation Virgo Cluster Survey |
| P96b | Papaderos et al. (1996b) |
| P08 | Papaderos et al. (2008) |
| pst | Point spread function |
| RPS | Ram pressure stripping |
| SAM | Semi analytical model |
| SB | Starburst |
| SBP | Surface brightness profile |
| SDSS | Sloan Digital Sky Survey |
| SED | Spectral energy distribution |
| SF | Star formation |
| SFDG | Star forming dwarf galaxy |
| SFH | Star formation history |
| SFR | Star formation rate |
| sSFR | Specific star formation rate |
| S/N | Signal to noise ratio |
| SNe | Supernovae |
| VCC | Virgo Cluster Catalog |
| Z | Metallicity |
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Abbreviations

### 6.2 Structural parameters of the sample

In Chapter 3 the structural properties of the late-type galaxies were derived. Table 6.2 summarised the results of this analysis.

Table 6.2: Derived structural parameters of the late-type galaxies.

| VCC | ms | Type | $(\mathrm{g}-\mathrm{i})$ | $\mathrm{M}_{\mathrm{r}}$ | $\mathrm{R}_{\text {eff,r }, \mathrm{kpc}]}$ | $\langle\mu\rangle_{\text {eff,r }}$ | $\mathrm{b} / \mathrm{a}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $[1]$ | $[2]$ | $[3]$ | $[4]$ | $[5]$ | $[6]$ | $[7]$ | $[8]$ |
| 0015 | 1 | Sa? | 0.79 | 14.67 | 1.26 | 22.65 | 0.73 |
| 0020 | 1 | $?$ | 0.73 | 17.38 | 0.45 | 23.14 | 0.85 |
| 0025 | 1 | Sc(r)I. 4 | 0.59 | 11.92 | 1.25 | 19.89 | 0.87 |
| 0031 | 1 | $?$ | 0.56 | 15.20 | 0.38 | 20.61 | 0.47 |
| 0034 | 2 | Sc: | 0.56 | 14.36 | 1.18 | 22.21 | 0.61 |
| 0048 | 2 | Sd(s) / SmIII | 0.61 | 14.03 | 1.84 | 22.85 | 0.87 |
| 0058 | 1 | SBb(r)I-II | 0.73 | 12.50 | 2.18 | 21.68 | 0.51 |
| 0066 | 2 | SBc(s)II | 0.65 | 11.17 | 3.72 | 21.52 | 0.29 |
| 0067 | 2 | Sc(s)pec | 0.53 | 13.51 | 1.59 | 22.02 | 0.36 |
| 0073 | 1 | Sb: | 1.10 | 12.20 | 0.76 | 19.09 | 0.40 |
| 0081 | 2 | d:Sc | 0.58 | 14.43 | 1.77 | 23.16 | 0.87 |
| 0087 | 2 | SmIII | 0.55 | 14.29 | 1.05 | 21.88 | 0.50 |
| 0089 | 1 | SBc(sr)II.2 | 0.78 | 11.45 | 2.49 | 20.92 | 0.79 |
| 0092 | 2 | SbII: | 1.05 | 9.71 | 4.34 | 20.39 | 0.28 |
| 0094 | 1 | S0/a | 1.11 | 12.49 | 0.73 | 19.30 | 0.72 |
| 0097 | 1 | Sc(s)II | 1.00 | 12.05 | 1.49 | 20.41 | 0.44 |
| 0099 | 1 | Sa? | 0.97 | 14.36 | 0.64 | 20.90 | 0.35 |
| 0104 | 2 | dE3 or ImV | 0.67 | 17.04 | 1.51 | 25.43 | 0.73 |
| 0105 | 2 | SBdIV | 0.64 | 13.14 | 2.94 | 22.98 | 0.87 |
| 0113 | 1 | $?$ | 0.37 | 16.61 | 0.82 | 23.68 | 0.53 |
| 0119 | 2 | Sc | 0.47 | 14.85 | 1.24 | 22.82 | 0.31 |
| 0120 | 2 | Scd(on-edge) | 0.67 | 12.58 | 1.30 | 20.64 | 0.26 |
| 0122 | 1 | S014 | 1.05 | 12.75 | 1.04 | 20.33 | 0.60 |
| 0126 | 2 | SBd | 0.56 | 14.01 | 2.19 | 23.20 | 0.67 |
| 0131 | 1 | Sc | 0.94 | 13.45 | 0.88 | 20.66 | 0.20 |
| 0135 | 2 | Spec / BCD | 0.98 | 13.75 | 0.67 | 20.36 | 0.51 |
| 0137 | 1 | $?$ | 0.88 | 16.19 | 0.54 | 22.34 | 0.68 |
| 0143 | 2 | Sc(on-edge): | 0.74 | 14.78 | 0.68 | 21.44 | 0.39 |
| 0145 | 2 | Sc(s) | 0.86 | 12.02 | 2.03 | 21.05 | 0.17 |
|  |  |  |  |  |  |  |  |

continued

| VCC | ms | Type | $(\mathrm{g}-\mathrm{i})$ | $\mathrm{M}_{\mathrm{r}}$ | $\mathrm{R}_{\text {eff,r }}[\mathrm{kpc}]$ | $\langle\mu\rangle_{\text {eff,r }}$ | $\mathrm{b} / \mathrm{a}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $[1]$ | $[2]$ | $[3]$ | $[4]$ | $[5]$ | $[6]$ | $[7]$ | $[8]$ |
| 0152 | 2 | Scd(on-edge) | 1.09 | 12.35 | 1.14 | 20.13 | 0.49 |
| 0157 | 2 | Sc(s)II-III | 0.89 | 10.82 | 2.33 | 20.15 | 0.57 |
| 0162 | 2 | Sd(on-edge) | 0.50 | 14.19 | 1.06 | 21.81 | 0.14 |
| 0166 | 1 | S029 | 1.10 | 11.89 | 0.74 | 18.74 | 0.39 |
| 0167 | 2 | Sb(s) | 1.29 | 9.69 | 2.47 | 19.15 | 0.25 |
| 0187 | 2 | Scd(on-edge $)$ | 0.96 | 12.87 | 1.42 | 21.13 | 0.21 |
| 0199 | 1 | Sa | 1.23 | 11.46 | 1.59 | 19.96 | 0.42 |
| 0213 | 2 | dS? / BCD? | 0.80 | 13.59 | 0.61 | 20.01 | 0.86 |
| 0221 | 1 | SBcIII.4 | 0.65 | 12.56 | 1.45 | 20.86 | 0.85 |
| 0222 | 1 | Sa | 1.19 | 11.32 | 1.50 | 19.69 | 0.28 |
| 0226 | 2 | Sc(r)II.8 | 1.01 | 11.35 | 1.67 | 19.95 | 0.62 |
| 0234 | 1 | Sa | 1.14 | 11.68 | 1.90 | 20.56 | 0.49 |
| 0237 | 1 | $?$ | 0.38 | 16.56 | 0.65 | 23.13 | 0.73 |
| 0241 | 2 | Sd(on-edge) | 0.21 | 14.19 | 1.29 | 22.23 | 0.30 |
| 0267 | 2 | SBbc(s)I-II | 0.62 | 12.92 | 2.47 | 22.38 | 0.84 |
| 0289 | 2 | Sc(on-edge) | 0.55 | 13.99 | 0.74 | 20.84 | 0.31 |
| 0307 | 2 | Sc(s)I.3 | 0.74 | 9.45 | 4.51 | 20.22 | 0.72 |
| 0312 | 1 | S016 | 1.16 | 12.35 | 0.58 | 18.65 | 0.47 |
| 0318 | 2 | SBcd(s)III | 0.46 | 13.74 | 1.67 | 22.35 | 0.56 |
| 0323 | 1 | Sa | 1.12 | 13.52 | 0.61 | 19.93 | 0.31 |
| 0341 | 1 | SBa(s) | 1.15 | 11.35 | 1.75 | 20.05 | 0.40 |
| 0342 | 1 | S017 | 1.11 | 13.36 | 0.40 | 18.85 | 0.44 |
| 0343 | 1 | SBd(s)II | 0.69 | 14.73 | 1.06 | 22.34 | 0.85 |
| 0358 | 1 | SBa(s) | 1.09 | 12.53 | 0.80 | 19.55 | 0.91 |
| 0366 | 1 | S016 | 1.14 | 13.19 | 1.11 | 20.90 | 0.58 |
| 0371 | 1 | S026 | 1.17 | 12.42 | 0.70 | 19.13 | 0.41 |
| 0373 | 1 | S012 | 1.12 | 12.26 | 1.12 | 20.01 | 0.74 |
| 0375 | 1 | S016 | 1.09 | 11.85 | 0.94 | 19.22 | 0.39 |
| 0382 | 1 | SBc(s)II | 0.67 | 11.52 | 1.50 | 19.89 | 0.63 |
| 0386 | 1 | SBa | 1.08 | 13.13 | 1.12 | 20.87 | 0.76 |
| 0393 | 1 | Sc(s)II | 0.72 | 12.71 | 1.72 | 21.39 | 0.80 |
| 0404 | 1 | Scd(on-edge) | 1.04 | 13.96 | 0.87 | 21.14 | 0.18 |
|  |  |  |  |  |  |  |  |

continued

| VCC | ms | Type | $(\mathrm{g}-\mathrm{i})$ | $\mathrm{M}_{\mathrm{r}}$ | $\mathrm{R}_{\text {eff, },[\mathrm{kpc}]}$ | $\langle\mu\rangle_{\text {eff }, \mathrm{r}}$ | $\mathrm{b} / \mathrm{a}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $[1]$ | $[2]$ | $[3]$ | $[4]$ | $[5]$ | $[6]$ | $[7]$ | $[8]$ |
| 0408 | 1 | S036 | 1.19 | 10.94 | 1.33 | 19.05 | 0.46 |
| 0414 | 2 | dE2 or ImV | 0.86 | 16.68 | 0.70 | 23.41 | 0.62 |
| 0415 | 1 | Sd: | 0.77 | 14.53 | 1.14 | 22.31 | 0.56 |
| 0429 | 1 | $?$ | 0.65 | 15.99 | 0.73 | 22.81 | 0.53 |
| 0446 | 2 | Im / BCD: | 0.75 | 15.15 | 0.87 | 22.34 | 0.58 |
| 0449 | 1 | Sbc(on-edge) | 1.24 | 13.34 | 0.93 | 20.66 | 0.37 |
| 0450 | 1 | SOpec | 0.97 | 14.76 | 0.78 | 21.70 | 0.47 |
| 0453 | 2 | Sm | 0.71 | 15.73 | 0.49 | 21.69 | 0.34 |
| 0460 | 2 | Sapec | 1.06 | 10.11 | 3.78 | 20.49 | 0.44 |
| 0465 | 2 | SBc(s)II-III | 0.54 | 11.82 | 1.86 | 20.66 | 0.37 |
| 0472 | 2 | ImIV-V or dE4 | 1.04 | 16.42 | 0.91 | 23.71 | 0.52 |
| 0483 | 2 | Sc(s)III | 0.00 | 0.00 | 0.00 | -inf | 0.00 |
| 0491 | 2 | Scd(s)III | 0.40 | 12.16 | 1.81 | 20.94 | 0.97 |
| 0497 | 2 | Sc(on-edge) | 0.00 | 0.00 | 0.00 | -inf | 0.00 |
| 0509 | 1 | Sd or SmIV | 0.58 | 14.80 | 1.27 | 22.82 | 0.43 |
| 0512 | 2 | SBmIV | 0.63 | 15.13 | 1.52 | 23.53 | 0.39 |
| 0514 | 2 | Sc(s)pec: | 0.77 | 14.12 | 1.77 | 22.86 | 0.90 |
| 0522 | 2 | Sa | 0.99 | 12.22 | 1.79 | 20.98 | 0.49 |
| 0524 | 2 | Sbc(on-edge) | 1.21 | 11.67 | 1.58 | 20.16 | 0.21 |
| 0534 | 1 | SBapec | 1.10 | 12.49 | 1.62 | 21.03 | 0.48 |
| 0559 | 2 | Sab | 1.01 | 11.49 | 2.04 | 20.53 | 0.26 |
| 0566 | 1 | SBmIII: | 0.24 | 15.97 | 0.59 | 22.31 | 0.54 |
| 0567 | 1 | ScdIII | 0.76 | 13.51 | 1.30 | 21.57 | 0.32 |
| 0570 | 2 | Sab | 1.12 | 11.29 | 2.18 | 20.47 | 0.55 |
| 0574 | 2 | dE3 or ImV | 0.92 | 16.48 | 1.22 | 24.41 | 0.78 |
| 0576 | 2 | Sbc(on-edge) | 1.21 | 12.46 | 1.11 | 20.17 | 0.27 |
| 0593 | 1 | S | 0.69 | 14.43 | 0.74 | 21.28 | 0.30 |
| 0596 | 2 | Sc(s)I | 0.87 | 9.02 | 7.62 | 20.92 | 0.87 |
| 0620 | 2 | SmIII | 0.52 | 15.36 | 0.63 | 21.84 | 0.37 |
| 0630 | 2 | Sd(on-edge) | 0.89 | 12.23 | 2.00 | 21.23 | 0.14 |
| 0654 | 2 | RSB025 | 1.13 | 10.91 | 2.28 | 20.19 | 0.70 |
| 0656 | 1 | Sb | 1.21 | 11.99 | 1.13 | 19.74 | 0.61 |
|  |  |  |  |  |  |  |  |

continued

| VCC | ms | Type | $(\mathrm{g}-\mathrm{i})$ | $\mathrm{M}_{\mathrm{r}}$ | $\mathrm{R}_{\text {eff,r }}[\mathrm{kpc}]$ | $\langle\mu\rangle_{\text {eff,r }}$ | $\mathrm{b} / \mathrm{a}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $[1]$ | $[2]$ | $[3]$ | $[4]$ | $[5]$ | $[6]$ | $[7]$ | $[8]$ |
| 0657 | 1 | S017 | 1.19 | 12.10 | 0.27 | 16.76 | 0.46 |
| 0664 | 2 | ScIII-IV | 0.38 | 13.10 | 2.01 | 22.10 | 0.64 |
| 0667 | 1 | Sc(s) | 0.89 | 13.20 | 1.60 | 21.71 | 0.48 |
| 0672 | 1 | S018 | 0.97 | 13.11 | 0.70 | 19.82 | 0.31 |
| 0679 | 1 | $?$ | 0.74 | 15.17 | 0.86 | 22.33 | 0.44 |
| 0688 | 2 | Sc(s)II-III | 0.87 | 13.16 | 1.05 | 20.76 | 0.64 |
| 0692 | 2 | Sc(s)II.3 | 0.66 | 12.21 | 2.05 | 21.27 | 0.72 |
| 0697 | 1 | Sc(s)II.2 | 0.88 | 13.08 | 1.75 | 21.79 | 0.93 |
| 0699 | 2 | Sc(s)III-IV or SmI | 0.57 | 12.99 | 1.00 | 20.48 | 0.92 |
| 0713 | 2 | Sc(on-edge) | 1.13 | 12.25 | 1.77 | 20.98 | 0.22 |
| 0739 | 1 | SdIII-IV | 0.50 | 13.44 | 1.95 | 22.38 | 0.93 |
| 0740 | 2 | SBmIII | 0.55 | 15.40 | 0.80 | 22.41 | 0.53 |
| 0768 | 1 | SBc | 0.63 | 14.47 | 0.69 | 21.15 | 0.27 |
| 0792 | 2 | Sab(s) | 1.07 | 11.20 | 3.10 | 21.15 | 0.52 |
| 0801 | 2 | Amorphu | 0.69 | 11.86 | 0.82 | 18.92 | 0.56 |
| 0809 | 2 | Sc(on-edge) | 0.82 | 13.80 | 0.98 | 21.26 | 0.19 |
| 0827 | 2 | Sc(on-edge) | 0.91 | 12.64 | 1.48 | 20.98 | 0.19 |
| 0836 | 2 | Sab | 0.94 | 10.76 | 2.27 | 20.03 | 0.30 |
| 0849 | 2 | Sbc(s)II | 0.60 | 12.56 | 1.70 | 21.21 | 0.74 |
| 0851 | 2 | Sc(on-edge) | 0.90 | 13.13 | 1.33 | 21.25 | 0.21 |
| 0857 | 2 | SBb(sr)I-II | 1.10 | 10.48 | 3.48 | 20.69 | 0.91 |
| 0859 | 1 | Sc(on-edge) | 1.21 | 13.23 | 1.41 | 21.47 | 0.55 |
| 0865 | 2 | Sc(s)II | 0.72 | 12.17 | 2.50 | 21.65 | 0.30 |
| 0869 | 2 | ImV or dE0 | 0.86 | 14.71 | 2.36 | 24.07 | 0.80 |
| 0873 | 2 | Sc(on-edge) | 1.14 | 11.31 | 2.49 | 20.79 | 0.26 |
| 0874 | 1 | Sc(s) / S0 | 0.95 | 11.94 | 1.40 | 20.17 | 0.71 |
| 0905 | 2 | SBc(s)II | 0.68 | 12.74 | 2.61 | 22.32 | 0.91 |
| 0912 | 2 | SBbc(rs) | 0.91 | 11.85 | 2.06 | 20.92 | 0.67 |
| 0938 | 2 | SBc(s)II.2 | 0.83 | 12.22 | 1.90 | 21.10 | 0.90 |
| 0939 | 2 | Sc(s)II | 0.66 | 12.19 | 3.32 | 22.29 | 0.95 |
| 0945 | 2 | SBmIII | 0.31 | 14.86 | 1.29 | 22.91 | 0.48 |
| 0950 | 2 | SmIV | 0.47 | 14.73 | 2.28 | 24.01 | 0.77 |
|  |  |  |  |  |  |  |  |

continued

| VCC | ms | Type | $(\mathrm{g}-\mathrm{i})$ | $\mathrm{M}_{\mathrm{r}}$ | $\mathrm{R}_{\text {eff,r }}[\mathrm{kpc}]$ | $\langle\mu\rangle_{\text {eff,r }}$ | $\mathrm{b} / \mathrm{a}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $[1]$ | $[2]$ | $[3]$ | $[4]$ | $[5]$ | $[6]$ | $[7]$ | $[8]$ |
| 0957 | 1 | Sc(s)III | 0.65 | 12.00 | 1.21 | 19.91 | 0.46 |
| 0958 | 2 | Sa | 1.17 | 10.86 | 1.29 | 18.91 | 0.31 |
| 0975 | 2 | Scd(s)II | 0.00 | 0.00 | 0.00 | -inf | 0.00 |
| 0979 | 2 | Sapec | 0.90 | 11.27 | 2.58 | 20.82 | 0.53 |
| 0980 | 2 | Scd(s)III | 0.44 | 13.69 | 1.71 | 22.34 | 0.41 |
| 0984 | 2 | SBa | 1.08 | 11.72 | 1.14 | 19.50 | 0.33 |
| 0989 | 1 | Sc or Im | 0.56 | 16.04 | 0.79 | 23.01 | 0.97 |
| 0995 | 2 | Sc(on-edge) | 0.49 | 14.64 | 0.67 | 21.27 | 0.28 |
| 1011 | 2 | SdmIII | 0.72 | 14.32 | 1.25 | 22.29 | 0.51 |
| 1020 | 2 | dE4,N: or ImIV | 0.99 | 16.28 | 0.66 | 22.88 | 0.55 |
| 1047 | 2 | SBa(sr) | 1.11 | 11.54 | 1.28 | 19.57 | 0.88 |
| 1060 | 1 | SmIII-IV | 0.34 | 15.34 | 1.00 | 22.82 | 0.50 |
| 1086 | 2 | S(on-edge) | 1.11 | 12.54 | 1.16 | 20.35 | 0.23 |
| 1091 | 2 | Sbc(s)I.8 | 0.47 | 13.68 | 0.77 | 20.62 | 0.44 |
| 1110 | 2 | Sabpec | 1.11 | 9.70 | 3.82 | 20.10 | 0.64 |
| 1118 | 2 | Sc(s)III | 0.89 | 12.41 | 0.92 | 19.72 | 0.60 |
| 1126 | 2 | Sc / Sa | 1.02 | 11.91 | 1.61 | 20.44 | 0.44 |
| 1145 | 1 | RSb(rs)II | 1.07 | 10.47 | 1.46 | 18.78 | 0.71 |
| 1156 | 1 | SBcd(s)II | 0.46 | 13.97 | 2.07 | 23.04 | 0.47 |
| 1158 | 2 | Sa | 1.19 | 10.89 | 1.28 | 18.92 | 0.67 |
| 1165 | 2 | dE3 or ImV | 0.82 | 17.40 | 0.91 | 24.70 | 0.55 |
| 1179 | 2 | ImIII / BCD | 0.65 | 14.75 | 0.97 | 22.18 | 0.34 |
| 1186 | 2 | dE0 or ImV | 0.72 | 17.16 | 0.82 | 24.23 | 0.81 |
| 1188 | 2 | Amorphu | 0.90 | 12.80 | 0.83 | 19.89 | 0.44 |
| 1189 | 2 | Sc(s)II | 0.68 | 13.21 | 1.51 | 21.59 | 0.62 |
| 1190 | 2 | Sa | 1.18 | 10.95 | 1.98 | 19.92 | 0.37 |
| 1193 | 2 | Sc: | 0.72 | 13.52 | 0.88 | 20.73 | 0.30 |
| 1205 | 1 | ScIIIpec | 0.67 | 12.25 | 1.13 | 20.01 | 0.68 |
| 1208 | 1 | SmIII | 0.56 | 15.13 | 0.61 | 21.55 | 0.84 |
| 1217 | 2 | SBmIV | 0.62 | 14.02 | 2.62 | 23.60 | 0.74 |
| 1227 | 2 | dE0 or ImV | 1.52 | 16.52 | 1.45 | 24.81 | 0.88 |
| 1237 | 1 | $?$ | 0.61 | 15.22 | 0.66 | 21.80 | 0.76 |
|  |  |  |  |  |  |  |  |

continued

| VCC | ms | Type | $(\mathrm{g}-\mathrm{i})$ | $\mathrm{M}_{\mathrm{r}}$ | $\mathrm{R}_{\text {eff,r }, \mathrm{r}}[\mathrm{kpc}]$ | $\langle\mu\rangle_{\text {eff,r }}$ | $\mathrm{b} / \mathrm{a}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $[1]$ | $[2]$ | $[3]$ | $[4]$ | $[5]$ | $[6]$ | $[7]$ | $[8]$ |
| 1266 | 1 | SdmIII-IV | 0.35 | 14.91 | 1.19 | 22.79 | 0.83 |
| 1290 | 1 | Sb(r)II | 0.78 | 12.25 | 1.58 | 20.74 | 0.58 |
| 1294 | 1 | S0: | 0.94 | 14.69 | 0.91 | 21.98 | 0.64 |
| 1326 | 2 | SBa(s) | 0.96 | 12.26 | 1.38 | 20.44 | 0.52 |
| 1330 | 2 | Sa | 1.13 | 11.28 | 3.09 | 21.23 | 0.92 |
| 1331 | 2 | dE3 or ImV | 0.78 | 16.57 | 1.50 | 24.94 | 0.84 |
| 1336 | 2 | dE1 or ImV | 1.06 | 16.84 | 1.23 | 24.78 | 0.91 |
| 1356 | 2 | SmIII / BCD | 0.48 | 14.75 | 0.54 | 20.89 | 0.56 |
| 1358 | 1 | Sa: | 0.80 | 14.46 | 1.14 | 22.23 | 0.68 |
| 1375 | 1 | SBcIII-IV | 0.58 | 11.40 | 3.54 | 21.64 | 0.83 |
| 1379 | 2 | SBc(s)II | 0.68 | 11.93 | 2.16 | 21.09 | 0.48 |
| 1393 | 1 | SBc(s)II-III | 0.68 | 12.72 | 1.60 | 21.23 | 0.66 |
| 1408 | 2 | dE1 or ImV | 0.95 | 16.99 | 1.02 | 24.52 | 0.74 |
| 1410 | 2 | SmIII | 0.60 | 13.77 | 0.94 | 21.13 | 0.54 |
| 1412 | 2 | Sa | 1.18 | 10.71 | 1.59 | 19.21 | 0.48 |
| 1413 | 2 | dE2 or ImIV-V | 0.96 | 17.20 | 0.85 | 24.33 | 0.63 |
| 1419 | 2 | S(dust)pec | 0.97 | 12.54 | 1.66 | 21.13 | 0.73 |
| 1427 | 2 | Im / BCD: | 0.68 | 14.64 | 0.93 | 21.98 | 0.58 |
| 1442 | 1 | Sd(on-edge) | 0.63 | 13.85 | 1.27 | 21.86 | 0.20 |
| 1448 | 2 | ImIV or dE1pec | 0.86 | 13.13 | 3.27 | 23.19 | 0.77 |
| 1450 | 2 | Sc(s)II.2 | 0.57 | 12.44 | 1.98 | 21.42 | 0.68 |
| 1507 | 2 | SmIV: | 0.76 | 14.81 | 1.21 | 22.71 | 0.53 |
| 1508 | 2 | SBc(rs)II.2 | 0.52 | 11.72 | 2.57 | 21.27 | 0.73 |
| 1516 | 2 | Sc $/$ Sb: | 0.78 | 12.10 | 2.08 | 21.19 | 0.24 |
| 1524 | 2 | SBd(s)III | 0.57 | 12.64 | 3.39 | 22.78 | 0.88 |
| 1529 | 1 | SdmIII-IV | 0.74 | 13.93 | 1.72 | 22.60 | 0.82 |
| 1532 | 1 | SBcpec | 0.76 | 12.80 | 1.87 | 21.65 | 0.71 |
| 1540 | 1 | Sb(s)II | 1.22 | 10.17 | 3.07 | 20.10 | 0.30 |
| 1552 | 2 | Sapec | 1.05 | 11.25 | 2.48 | 20.71 | 0.67 |
| 1554 | 2 | SmIII | 0.45 | 11.57 | 1.65 | 20.14 | 0.47 |
| 1555 | 2 | SBc(s)I.3 | 0.87 | 9.69 | 6.67 | 21.31 | 0.67 |
| 1557 | 1 | Scd(on-edge) | 0.79 | 13.39 | 1.05 | 20.99 | 0.25 |
|  |  |  | 0 |  |  |  |  |

continued

| VCC | ms | Type | $(\mathrm{g}-\mathrm{i})$ | $\mathrm{M}_{\mathrm{r}}$ | $\mathrm{R}_{\text {eff, }, \mathrm{r}}[\mathrm{kpc}]$ | $\langle\mu\rangle_{\text {eff }, \mathrm{r}}$ | $\mathrm{b} / \mathrm{a}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $[1]$ | $[2]$ | $[3]$ | $[4]$ | $[5]$ | $[6]$ | $[7]$ | $[8]$ |
| 1562 | 1 | Sc(s)I | 0.87 | 10.16 | 4.52 | 20.93 | 0.37 |
| 1566 | 2 | SdIV | 0.69 | 14.48 | 1.35 | 22.63 | 0.61 |
| 1569 | 2 | Scd: | 0.62 | 14.88 | 1.23 | 22.82 | 0.60 |
| 1574 | 1 | $?$ | 0.89 | 16.45 | 0.44 | 22.16 | 0.80 |
| 1575 | 2 | SBmpec | 0.90 | 12.76 | 1.51 | 21.15 | 0.73 |
| 1581 | 2 | SmIV | 0.70 | 13.93 | 2.08 | 23.01 | 0.78 |
| 1582 | 2 | ImV or dE2 | 0.64 | 16.14 | 1.34 | 24.26 | 0.92 |
| 1588 | 2 | Scd(s)III-IV | 0.90 | 11.54 | 2.09 | 20.63 | 0.82 |
| 1605 | 2 | Sd: | 0.60 | 16.61 | 0.61 | 23.05 | 0.23 |
| 1615 | 2 | SBb(rs)I-II | 1.12 | 9.72 | 5.31 | 20.84 | 0.76 |
| 1623 | 1 | $?$ | 0.63 | 15.99 | 0.58 | 22.31 | 0.60 |
| 1624 | 1 | Sc(on-edge | 1.07 | 12.78 | 1.08 | 20.45 | 0.37 |
| 1644 | 2 | SmIV | 0.56 | 16.80 | 0.53 | 22.90 | 0.29 |
| 1656 | 2 | dE3 or ImIV-V | 0.90 | 15.88 | 1.24 | 23.84 | 0.66 |
| 1675 | 2 | pec | 0.83 | 13.72 | 1.50 | 22.10 | 0.62 |
| 1678 | 2 | SBdIV | 0.44 | 13.55 | 2.48 | 23.02 | 0.82 |
| 1685 | 1 | SBd(on-edge) | 0.46 | 15.42 | 0.97 | 22.84 | 0.24 |
| 1686 | 2 | SmIII | 0.57 | 12.61 | 2.34 | 21.95 | 0.52 |
| 1690 | 2 | Sab(s)I-II | 0.90 | 9.27 | 5.26 | 20.37 | 0.41 |
| 1696 | 2 | Sc(s)II-III | 0.96 | 10.81 | 4.36 | 21.50 | 0.85 |
| 1699 | 2 | SBmIII | 0.43 | 13.73 | 1.13 | 21.49 | 0.54 |
| 1713 | 1 | $?$ | 0.56 | 15.43 | 0.75 | 22.30 | 0.61 |
| 1725 | 2 | SmIII / BCD | 0.54 | 13.76 | 1.56 | 22.22 | 0.66 |
| 1726 | 2 | SdmIV | 0.28 | 14.63 | 1.83 | 23.44 | 0.74 |
| 1727 | 2 | Sab(s)II | 1.11 | 9.18 | 4.36 | 19.87 | 0.77 |
| 1730 | 2 | Sc / Sa | 1.08 | 11.49 | 2.02 | 20.51 | 0.71 |
| 1756 | 2 | dE5 or Im | 0.76 | 17.41 | 0.60 | 23.80 | 0.49 |
| 1757 | 2 | Sa(s)pec | 0.88 | 12.77 | 1.40 | 20.99 | 0.67 |
| 1758 | 2 | Sc(on-edge) | 0.74 | 14.09 | 0.85 | 21.24 | 0.14 |
| 1760 | 1 | Sa | 1.18 | 11.36 | 2.23 | 20.59 | 0.61 |
| 1776 | 2 | dE2 or ImIV | 1.15 | 16.78 | 1.89 | 25.66 | 0.89 |
| 1778 | 2 | Amorphu? | 0.65 | 13.45 | 1.10 | 21.15 | 0.50 |
|  |  |  |  |  |  |  |  |

continued

| VCC | ms | Type | $(\mathrm{g}-\mathrm{i})$ | $\mathrm{M}_{\mathrm{r}}$ | $\mathrm{R}_{\text {eff,r }}[\mathrm{kpc}]$ | $\langle\mu\rangle_{\text {eff,r }}$ | $\mathrm{b} / \mathrm{a}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $[1]$ | $[2]$ | $[3]$ | $[4]$ | $[5]$ | $[6]$ | $[7]$ | $[8]$ |
| 1791 | 2 | SBmIII / BCD | 0.34 | 13.89 | 1.43 | 22.15 | 0.53 |
| 1804 | 2 | ImIII / BCD | 0.82 | 15.05 | 0.92 | 22.37 | 0.49 |
| 1811 | 2 | Sc(s)II.8 | 0.60 | 12.13 | 1.53 | 20.55 | 0.66 |
| 1813 | 2 | SBa | 1.14 | 10.03 | 3.22 | 20.07 | 0.78 |
| 1821 | 1 | $?$ | 0.79 | 16.88 | 0.48 | 22.75 | 0.56 |
| 1825 | 2 | dE2 or ImIV | 0.82 | 15.29 | 1.08 | 22.96 | 0.87 |
| 1834 | 1 | S016,N | 1.07 | 12.20 | 1.29 | 20.25 | 0.75 |
| 1837 | 1 | $?$ | 0.81 | 15.97 | 0.99 | 23.44 | 0.49 |
| 1855 | 1 | S0: | 0.86 | 14.72 | 0.84 | 21.83 | 0.69 |
| 1859 | 2 | Sapec | 0.94 | 11.66 | 1.91 | 20.56 | 0.49 |
| 1860 | 1 | $?$ | 1.05 | 16.22 | 0.84 | 23.33 | 0.68 |
| 1868 | 2 | Scd(on-edge) | 1.14 | 12.57 | 1.57 | 21.04 | 0.20 |
| 1873 | 1 | $?$ | 0.39 | 16.27 | 0.59 | 22.61 | 0.43 |
| 1883 | 2 | RSB01/2 | 1.05 | 11.14 | 1.40 | 19.36 | 0.76 |
| 1884 | 2 | dE1 or ImV | 0.70 | 15.83 | 2.80 | 25.55 | 0.92 |
| 1898 | 1 | $?$ | 0.90 | 14.99 | 1.05 | 22.59 | 0.72 |
| 1900 | 2 | ImV or dE3 | 0.74 | 15.48 | 1.69 | 24.12 | 0.78 |
| 1902 | 2 | S0 / Sa | 0.93 | 12.34 | 1.82 | 21.13 | 0.88 |
| 1905 | 2 | dE2 or ImIV | 1.17 | 16.74 | 2.39 | 26.13 | 0.92 |
| 1906 | 1 | S0: | 0.94 | 15.00 | 0.51 | 21.02 | 0.79 |
| 1920 | 1 | S0? | 0.95 | 14.68 | 0.70 | 21.39 | 0.86 |
| 1923 | 1 | Sbc(s)II-III | 0.84 | 12.06 | 1.70 | 20.71 | 0.76 |
| 1929 | 2 | Scd(s) | 0.73 | 12.88 | 1.82 | 21.68 | 0.40 |
| 1932 | 2 | Sc(on-edge) | 0.89 | 12.13 | 1.42 | 20.39 | 0.17 |
| 1933 | 1 | Sab? | 0.53 | 15.98 | 1.11 | 23.70 | 0.94 |
| 1943 | 2 | SBb(r)II | 0.91 | 11.29 | 1.73 | 19.97 | 0.79 |
| 1944 | 1 | BCD? | 1.16 | 16.19 | 0.34 | 21.37 | 0.96 |
| 1955 | 2 | Spec / BCD | 0.85 | 12.74 | 2.33 | 22.07 | 0.88 |
| 1960 | 1 | ImIII / BCD? | 0.74 | 17.39 | 0.37 | 22.73 | 0.58 |
| 1987 | 2 | SBc(rs)II | 0.80 | 10.23 | 3.94 | 20.70 | 0.61 |
| 1994 | 2 | dE2 or ImV | 0.76 | 16.44 | 0.99 | 23.92 | 0.90 |
| 1999 | 2 | Sa | 1.06 | 11.93 | 1.14 | 19.72 | 0.75 |
|  |  |  |  |  |  |  |  |

continued

| VCC | ms | Type | $(\mathrm{g}-\mathrm{i})$ | $\mathrm{M}_{\mathrm{r}}$ | $\mathrm{R}_{\text {eff,r } \mathrm{r}}[\mathrm{kpc}]$ | $\langle\mu\rangle_{\text {eff,r }}$ | $\mathrm{b} / \mathrm{a}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $[1]$ | $[2]$ | $[3]$ | $[4]$ | $[5]$ | $[6]$ | $[7]$ | $[8]$ |
| 2006 | 2 | Amorphu | 0.76 | 13.06 | 1.62 | 21.60 | 0.30 |
| 2007 | 2 | ImIII / BCD: | 0.69 | 15.20 | 0.68 | 21.85 | 0.79 |
| 2023 | 2 | SBc(s)II | 0.57 | 13.32 | 1.49 | 21.68 | 0.43 |
| 2058 | 2 | Sc(s)II.3 | 0.92 | 10.59 | 4.41 | 21.31 | 0.78 |
| 2066 | 2 | Amorphu | 0.83 | 11.25 | 1.70 | 19.90 | 0.50 |
| 2070 | 2 | Sa | 1.16 | 10.28 | 2.48 | 19.75 | 0.73 |
| 2089 | 1 | BCD? | 1.35 | 15.04 | 0.35 | 20.28 | 0.79 |

Derived structural parameters of the late-type galaxies.

### 6.3 Results of GAZELLE

Table 6.4 summarised the results for the galaxies from the GALEV/GAZELLE runs of our study, which were analysed in the course of Chapter 4. The complete latetype morphological sequence was included in the GALEV/GAZELLE runs and is summarised in Tab. 6.3. Missing internal identification numbers (Type-ID; e.g. BCDs $=1$ and 2) in Tab. 6.3 have been used to mark the possible membership to the Virgo cluster in the course of the analysis. Since the differences between certain and possible members are not significant, we did not differentiate the membership within the GALEV/GAZELLE runs.

Table 6.3: The different morphological types according to the VCC and their internal identification number.

| Type-ID | Description |
| :---: | :--- |
| 1 | BCD |
| 3 | dI |
| 5 | Sb |
| 7 | Sc |
| 9 | Sd |
| 11 | Sm |
| 15 | Amorphous |
| 17 | Sa |
| 19 | S 0 |
| 23 | S |
| 24 | unknown |
| 25 | dE/dI |

Table 6.4: Results from GALEV/GAZELLE runs (Part I).
Upper and lower values correspond to the minimum and maximum values of the GAZELLE output.

| VCC | Type- <br> ID | $\begin{gathered} \mathrm{mag}_{\mathrm{r}, \mathrm{obs}} \\ {[\mathrm{mag}]} \end{gathered}$ | mag $_{\text {r,mod }}$ [mag] | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \mathrm{M}_{\text {nonvis }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 0001 | 1 | $14.82 \pm 0.05$ | $14.85 \pm 0.10$ | $8.73 e+08_{7.88 e+09}^{8.73 e+08}$ | $5.79 e+08_{5.83 e+08}^{4.17 e+08}$ | $2.94 e+08_{7.41 e+09}^{2.94++08}$ | $8.33 e+08_{9}^{8.53 e+e+09}$ |
| 0004 | 3 | $16.29 \pm 0.06$ | $16.32 \pm 0.10$ | $8.83 e+08_{8.83 e+08}^{1.66+08}$ | $4.71 e+07_{5.24 e+07}^{4.64+07}$ | $8.36 e+08_{8.36 e+08}^{1.19 e+08}$ | $6.62 e+07_{4.74 e+08}^{6.62 e+07}$ |
| 0010 | 1 | $15.07 \pm 0.05$ | $15.09 \pm 0.10$ | $4.79 e+09_{4.79 e+09}^{6.11 e+08}$ | $2.55 e+08_{3.69 e+08}^{2.55 e+08}$ | $4.53 e+09_{4.53 e+09}^{2.42 e+08}$ | $3.59 e+08_{2.13 e+09}^{3.59+08}$ |
| 0015 | 7 | $14.67 \pm 0.05$ | $14.71 \pm 0.10$ | $7.09 e+09_{7.09 e+09}^{8.75 e+08}$ | $3.78 e+08_{5.28 e+08}^{3.78 e+08}$ | $6.71 e+09_{6.71 e+09}^{3.47 e+08}$ | $5.32 e+08_{6.30 e+09}^{5.32 e+08}$ |
| 0017 | 3 | $15.30 \pm 0.05$ | $15.25 \pm 0.10$ | $9.99 e+08_{2.52 e+09}^{4.59+08}$ | $1.46 e+08_{1.46 e+08}^{1.29 e+08}$ | $8.53 e+08_{2.38 e+09}^{3.31 e+08}$ | $2.05 e+08_{1.16 e+09}^{2.05+08}$ |
| 0020 | 24 | $17.38 \pm 0.08$ | $17.36 \pm 0.10$ | $1.01 e+08_{2.22 e+08}^{8.76 e+07}$ | $2.82 e+07_{3.86 e+07}^{2.82 e+07}$ | $7.24 e+07_{1.89 e+08}^{4.90++7}$ | $3.87 e+07_{4.18 e+08}^{3.87+07}$ |
| 0022 | 1 | $15.65 \pm 0.05$ | $15.67 \pm 0.10$ | $2.47 e+09^{3.4720 e+08}$ | $1.32 e+08_{1.72 e+08}^{1.32+08}$ | $2.34 e+09_{2.34 e+09}^{2.18 t+08}$ | $1.85 e+08_{2.53 e+09}^{1.85+08}$ |
| 0024 | 1 | $14.88 \pm 0.05$ | $14.90 \pm 0.10$ | $4.95 e+09_{4.95 e+09}^{7.83 e+08}$ | $2.64 e+08_{3.45 e+08}^{2.64+08}$ | $4.68 e+09_{4.68 e+09}^{4.38++8}$ | $3.71 e+08^{3.709 e+09} 3$ |
| 0025 | 7 | $11.92 \pm 0.05$ | $11.92 \pm 0.10$ | $6.67 e+10{ }_{6}^{6.667 e+10}$ | $3.56 e+09^{3} 3.566++09$ | $6.31 e+10_{6.31 e+10}^{6.310+10}$ | $5.00 e+09_{1.50 e+10}^{5.00 e+09}$ |
| 0026 | 3 | $17.34 \pm 0.08$ | $17.30 \pm 0.10$ | $6.62 e+07_{3.53 e+08}^{6.62 e+07}$ | $1.85 e+07_{2.80 e+07}^{1.85 e+07}$ | $4.76 e+07_{3.34 e+08}^{4.76+07}$ | $2.54 e+07_{2.30 e+08}^{2.54 e+07}$ |
| 0030 | 3 | $16.03 \pm 0.06$ | $16.06 \pm 0.10$ | $2.09 e+08_{1.13 e+09}^{2.09+08}$ | $5.86 e+07_{8.73 e+07}^{5.86+07}$ | $1.51 e+08_{1.07 e+09}^{1.51 e+08}$ | $8.04 e+07_{7.24 e+08}^{8.04 e+07}$ |
| 0031 | 24 | $15.20 \pm 0.05$ | $15.22 \pm 0.10$ | $3.23 e+09_{3.23 e+09}^{5.19 e+08}$ | $1.72 e+08_{2.29 e+08}^{1.72 e+08}$ | $3.06 e+09_{3.06 e+09}^{2.90 e+08}$ | $2.42 e+08_{2.77 e+09}^{2.42 e+08}$ |
| 0034 | 7 | $14.36 \pm 0.05$ | $14.39 \pm 0.10$ | $6.08 e+09_{6.088 e+09}^{1.28 e+09}$ | $3.24 e+08_{3.98 e+08}^{3.24 e+08}$ | $5.75 e+09_{5.76 e+09}^{8.86+08}$ | $4.56 e+08_{3.48 e+09}^{4.56 e+08}$ |
| 0041 | 3 | $16.34 \pm 0.09$ | $16.23 \pm 0.10$ | $1.35 e+09_{1.35 e+09}^{2.212 e+08}$ | $7.21 e+07_{1.62 e+08}^{7.900+07}$ | $1.28 e+09_{1.28 e+09}^{8.2 l e+07}$ | $1.01 e+08_{2.06 e+09}^{1.01 e+08}$ |
| 0048 | 9 | $14.03 \pm 0.05$ | $14.03 \pm 0.10$ | $7.95 e+09_{7.95 e+09}^{1.78+09}$ | $4.24 e+00_{5.52 e+08}^{8.24 e+08}$ | $7.53 e+09_{7.53 e+09}^{1.2 e x+09}$ | $5.96 e+08_{4.64 e+09}^{5.96 e+08}$ |

continued

| VCC | Type- <br> ID | mag $_{\mathrm{r}, \mathrm{obs}}$ <br> [mag] | $\operatorname{mag}_{r, \bmod }$ <br> [mag] | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \hline \mathrm{M}_{\text {nonvis }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 0052 | 3 | $17.04 \pm 0.09$ | $17.04 \pm 0.10$ | $6.75 e+08_{6.76 e+08}^{1.06+08}$ | $3.60 e+07_{4.69 e+07}^{3.60 e+07}$ | $6.39 e+08_{6.40 e+08}^{5.95+07}$ | $5.07 e+07_{5.20 e+08}^{5.07 e+07}$ |
| 0058 | 5 | $12.50 \pm 0.05$ | $12.53 \pm 0.10$ | $5.12 e+10_{5.12 e+10}^{66.49+09}$ | $2.73 e+09_{3.92 e+09}^{2.73 e+09}$ | $4.84 e+10_{4.85 e+10}^{2.58 e+09}$ | $3.84 e+09_{3.10 e+10}^{3.84 e+09}$ |
| 0066 | 7 | $11.17 \pm 0.05$ | $11.17 \pm 0.10$ | $1.33 e+11_{1.33 e+11}^{2.16 e+10}$ | $7.08 e+09_{9.52 e+09}^{7.08 e+09}$ | $1.26 e+11_{1.26 e+11}^{1.21 e+10}$ | $9.96 e+09_{8.12 e+10}^{99.96+09}$ |
| 0067 | 7 | $13.51 \pm 0.05$ | $13.53 \pm 0.10$ | $1.49 e+10_{1.49 e+10}^{3.12 e+09}$ | $7.95 e+08_{9.66 e+08}^{7.95 e+08}$ | $1.41 e+10_{1.41 e+10}^{2.15 e+09}$ | $1.12 e+09_{7.13 e+09}^{1.12 e+09}$ |
| 0073 | 5 | $12.20 \pm 0.05$ | $12.19 \pm 0.10$ | $1.48 e+10_{1.33 e+11}^{1.43 e+10}$ | $9.84 e+09_{9.89 e+09}^{7.07 e+09}$ | $4.99 e+09_{1.25 e+11}^{4.47 e+09}$ | $1.42 e+10_{1.19 e+11}^{1.42 e+10}$ |
| 0074 | 1 | $15.88 \pm 0.05$ | $15.90 \pm 0.10$ | $2.19 e+09_{2.19 e+09}^{3.45+08}$ | $1.17 e+08_{1.52 e+08}^{1.17 e+08}$ | $2.07 e+09^{1.07 e++09}$ | $1.64 e+08_{9.24 e+08}^{1.64 e+08}$ |
| 0081 | 7 | $14.43 \pm 0.05$ | $14.44 \pm 0.10$ | $5.41 e+09_{5.41 e+09}^{1.22 e+09}$ | $2.89 e+08_{3.78 e+08}^{2.89+08}$ | $5.12 e+09_{5.12 e+09}^{8.41 e+08}$ | $4.06 e+08_{3.16 e+09}^{4.06+08}$ |
| 0083 | 3 | $15.11 \pm 0.05$ | $15.12 \pm 0.10$ | $3.91 e+09_{3.91 e+09}^{6.19 e+08}$ | $2.09 e+08_{2.73 e+08}^{2.09 e+08}$ | $3.70 e+09^{3.746 e+08}$ | $2.93 e+08_{2.37 e+09}^{2.93+08}$ |
| 0085 | 3 | $16.80 \pm 0.10$ | $16.83 \pm 0.10$ | $1.42 e+08_{6.75 e+08}^{1.19 e+08}$ | $4.39 e+07_{5.27 e+07}^{3.60+07}$ | $9.78 e+07_{6.39 e+08}^{6.68 e+07}$ | $6.21 e+07_{5.88 e+08}^{6.21 e+07}$ |
| 0087 | 11 | $14.29 \pm 0.05$ | $14.26 \pm 0.10$ | $6.44 e+09_{6.44 e+09}^{1.44+09}$ | $3.44 e+08_{4.45 e+08}^{3.44 e+08}$ | $6.10 e+09_{6.10 e+09}^{99.91 e+08}$ | $4.83 e+08_{3.12 e+09}^{4.83+08}$ |
| 0089 | 7 | $11.45 \pm 0.05$ | $11.50 \pm 0.10$ | $1.43 e+11_{1.43 e+11}^{1.77 e+10}$ | $7.64 e+09_{1.07 e+10}^{7.6 e+09}$ | $1.36 e+11_{1.36 e+11}^{7.02 e+09}$ | $1.07 e+10_{1.14 e+11}^{1.02 e+10}$ |
| 0092 | 5 | $9.71 \pm 0.05$ | $9.68 \pm 0.10$ | $1.30 e+11_{1.61 e+11}^{1.30+11}$ | $8.90 e+10_{1.10 e+11}^{8.90 e+10}$ | $4.06 e+10_{5.07 e+10}^{4.06 e+10}$ | $1.29 e+11_{1.10 e+12}^{1.29 e+11}$ |
| 0093 | 25 | $16.35 \pm 0.05$ | $16.39 \pm 0.10$ | $8.83 e+08_{8.83 e+08}^{1.64++08}$ | $4.71 e+07_{7.82 e+07}^{4.59+07}$ | $8.36 e+08_{8.36 e+08}^{9.93 e+07}$ | $6.62 e+07_{7.60 e+08}^{6.62 e+07}$ |
| 0094 | 19 | $12.49 \pm 0.05$ | $12.47 \pm 0.10$ | $1.23 e+10_{1.45 e+10}^{1.23 e+10}$ | $8.47 e+09_{9.88 e+09}^{8.47 e+09}$ | $3.87 e+09_{4.60 e+09}^{3.87 e+09}$ | $1.23 e+10_{1.49 e+11}^{1.23 e+10}$ |
| 0097 | 7 | $12.05 \pm 0.05$ | $12.04 \pm 0.10$ | $1.37 e+10_{1.23 e+11}^{1.37 e+10}$ | $9.08 e+09_{9.78 e+09}^{6.54 e+09}$ | $4.60 e+09_{1.16 e+11}^{4.46 e+09}$ | $1.31 e+10_{9.89 e+10}^{1.31 e+10}$ |
| 0099 | 17 | $14.36 \pm 0.05$ | $14.35 \pm 0.10$ | $1.63 e+09_{1.46 e+10}^{1.63 e+09}$ | $1.08 e+09_{1.17 e+09}^{77.78+08}$ | $5.49 e+08_{1.38 e+10}^{5.36 e+08}$ | $1.56 e+09_{1.18 e+10}^{1.56+09}$ |
| 0104 | 25 | $17.04 \pm 0.14$ | $17.07 \pm 0.10$ | $4.19 e+08_{4.19 e+08}^{6.67 e+07}$ | $2.24 e+07_{9.79 e+07}^{1.87 e+07}$ | $3.97 e+08_{3.97 e+08}^{3.89+07}$ | $3.14 e+07_{9,72 e+08}^{3.14 e+07}$ |
| 0105 | 9 | $13.14 \pm 0.05$ | $13.14 \pm 0.10$ | $2.48 e+10_{2.48 e+10}^{9,75 e+10}$ | $1.32 e+09_{1.42 e+09}^{1.32 e+09}$ | $2.34 e+10_{2.34 e+10}^{8.32 e+09}$ | $1.86 e+09_{9.58 e+09}^{1.86+09}$ |
| 0113 | 24 | $16.61 \pm 0.07$ | $16.57 \pm 0.10$ | $1.42 e+08_{3.09 e+08}^{1.42 e+08}$ | $3.98 e+07_{5.50 e+07}^{3.98 e+07}$ | $1.03 e+08_{2.64 e+08}^{1.02 e+08}$ | $5.47 e+07_{4.47 e+08}^{5.47 e}$ |
| 0114 | 3 | $16.14 \pm 0.07$ | $15.93 \pm 0.10$ | $2.31 e+09_{2.31 e+09}^{2.85 e+08}$ | $1.23 e+08_{2.15 e+08}^{1.17 e+08}$ | $2.18 e+09_{2.18 e+09}^{1.09 e+08}$ | $1.73 e+08_{2.89 e+09}^{1.73 e+08}$ |

continued

| VCC | $\begin{gathered} \hline \hline \text { Type- } \\ \text { ID } \end{gathered}$ | $\begin{gathered} \text { mag }_{\mathrm{r}, \mathrm{bb}} \\ {[\mathrm{mag}]} \end{gathered}$ | $\begin{gathered} \mathrm{mag}_{\mathrm{r}, \mathrm{mod}} \\ {[\mathrm{mag}]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \mathrm{M}_{\text {nonvis }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 0117 | 3 | $15.97 \pm 0.05$ | $15.98 \pm 0.10$ | $1.15 e+09_{1.15 e+09}^{1.92 e+08}$ | $6.11 e+07_{6.11 e+07}^{5.37+07}$ | $1.08 e+09_{1.08 e+09}^{1.38++08}$ | $8.59 e+07_{4.05 e+08}^{8.59 e+07}$ |
| 0119 | 7 | $14.85 \pm 0.05$ | $14.90 \pm 0.10$ | $3.11 e+09_{3.11 e+09}^{5.38+08}$ | $1.66 e+08_{1.92 e+08}^{1.51 e+08}$ | $2.94 e+09_{2.94 e+09}^{3.88+08}$ | $2.33 e+08_{1.45 e+09}^{2.33 e+08}$ |
| 0120 | 7 | $12.58 \pm 0.05$ | $12.59 \pm 0.10$ | $4.09 e+10_{4.09 e+10}^{4.09 e+10}$ | $2.18 e+09_{2.18 e+09}^{2.18+09}$ | $3.87 e+10_{3.88 e+10}^{3.87 e+10}$ | $3.07 e+09_{9.2 l e+09}^{3.37 e+09}$ |
| 0122 | 19 | $12.75 \pm 0.05$ | $12.74 \pm 0.10$ | $8.64 e+09_{9.93 e+09}^{7.84+09}$ | $5.93 e+09_{6.78 e+09}^{5.38+09}$ | $2.71 e+09_{3.15 e+09}^{2.46 e+09}$ | $8.64 e+00_{1.00 e+11}^{89.64 e+09}$ |
| 0126 | 9 | $14.01 \pm 0.05$ | $14.03 \pm 0.10$ | $1.05 e+10_{1.05 e++10}^{1.65 s+}$ | $5.59 e+08_{7.29 e+08}^{5.5 e+08}$ | $9.92 e+09_{9.92 e+08}^{99.25 e+08}$ | $7.85 e+08_{7.08 e+09}^{1.85+08}$ |
| 0128 | 25 | $14.78 \pm 0.06$ | $14.78 \pm 0.10$ | $8.54 e+09_{8.55 e+09}^{1.05+09}$ | $4.56 e+08_{6.36 e+08}^{4.56+08}$ | $8.09 e+09_{8.09 e+09}^{4.18 e+08}$ | $6.41 e+08_{4.58 e+09}^{6.41 e+08}$ |
| 0130 | 1 | $16.44 \pm 0.05$ | $16.45 \pm 0.10$ | $7.80 e+08_{7.80 e+08}^{1.36+08}$ | $4.16 e+07_{4.97 e+07}^{4.007}$ | $7.38 e+08_{7.39 e+08}^{8.80+07}$ | $5.85 e+076.352 e+08$ |
| 0131 | 7 | $13.45 \pm 0.05$ | $13.44 \pm 0.10$ | $3.40 e+09_{3.40 e++09}^{3.409}$ | $2.26 e+09_{2.26 e+09}^{2.26 e+09}$ | $1.15 e+09_{1.15 e+09}^{1.15 e+09}$ | $3.25 e+00_{9.74 e+09}^{3.25 e+09}$ |
| 0132 | 3 | $16.01 \pm 0.08$ | $15.94 \pm 0.10$ | $1.33 e+09_{1.33 e+09}^{2.47 e+08}$ | $7.11 e+07_{9.55 e+07}^{6.91 e+07}$ | $1.26 e+09_{1.26 e+09}^{1.78+08}$ | $9.99 e+079.9890 e+07$ |
| 0135 | 23 | $13.75 \pm 0.05$ | $13.75 \pm 0.10$ | $2.81 e+09_{2.52 e+10}^{2.81 e+09}$ | $1.87 e+09_{2.03 e+09}^{1.34+09}$ | $9.47 e+08_{2.38 e+10}^{99.28 e+10}$ | $2.69 e+0{ }_{2}^{2.049++09}$ |
| 0137 | 25 | $16.19 \pm 0.05$ | $16.22 \pm 0.10$ | $2.36 e+08_{1.92 e+09}^{2.36 e+08}$ | $1.43 e+08_{1.43 e+08}^{1.02 e+08}$ | $9.37 e+07_{1.81 e+09}^{9.37+07}$ | $2.04 e+08_{1.63 e+09}^{2.04+08}$ |
| 0143 | 7 | $14.78 \pm 0.05$ | $14.84 \pm 0.10$ | $6.63 e+09_{6.63 e+09}^{8.19 e+08}$ | $3.54 e+08_{4.94 e+08}^{3.54+08}$ | $6.28 e+09_{6.28 e+09}^{3.25 e+08}$ | $4.97 e+085.8579 e+09$ |
| 0144 | 1 | $14.63 \pm 0.05$ | $14.73 \pm 0.10$ | $4.02 e+08_{5.27 e+08}^{3.41 e+08}$ | $1.43 e+08_{1.86 e+08}^{1.43 e+08}$ | $2.59 e+08_{3.43 e+08}^{1.73 e+08}$ | $1.97 e+08_{2.05 e+09}^{1.97 e+08}$ |
| 0145 | 7 | $12.02 \pm 0.05$ | $12.04 \pm 0.10$ | $9.13 e+10_{9.13 e+10}^{1.13 e+10}$ | $4.87 e+09_{6.81 e+09}^{4.87 e+09}$ | $8.64 e+10_{8.64 e+109}^{4.47 e+9}$ | $6.84 e+00_{5.80 e+10}^{69.64 t+09}$ |
| 0152 | 7 | $12.35 \pm 0.05$ | $12.32 \pm 0.10$ | $1.17 e+10_{1.28 e+10}^{1.17 e+10}$ | $8.02 e+09_{8.78 e+09}^{8.02 e+09}$ | $3.66 e+09_{4.09 e+09}^{3.66 e+09}$ | $1.16 e+10_{8.36 e+10}^{1.16 e+10}$ |
| 0155 | 25 | $14.47 \pm 0.05$ | $14.49 \pm 0.10$ | $9.65 e+09_{9.65 e+09}^{1.19 e+09}$ | $5.15 e+08_{8.16 e+08}^{5.15 e+08}$ | $9.14 e+09_{9.14 e+09}^{4.14 e+08}$ | $7.23 e+08_{1.00 e+10}^{7.23 e+08}$ |
| 0157 | 7 | $10.82 \pm 0.05$ | $10.81 \pm 0.10$ | $3.57 e+10_{3.13 e+11}^{3.57 e+10}$ | $2.37 e+10_{2.37 e+10}^{1.67+10}$ | $1.20 e+10_{2.96 e+11}^{1.20 e+10}$ | $3.41 e+10_{1.83 e+11}^{3.41 e+10}$ |
| 0159 | 3 | $15.38 \pm 0.05$ | $15.37 \pm 0.10$ | $3.94 e+08_{8.81 e+08}^{3.94 e+08}$ | $1.10 e+08_{1.29 e+08}^{1.10 e+08}$ | $2.84 e+08_{7.52 e+08}^{2.83 e+08}$ | $1.51 e+08_{8.16 e+08}^{1.51+08}$ |
| 0162 | 9 | $14.19 \pm 0.05$ | $14.21 \pm 0.10$ | $7.21 e+09_{7.21 e+09}^{2.84 e+09}$ | $3.85 e+08_{4.16+0+08}^{4.35+08}$ | $6.82 e+09_{6.83 e+09}^{2.432+09}$ | $5.41 e+00_{2.79 e+09}^{8.412+08}$ |
| 0166 | 19 | $11.89 \pm 0.05$ | $11.89 \pm 0.10$ | $2.24 e+10_{2.73 e+10}^{2.24 e+10}$ | $1.54 e+10_{1.86 e+10}^{1.54+10}$ | $7.03 e+09_{8.72 e+09}^{7.03 e+09}$ | $2.24 e+10_{3.25 e+11}^{2.24 e+10}$ |

continued

| VCC | TypeID | mag $_{\mathrm{r}, \mathrm{obs}}$ <br> [mag] | $\operatorname{mag}_{\mathrm{r}, \mathrm{mod}}$ <br> [mag] | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \hline \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \hline \hline \mathrm{M}_{\text {nonvis }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 0167 | 5 | $9.69 \pm 0.05$ | $9.67 \pm 0.10$ | $2.25 e+11_{2.64 e+11}^{2.25 e+11}$ | $1.55 e+11_{1.80 e+11}^{1.54 e+11}$ | $7.05 e+10_{8.39 e+10}^{7.05 e+10}$ | $2.25 e+11_{2.72 e+12}^{2.25 e+11}$ |
| 0168 | 3 | $16.40 \pm 0.06$ | $16.39 \pm 0.10$ | $1.61 e+09_{1.61 e+09}^{1.99 e+08}$ | $8.60 e+07_{1.20 e+08}^{8.80 e+07}$ | $1.53 e+09_{1.53 e+09}^{7.89+07}$ | $1.21 e+08_{1.26 e+09}^{1.21 e+08}$ |
| 0169 | 3 | $16.29 \pm 0.11$ | $16.22 \pm 0.10$ | $1.49 e+08_{9.39 e+08}^{1.49 e+08}$ | $4.17 e+07_{5.86 e+07}^{4.17 e+07}$ | $1.07 e+08_{8.89 e+08}^{1.07 e+08}$ | $5.73 e+0{ }^{5} 5.783 e+08$ |
| 0172 | 1 | $14.32 \pm 0.05$ | $14.34 \pm 0.10$ | $5.95 e+09_{5.95 e+09}^{1.03+09}$ | $3.17 e+08_{4}^{3.014 e+08}$ | $5.63 e+09_{5}^{6.64 e++08}$ | $4.46 e+08_{6}^{4.455+0+08}$ |
| 0181 | 25 | $17.48 \pm 0.06$ | $17.52 \pm 0.10$ | $3.97 e+08_{3.97 e+08}^{6.28 e+07}$ | $2.12 e+07_{2.77 e+07}^{2.12 e+07}$ | $3.76 e+08_{3.76 e+08}^{3.51 e+07}$ | $2.97 e+07_{2.68 e+08}^{2.97 e+07}$ |
| 0187 | 7 | $12.87 \pm 0.05$ | $12.88 \pm 0.10$ | $5.23 e+10_{5.23 e+10}^{6.46 e+09}$ | $2.79 e+09_{3.90 e+09}^{2.79+09}$ | $4.95 e+10_{4.95 e+10}^{2.56 e+09}$ | $3.92 e+00_{2.29 e+10}^{3.92 e+09}$ |
| 0190 | 25 | $16.44 \pm 0.09$ | $16.44 \pm 0.10$ | $5.43 e+08_{1.35 e+09}^{1.79+08}$ | $7.94 e+07_{1.36 e+08}^{7.05 e+07}$ | $4.64 e+08_{1.28 e+09}^{6.90 e+07}$ | $1.12 e+08_{1.75 e+09}^{1.12 e+08}$ |
| 0199 | 17 | $11.46 \pm 0.05$ | $11.45 \pm 0.10$ | $3.85 e+10_{4.51 e+10}^{3.85 e+10}$ | $2.64 e+10_{3.08 e+10}^{2.64 e+10}$ | $1.21 e+10_{1.43 e+10}^{1.21 e+10}$ | $3.85 e+10_{4.64 e+11}^{3.85 e+10}$ |
| 0207 | 1 | $16.55 \pm 0.05$ | $16.58 \pm 0.10$ | $9.06 e+07_{1.04 e+08}^{8.08 t+07}$ | $3.20 e+07_{3.48 e+07}^{2.86 e+07}$ | $5.86 e+07_{6}^{5.29 e++07}$ | $4.45 e+07_{2.63 e+08}^{4.45 e 07}$ |
| 0213 | 23 | $13.59 \pm 0.05$ | $13.62 \pm 0.10$ | $2.07 e+10_{2.07 e+10}^{2.56 e+09}$ | $1.11 e+09_{1.54 c+09}^{1.11 e+09}$ | $1.96 e+10_{1.96 e+10}^{1.02 e+09}$ | $1.55 e+0{ }_{1.45 e+10}^{1.55 e+09}$ |
| 0217 | 3 | $15.40 \pm 0.07$ | $15.35 \pm 0.10$ | $3.66 e+08_{2.088 e+09}^{3.666+08}$ | $1.03 e+08_{1.68 e+08}^{1.032+08}$ | $2.64 e+08_{1.97 e+09}^{2.6 e+08}$ | $1.41 e+08_{1.34 e+09}^{1.45 e+08}$ |
| 0221 | 7 | $12.56 \pm 0.05$ | $12.58 \pm 0.10$ | $4.28 e+10_{4.28 e+10}^{1.68 e+10}$ | $2.28 e+09_{2.46 e+09}^{2.28 e+09}$ | $4.05 e+10_{4.05 e+10}^{1.44 e+10}$ | $3.21 e+00_{1.65 e+10}^{3.21 e+09}$ |
| 0222 | 17 | $11.32 \pm 0.05$ | $11.33 \pm 0.10$ | $4.27 e+10_{5.20 e+10}^{4.27 e+10}$ | $2.93 e+10_{3.54 e+10}^{2.93 e+10}$ | $1.34 e+10_{1.66 e+10}^{1.34+10}$ | $4.27 e+10_{5.67 e+11}^{4.27 e+10}$ |
| 0223 | 1 | $15.77 \pm 0.05$ | $15.80 \pm 0.10$ | $2.04 e+09^{3} 3.243 e+09$ | $1.09 e+08_{1.42 e+08}^{1.09 e+08}$ | $1.93 e+09^{1.83 e++09}$ | $1.53 e+08_{1.03 e+09}^{1.53 e+08}$ |
| 0226 | 7 | $11.35 \pm 0.05$ | $11.34 \pm 0.10$ | $2.64 e+10_{2.37 e+11}^{2.64 e+10}$ | $1.75 e+10_{1.75 e+10}^{1.26 e+10}$ | $8.89 e+09_{2.24 e+11}^{8.89+09}$ | $2.52 e+10_{1.11 e+11}^{2.52 e+10}$ |
| 0234 | 17 | $11.68 \pm 0.05$ | $11.66 \pm 0.10$ | $2.87 e+10_{3.07 e+10}^{2.62 e+10}$ | $1.96 e+10_{2.09 e+10}^{1.80 e+10}$ | $9.06 e+09^{8.2 .76 e++09}$ | $2.87 e+10_{3.18 e+11}^{2.87 e+10}$ |
| 0237 | 29 | $16.56 \pm 0.07$ | $16.55 \pm 0.10$ | $1.24 e+08_{2.94+08}^{1.24+08}$ | $3.46 e+07_{4.29 e+07}^{3.46 e+07}$ | $8.92 e+07_{2.51+08}^{8.90 e+07}$ | $4.75 e+07_{2.63 e+08}^{4.75 e+07}$ |
| 0241 | 9 | $14.19 \pm 0.05$ | $14.22 \pm 0.10$ | $9.43 e+08_{5.96 e+09}^{99.42 e+08}$ | $2.64 e+08_{3.18 e+08}^{2.64+08}$ | $6.79 e+08_{5.64 e+09}^{6.78 e+08}$ | $3.62 e+08_{1.53 e+09}^{8.62 e+08}$ |
| 0247 | 3 | $17.15 \pm 0.12$ | $16.89 \pm 0.10$ | $4.84 e+08_{1.84 e+09}^{1.43+08}$ | $7.07 e+07_{1.26 e+09}^{6.00 e+09}$ | $4.13 e+08_{1.16 e+09}^{4.82 e+07}$ | $9.95 e+0777_{7.34 e+09}^{9.95}$ |
| 0252 | 25 | $17.95 \pm 0.07$ | $17.96 \pm 0.10$ | $4.94 e+07_{4.24 e+08}^{4.94 e+07}$ | $3.28 e+07_{4.14 e+07}^{2.15 e+07}$ | $1.66 e+07_{4.01 e+08}^{1.66 e+07}$ | $4.72 e+07_{7.05 e+08}^{4.72 e+07}$ |

continued

| VCC | $\begin{gathered} \hline \hline \text { Type- } \\ \text { ID } \end{gathered}$ | $\begin{gathered} \hline \mathrm{mag}_{\mathrm{r}, \text { obs }} \\ {[\mathrm{mag}]} \end{gathered}$ | $\operatorname{mag}_{\mathrm{r}, \bmod }$ <br> [mag] | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \hline \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \mathrm{M}_{\text {nonvis }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 0260 | 3 | $16.11 \pm 0.06$ | $16.12 \pm 0.10$ | $1.38 e+09_{1.38 e+09}^{2.28 e+08}$ | $7.34 e+07_{1.01 e+08}^{7.34 e+07}$ | $1.30 e+09_{1.30 e+09}^{1.28 e+08}$ | $1.03 e+08_{9.63 e+08}^{1.03+08}$ |
| 0267 | 5 | $12.92 \pm 0.05$ | $12.94 \pm 0.10$ | $2.79 e+10_{2.80 e+10}^{4.42 e+09}$ | $1.49 e+09_{1.95 e+09}^{1.49 e+09}$ | $2.65 e+10_{2.65 e+10}^{2.47 e+09}$ | $2.10 e+00_{1.18 e+10}^{2.10 e+09}$ |
| 0274 | 1 | $16.72 \pm 0.06$ | $16.75 \pm 0.10$ | $5.88 e+08_{5.88 e+08}^{7.28+07}$ | $3.14 e+07_{5.43 e+07}^{3.05 e+07}$ | $5.56 e+08_{5.57 e+08}^{2.43 e+07}$ | $4.41 e+07_{1.09 e+09}^{4.41 e+07}$ |
| 0275 | 25 | $13.92 \pm 0.05$ | $13.95 \pm 0.10$ | $1.24 e+10_{1.24 e+10}^{1.96 e+9}$ | $6.62 e+08_{8.65 e+08}^{6.62 e+08}$ | $1.17 e+10_{1.17 e+10}^{1.10 e+9}$ | $9.31 e+08_{9.54 e+09}^{9.31 e+08}$ |
| 0282 | 25 | $16.93 \pm 0.07$ | $16.94 \pm 0.10$ | $7.26 e+08_{7.27 e+08}^{1.124+08}$ | $3.88 e+07_{5.05 e+07}^{3+.87 e+07}$ | $6.87 e+08_{6.88 e+07}^{6.4 e+07}$ | $5.45 e+07_{4.95 e+08}^{5}$ |
| 0286 | 3 | $15.94 \pm 0.05$ | $15.96 \pm 0.10$ | $1.74 e+09_{1.75 e+09}^{2.76 e+08}$ | $9.31 e+07_{1.22 e+08}^{99.31 e+07}$ | $1.65 e+09_{1.55 e+09}^{1.54+08}$ | $1.31 e+08_{7.38 e+08}^{1.31+08}$ |
| 0289 | 7 | $13.99 \pm 0.05$ | $14.01 \pm 0.10$ | $9.38 e+09_{9.33 e+09}^{3.70 e+09}$ | $5.01 e+08_{5.40 e+08}^{5.01 e+08}$ | $8.88 e+09_{8.88 e+09}^{3.16 e+09}$ | $7.04 e+08_{3.63 e+09}^{7.04 e}$ |
| 0293 | 25 | $15.26 \pm 0.06$ | $15.24 \pm 0.10$ | $7.49 e+08_{4.79 e+09}^{5.91 e+08}$ | $3.30 e+08_{3.57 e+08}^{2.37 e+8}$ | $4.19 e+08_{4.53 e+09}^{2.34+08}$ | $4.68 e+08_{5.43 e+09}^{4.68+08}$ |
| 0304 | 25 | $15.98 \pm 0.06$ | $16.03 \pm 0.10$ | $1.70 e+09_{1.70 e+09}^{2.68 e+08}$ | $9.09 e+07_{1.18 e+07}^{99.09 e+07}$ | $1.61 e+09_{1.61 e+08}^{1.50 e+08}$ | $1.28 e+08_{1.15 e+09}^{1.28+08}$ |
| 0307 | 7 | $9.45 \pm 0.05$ | $9.45 \pm 0.10$ | $7.72 e+11_{7.72 e+11}^{1.22 e+11}$ | $4.12 e+10_{5.39 e+10}^{4.12 e+10}$ | $7.31 e+11_{7.31 e+11}^{6.84 e}$ | $5.79 e+10_{4.68 e+11}^{5.79 e+10}$ |
| 0309 | 3 | $15.76 \pm 0.06$ | $15.75 \pm 0.10$ | $1.60 e+09_{1.60 e+09}^{3.69++88}$ | $8.53 e+07_{1.14 e+08}^{8.53 e+07}$ | $1.51 e+09_{1.51 e+09}^{2.55+08}$ | $1.20 e+08_{9.47 e+08}^{1.20 e+08}$ |
| 0312 | 19 | $12.35 \pm 0.05$ | $12.35 \pm 0.10$ | $1.58 e+10_{1.93 e+10}^{1.58 e+10}$ | $1.09 e+10_{1.31 e+10}^{1.08 e+10}$ | $4.95 e+09_{6.15 e+09}^{4.95 e+09}$ | $1.58 e+10_{2.10 e+11}^{1.58 e+10}$ |
| 0318 | 7 | $13.74 \pm 0.05$ | $13.75 \pm 0.10$ | $3.88 e+09_{9.85 e+09}^{3.88+09}$ | $5.67 e+08_{5.67 e+08}^{5.25+08}$ | $3.31 e+09_{9.32 e+09}^{3.31 e+09}$ | $7.99 e+08_{3.87 e+09}^{7.99 e+08}$ |
| 0320 | 3 | $16.03 \pm 0.07$ | $15.96 \pm 0.10$ | $4.22 e+08_{2.34 e+09}^{3.658+08}$ | $1.18 e+08_{1.62 e+08}^{1.18 e+08}$ | $3.04 e+08_{2.21 e+09}^{2.06 e+08}$ | $1.62 e+08_{2.11 e+09}^{1.62 e+08}$ |
| 0322 | 3 | $14.36 \pm 0.05$ | $14.35 \pm 0.10$ | $8.21 e+09_{8.21 e+09}^{1.30 e+09}$ | $4.38 e+08_{5.72 e+08}^{4.38+08}$ | $7.78 e+09_{7.78 e+09}^{7.26 e+08}$ | $6.16 e+08_{4.80 e+09}^{6.16 e+08}$ |
| 0323 | 17 | $13.52 \pm 0.05$ | $13.49 \pm 0.10$ | $4.41 e+09_{5.54 e+09}^{4.41 e+09}$ | $3.03 e+09_{3.78 e+09}^{3.03 e+09}$ | $1.38 e+09_{1.76 e+09}^{1.38+09}$ | $4.39 e+09_{5.99 e+10}^{4.39 e+09}$ |
| 0324 | 1 | $13.64 \pm 0.05$ | $13.65 \pm 0.10$ | $1.13 e+10_{1.13 e+10}^{4.44 e+99}$ | $6.01 e+08_{6.49 e+08}^{6.01 e+08}$ | $1.07 e+10_{1.07 e+10}^{3.79 e+09}$ | $8.45 e+08_{3.45 e+09}^{8.45 e+08}$ |
| 0328 | 3 | $15.81 \pm 0.05$ | $15.85 \pm 0.10$ | $1.52 e+09_{1.52 e+09}^{3.35+08}$ | $8.10 e+07_{1.04 e+08}^{8.10 e+07}$ | $1.44 e+09_{1.44 e+09}^{2.31 e+08}$ | $1.14 e+08_{8.83 e+08}^{1.14 e+08}$ |
| 0329 | 3 | $17.56 \pm 0.08$ | $17.54 \pm 0.10$ | $2.72 e+08_{2.72 e+08}^{8.82 e+07}$ | $1.45 e+07_{1.70 e+07}^{1.37++07}$ | $2.58 e+08_{2.58 e+08}^{1.52 e+07}$ | $2.04 e+07_{1.47 e+08}^{7.04+07}$ |
| 0334 | 1 | $15.08 \pm 0.05$ | $15.08 \pm 0.10$ | $3.13 e+09_{3.13 e+09}^{3.13 e+09}$ | $1.67 e+08_{1.67 e+08}^{1.67+08}$ | $2.96 e+09^{2} 2.966 e+09$ | $2.34 e+08_{7.03 e+08}^{2.34 e+08}$ |

continued

| VCC | $\begin{gathered} \hline \hline \text { Type- } \\ \text { ID } \end{gathered}$ | mag $_{\mathrm{r}, \mathrm{obs}}$ <br> [mag] | $\operatorname{mag}_{\mathrm{r}, \mathrm{mod}}$ [mag] | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \mathrm{M}_{\mathrm{nonvis}} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 0340 | 1 | $14.16 \pm 0.05$ | $14.16 \pm 0.10$ | $9.16 e+09_{9.16 e+09}^{9.16 e+09}$ | $4.89 e+08_{4.89 e+08}^{4.89 e+08}$ | $8.67 e+09_{8.67 e+09}^{8.67+09}$ | $6.87 e+08_{2.06 e+09}^{6.87 e+08}$ |
| 0341 | 17 | $11.35 \pm 0.05$ | $11.34 \pm 0.10$ | $4.01 e+10_{4.89 e+10}^{4.01 e+10}$ | $2.75 e+10_{3.33 e+10}^{2.75 e+10}$ | $1.26 e+10_{1.56 e+10}^{1.26 e+10}$ | $4.01 e+10_{5.81 e+11}^{4.01 e+10}$ |
| 0342 | 19 | $13.36 \pm 0.05$ | $13.35 \pm 0.10$ | $5.74 e+09_{6.99 e+09}^{5.73+09}$ | $3.94 e+09_{4.76 e+09}^{3.94+09}$ | $1.80 e+09_{2.23 e+09}^{1.80 e+09}$ | $5.74 e+09_{7.62 e+10}^{5.74 e+10}$ |
| 0343 | 9 | $14.73 \pm 0.05$ | $14.74 \pm 0.10$ | $5.66 e+09_{5.66 e+09}^{8.94 t+08}$ | $3.02 e+08_{3.94 e+08}^{3.02 e+08}$ | $5.35 e+09_{5.36 e+09}^{5.00 e+08}$ | $4.24 e+08_{3.31 e+09}^{4.24++08}$ |
| 0350 | 3 | $16.03 \pm 0.07$ | $16.06 \pm 0.10$ | $1.47 e+09_{1.48 e+09}^{2.40 e+08}$ | $7.87 e+07_{1.06 e+08}^{7.87 e+07}$ | $1.40 e+09_{1.40 e+09}^{1.34+08}$ | $1.11 e+08_{9.01 e+08}^{1.11 e+08}$ |
| 0354 | 25 | $14.89 \pm 0.06$ | $14.89 \pm 0.10$ | $8.72 e+08_{7.75 e+09}^{8.72 e+08}$ | $5.78 e+08_{8.96 e+08}^{4.13 e+08}$ | $2.93 e+08_{7.34 e+09}^{2.93 e+08}$ | $8.32 e+08_{1.65 e+10}^{8.32 e+08}$ |
| 0358 | 17 | $12.53 \pm 0.05$ | $12.52 \pm 0.10$ | $1.27 e+10_{1.54 e+10}^{1.18 e+10}$ | $8.69 e+09_{1.05 e+10}^{8.10 e+09}$ | $3.97 e+09_{4.92 e+09}^{3.70 e+09}$ | $1.27 e+10_{1.95 e+11}^{1.27 e+10}$ |
| 0364 | 3 | $16.56 \pm 0.07$ | $16.56 \pm 0.10$ | $2.63 e+08_{1.49 e+09}^{1.83+08}$ | $7.36 e+07_{1.10 e+08}^{7.36+07}$ | $1.89 e+08_{1.41 e+09}^{7.25 e+07}$ | $1.01 e+08_{1.64 e+09}^{1.01 e+08}$ |
| 0366 | 19 | $13.19 \pm 0.05$ | $13.22 \pm 0.10$ |  | $2.68 e+09^{1.822 e+09}$ | $1.36 e+09_{3.44 e+10}^{1.36 e+9}$ | $3.86 e+09_{5.90 e+10}^{3.86 e+09}$ |
| 0367 | 3 | $14.91 \pm 0.07$ | $14.93 \pm 0.10$ | $1.40 e+09^{9} 9.097 e+08$ | $3.91 e+08_{8.37 e+08}^{3.91 e+08}$ | $1.01 e+09_{7}^{3.55 e++08}$ | $5.37 e+08_{1.38 e+10}^{5.37 e+10}$ |
| 0371 | 19 | $12.42 \pm 0.05$ | $12.43 \pm 0.10$ | $1.55 e+10_{1.89 e+10}^{1.55 e+10}$ | $1.07 e+10_{1.29 e+10}^{1.07 e+10}$ | $4.87 e+09_{6.04 e+09}^{4.87 e+09}$ | $1.55 e+10_{2.25 e+11}^{1.55 e+10}$ |
| 0373 | 19 | $12.26 \pm 0.05$ | $12.27 \pm 0.10$ | $1.62 e+10_{1.80 e+10}^{1.38+10}$ | $1.11 e+10_{1.23 e+10}^{99.43+09}$ | $5.12 e+09_{5.74 e+09}^{4.32 e+09}$ | $1.62 e+10_{2.57 e+11}^{1.62 e+10}$ |
| 0375 | 19 | $11.85 \pm 0.05$ | $11.84 \pm 0.10$ | $2.24 e+10_{2.73 e+10}^{2.24 e+10}$ | $1.54 e+10_{1.86 e+10}^{1.53 e+10}$ | $7.01 e+09_{8.69 e+09}^{7.00 e+09}$ | $2.24 e+10_{2.97 e+11}^{2.24 e+10}$ |
| 0379 | 3 | $0.00 \pm-99999.00$ | $198.00 \pm 99.00$ | $0.00 e+00_{0.00 e++00}^{0.000+00}$ | $0.00 e+00_{0.00 e++00}^{0.000+00}$ | $0.00 e+00_{0.00 e++00}^{0.000 e+00}$ | $6.95 e-310_{6.95 e-310}^{6.955-310}$ |
| 0381 | 3 | $16.19 \pm 0.09$ | $16.07 \pm 0.10$ | $1.42 e+09^{2}+1.40 e+088$ | $7.60 e+07_{1.50 e+08}^{7.32 e++7}$ | $1.35 e+09_{1.356++09}^{99.86 e+07}$ | $1.07 e+08_{1.53 e+09}^{1.0708}$ |
| 0382 | 7 | $11.52 \pm 0.05$ | $11.54 \pm 0.10$ | $1.08 e+11_{1.088 e+11}^{1.08 e+11}$ | $5.78 e+09_{5.78 e+09}^{5.77 e+09}$ | $1.02 e+11_{1.02 e+11}^{1.02 e+11}$ | $8.12 e+09_{2.44 e+10}^{8.12 e+09}$ |
| 0386 | 17 | $13.13 \pm 0.05$ | $13.12 \pm 0.10$ | $7.21 e+09_{8.02 e+09}^{6.58 e+09}$ | $4.93 e+09_{5.46 e+09}^{4.52 e+09}$ | $2.28 e+09_{2.56 e+09}^{2.06+09}$ | $7.21 e+09_{8.80 e+10}^{7.21 e+09}$ |
| 0393 | 7 | $12.71 \pm 0.05$ | $12.75 \pm 0.10$ | $4.21 e+10{ }_{4.21 e+10}^{5.31 e+09}$ | $2.25 e+09_{3.20 e+09}^{2.25 e+09}$ | $3.99 e+10_{3.99 e+10}^{2.11 e+09}$ | $3.16 e+09_{3.38 e+10}^{3.16 e+09}$ |
| 0404 | 7 | $13.96 \pm 0.05$ | $13.94 \pm 0.10$ | $2.58 e+09^{2.53 l e+10} 2$ | $1.71 e+09_{1.73 e+09}^{1.23+09}$ | $8.68 e+08_{2.18 e+10}^{7.88+08}$ | $2.46 e+09_{1.58 e+10}^{2.46+09}$ |
| 0408 | 19 | $10.94 \pm 0.05$ | $10.93 \pm 0.10$ | $6.30 e+10_{7.68 e+10}^{6.30 e+10}$ | $4.33 e+10_{5.23 e+10}^{4.32 e+10}$ | $1.97 e+10_{2.45 e+10}^{1.97 e+10}$ | $6.30 e+10_{9.13 e+11}^{6.30 e+10}$ |

continued

| VCC | $\begin{gathered} \hline \hline \text { Type- } \\ \text { ID } \end{gathered}$ | $\begin{gathered} \hline \mathrm{mag}_{\mathrm{r}, \text { obs }} \\ {[\mathrm{mag}]} \end{gathered}$ | $\operatorname{mag}_{\mathrm{r}, \bmod }$ <br> [mag] | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \hline \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \mathrm{M}_{\text {nonvis }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 0409 | 25 | $17.57 \pm 0.09$ | $17.61 \pm 0.10$ | $3.78 e+08_{3.78 e+08}^{5.95+07}$ | $2.02 e+07_{2.62 e+07}^{2.02 e+07}$ | $3.58 e+08_{3.58 e+08}^{3.33++07}$ | $2.84 e+07_{2.60 e+08}^{2.84 e+07}$ |
| 0410 | 1 | $16.68 \pm 0.05$ | $16.75 \pm 0.10$ | $7.97 e+07_{5.83 e+08}^{5.70 e+07}$ | $2.82 e+07_{3.57 e+07}^{2.53 e+07}$ | $5.15 e+07_{5.52 e+08}^{2.89+07}$ | $3.91 e+07_{6.17 e+08}^{7.91 e+07}$ |
| 0413 | 3 | $16.94 \pm 0.06$ | $16.94 \pm 0.10$ | $8.10 e+08_{8.10 e+08}^{1.11 e+08}$ | $4.32 e+07_{6.73 e+07}^{4.32 e+07}$ | $7.66 e+08_{7.66 e+08}^{4.42 e+07}$ | $6.07 e+07_{7.42 e+08}^{6.07 e+07}$ |
| 0414 | 25 | $16.68 \pm 0.07$ | $16.69 \pm 0.10$ | $1.58 e+08_{1.41 e+09}^{1.58+08}$ | $1.05 e+08_{1.05 e+08}^{7.52 e+07}$ | $5.33 e+07_{1.33 e+09}^{5.33+07}$ | $1.51 e+08_{1.33 e+09}^{1.51 e+08}$ |
| 0415 | 9 | $14.53 \pm 0.05$ | $14.53 \pm 0.10$ | $7.55 e+09_{7}^{1.15 e++09}$ | $4.03 e+08_{5.26 e+08}^{4.032+08}$ | $7.14 e+09_{7.15 e+09}^{6.667+08}$ | $5.66 e+08_{5.10 e+09}^{5.566+08}$ |
| 0423 | 3 | $15.31 \pm 0.12$ | $15.04 \pm 0.10$ | $1.52 e+09_{8.55 e+09}^{99.89+08}$ | $4.26 e+08_{9.28 e+08}^{4.26 e+08}$ | $1.10 e+09_{8.10 e+09}^{3.16 e+08}$ | $5.85 e+08_{1.74 e+10}^{5.85 e+08}$ |
| 0425 | 3 | $17.42 \pm 0.11$ | $17.49 \pm 0.10$ | $2.49 e+08_{2.27 e+09}^{2.35+08}$ | $1.71 e+08_{2.06 e+08}^{1.21 e+08}$ | $7.80 e+07_{2.15 e+09}^{7.38+07}$ | $2.49 e+08_{5.13 e+09}^{2.49 e+08}$ |
| 0428 | 1 | $16.68 \pm 0.05$ | $16.79 \pm 0.10$ | $5.51 e+07_{6.88 e+07}^{4.27 e+07}$ | $1.85 e+07_{3.28 e+07}^{1.85 e+07}$ | $3.66 e+07_{4.58 e+07}^{1.37 e+07}$ | $2.57 e+07_{2.64 e+08}^{7.57 e}$ |
| 0429 | 29 | $15.99 \pm 0.05$ | $16.02 \pm 0.10$ | $3.58 e+08_{7.83 e+08}^{2.288+08}$ | $1.00 e+08_{1.52 e+08}^{1.00+08}$ | $2.58 e+08_{6.68 e+08}^{7.61 e+07}$ | $1.37 e+08_{2.62 e+09}^{1.37 e}$ |
| 0446 | 29 | $15.15 \pm 0.05$ | $15.17 \pm 0.10$ | $6.52 e+08_{4.13 e+09}^{6.52 e+08}$ | $2.87 e+08_{2.87 e+08}^{2.22 e+08}$ | $3.65 e+08_{3.99 e+08}^{3.65 e+08}$ | $4.08 e+08_{2.93 e+09}^{4.088}$ |
| 0448 | 3 | $16.36 \pm 0.05$ | $16.39 \pm 0.10$ | $9.54 e+08_{9.54 e+08}^{2.04+08}$ | $5.09 e+07_{6.33 e+07}^{5.09++07}$ | $9.04 e+08_{9.04 e+08}^{1.41 e+08}$ | $7.15 e+07_{5.50 e+08}^{7.15+07}$ |
| 0449 | 5 | $13.34 \pm 0.05$ | $13.34 \pm 0.10$ | $8.45 e+09_{1.12 e+10}^{66.64+09}$ | $3.73 e+09_{4.01+e+09}^{3.46 e+09}$ | $4.73 e+09_{7.71 e+09}^{2.63 e+09}$ | $5.29 e+09_{3.14 e+10}^{5.29 e+09}$ |
| 0450 | 19 | $14.76 \pm 0.05$ | $14.76 \pm 0.10$ | $1.10 e+09_{9.81 e+09}^{1.10 e+09}$ | $7.30 e+08_{7.30 e+08}^{5.24+08}$ | $3.70 e+08_{9.29 e+09}^{3.70 e+08}$ | $1.05 e+09_{3.89 e+09}^{1.05 e+09}$ |
| 0453 | 11 | $15.73 \pm 0.05$ | $15.77 \pm 0.10$ | $2.59 e+09_{2.59 e+09}^{3.28 e+08}$ | $1.38 e+08_{1.98 e+08}^{1.38+08}$ | $2.45 e+09_{2.45 e+09}^{1.30 e+08}$ | $1.94 e+08_{2.08 e+09}^{1.94+08}$ |
| 0459 | 1 | $14.37 \pm 0.05$ | $14.38 \pm 0.10$ | $5.60 e+09_{5.60 e+09}^{5.600+09}$ | $2.99 e+08_{2.99 e+08}^{2.99 e+08}$ | $5.30 e+09_{5.30 e+09}^{5.30 e+09}$ | $4.20 e+08_{1.26 e+09}^{4.20 e+08}$ |
| 0460 | 17 | $10.11 \pm 0.05$ | $10.05 \pm 0.10$ | $1.07 e+11_{1.30 e++11}^{1.00+11}$ | $7.35 e+10_{8.89 e+10}^{6.86 e+10}$ | $3.35 e+10_{4.16 e+10}^{3.13 e+10}$ | $1.07 e+11_{1.62 e+12}^{1.07 e+11}$ |
| 0465 | 7 | $11.82 \pm 0.05$ | $11.83 \pm 0.10$ | $6.09 e+10_{6.09 e+10}^{2.40 e+10}$ | $3.25 e+09_{3.51+e+09}^{3.25 e+09}$ | $5.76 e+10_{5.76 e+10}^{2.05 e+10}$ | $4.56 e+09_{2.36 e+10}^{4.56 e+09}$ |
| 0472 | 25 | $16.42 \pm 0.06$ | $16.37 \pm 0.10$ | $2.59 e+08_{2.012 e+09}^{2.25 e+08}$ | $1.78 e+08_{2}^{1.57 e++08}$ |  | $2.57 e+08_{5.87 e+09}^{2.57 e+08}$ |
| 0476 | 3 | $17.23 \pm 0.09$ | $17.27 \pm 0.10$ | $5.29 e+08_{5.30 e+08}^{8.22 e+07}$ | $2.82 e+07_{6.27 e+07}^{\substack{2 \\ \hline 20}}$ | $5.01 e+08_{5.01 e+08}^{\substack{3 \\ 5}}$ | $3.97 e+07_{7}^{7.977 e+07}$ |
| 0477 | 3 | $16.51 \pm 0.08$ | $16.51 \pm 0.10$ | $1.94 e+08_{1.05 e+09}^{1.66+08}$ | $5.43 e+07_{7.32 e+07}^{5.43 e+07}$ | $1.40 e+08_{9.96 e+08}^{9} 9 .+20^{2}$ | $7.45 e+07_{8.75 e+08}^{7.45 e+07}$ |

continued

| VCC | Type- <br> ID | mag $_{\text {r,obs }}$ <br> [mag] | $\operatorname{mag}_{r, \bmod }$ <br> [mag] | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \hline \mathrm{M}_{\text {nonvis }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 0479 | 3 | $15.84 \pm 0.05$ | $15.89 \pm 0.10$ | $1.87 e+09_{1.87 e+09}^{2.95 e+08}$ | $9.97 e+07_{1.30 e+08}^{9.97 e+07}$ | $1.77 e+09_{1.77 e+09}^{1.65 e+08}$ | $1.40 e+08_{1.44 e+09}^{1.40 e+08}$ |
| 0483 | 7 | $0.00 \pm-99999.00$ | $198.00 \pm 99.00$ | $0.00 e+00_{0.00 e+00}^{0.00 e+00}$ | $0.00 e+00_{0.00 e+00}^{0.00 e+00}$ | $0.00 e+00_{0.00 e++00}^{0.00 e+00}$ | $6.95 e-310_{6.95 e-310}^{6.95 e-310}$ |
| 0488 | 3 | $16.72 \pm 0.06$ | $16.81 \pm 0.10$ | $8.75 e+08_{8.75 e+08}^{1.25 e+08}$ | $4.67 e+07_{7.53 e+07}^{4.67 e+07}$ | $8.28 e+08_{8.28 e+08}^{4.95+07}$ | $6.56 e+07_{8.87 e+08}^{6.56 e+07}$ |
| 0491 | 7 | $12.16 \pm 0.05$ | $12.18 \pm 0.10$ | $3.81 e+10_{3.81 e+10}^{6.47 e+09}$ | $2.03 e+09_{2.04 e+09}^{1.81 e+09}$ | $3.61 e+10_{3.61 e+10}^{4.66+09}$ | $2.86 e+09_{1.36 e+10}^{2.86 e+9}$ |
| 0494 | 25 | $15.08 \pm 0.06$ | $15.02 \pm 0.10$ | $5.68 e+09_{5.68 e+09}^{7.02 e+08}$ | $3.03 e+08_{7.46 e+08}^{2.90 e+08}$ | $5.38 e+09_{5.38 e+09}^{2.54 e+08}$ | $4.26 e+08_{1.18 e+10}^{4.26 e+08}$ |
| 0497 | 7 | $0.00 \pm-99999.00$ | $198.00 \pm 99.00$ | $0.00 e+00_{0.00 e+00}^{0.0 .0 e+00}$ | $0.00 e+00_{0.00 e+00}^{0.00 e+00}$ | $0.00 e+00_{0.00 e++00}^{0.000+00}$ | $6.95 e-310_{6.95 e-310}^{6.959-310}$ |
| 0509 | 9 | $14.80 \pm 0.05$ | $14.83 \pm 0.10$ | $8.99 e+08_{4.28 e+09}^{8.99 e+08}$ | $2.78 e+08_{2.78 e+08}^{2.29 e+08}$ | $6.20 e+08_{4.06 e+09}^{6.20 e+08}$ | $3.94 e+08_{2.20 e+09}^{3.94 e+08}$ |
| 0512 | 11 | $15.13 \pm 0.05$ | $15.10 \pm 0.10$ | $3.42 e+09_{3.42 e+09}^{5.76+08}$ | $1.82 e+08_{2.54 e+08}^{1.82 e+08}$ | $3.23 e+09_{3.24 e+09}^{3.22 e+08}$ | $2.56 e+08_{2.312+09}^{2.56+08}$ |
| 0513 | 1 | $14.62 \pm 0.05$ | $14.67 \pm 0.10$ | $8.56 e+09_{8.56 e+09}^{1.04+09}$ | $4.57 e+08_{6.99 e+08}^{4.57 e+08}$ | $8.10 e+09_{8.11 e+09}^{3.50 e+08}$ | $6.42 e+08_{1.112 e+10}^{6.42 e+10}$ |
| 0514 | 7 | $14.12 \pm 0.05$ | $14.11 \pm 0.10$ | $2.25 e+09_{4.23 e+09}^{1.70 e+09}$ | $6.97 e+08_{7.49 e+08}^{5.41 e+08}$ | $1.55 e+09_{3.62 e+09}^{9.51 e+08}$ | $9.85 e+08_{7.44 e+09}^{99.85+08}$ |
| 0520 | 3 | $17.86 \pm 0.11$ | $17.86 \pm 0.10$ | $2.55 e+08_{2.55 e+07}^{4.72 e+07}$ | $1.36 e+07_{1.66 e+07}^{1.32 e+07}$ | $2.41 e+08_{2.41 e+08}^{3.39+07}$ | $1.91 e+07_{1.64 e+07}^{1.92 e+08}$ |
| 0522 | 17 | $12.22 \pm 0.05$ | $12.23 \pm 0.10$ | $1.18 e+10_{1.59 e+10}^{1.18 e+10}$ | $8.10 e+09_{1.09 e+10}^{7.89 e+09}$ | $3.70 e+09_{5.04 e+09}^{3.70 e+09}$ | $1.17 e+10_{1.80 e+11}^{1.17 e+10}$ |
| 0524 | 5 | $11.67 \pm 0.05$ | $11.65 \pm 0.10$ | $3.13 e+10_{3.67 e+10}^{3.13 e+10}$ | $2.15 e+10_{2.50 e+10}^{2.15 e+10}$ | $9.81 e+09_{1.17 e+10}^{9.81+09}$ | $3.13 e+10_{3.78 e+11}^{3.13 e+10}$ |
| 0530 | 3 | $14.76 \pm 0.06$ | $14.72 \pm 0.10$ | $1.09 e+09_{5.21 e+09}^{8.32 e+08}$ | $3.38 e+08_{3.67 e+08}^{2.78 e+08}$ | $7.52 e+08_{4.93 e+09}^{4.66 e+08}$ | $4.77 e+08_{3.688+09}^{4.77 e+08}$ |
| 0534 | 17 | $12.49 \pm 0.05$ | $12.47 \pm 0.10$ | $1.24 e+10_{1.33 e+10}^{1.12 e+10}$ | $8.51 e+09_{9}^{7} 7.789++099$ | $3.93 e+09_{4.23 e+09}^{3.56+09}$ | $1.24 e+10_{1.38 e+11}^{1.24 e+10}$ |
| 0546 | 25 | $15.00 \pm 0.05$ | $15.02 \pm 0.10$ | $5.29 e+09_{5.29 e+09}^{6.35 z+08}$ | $2.82 e+08_{3.94 e+08}^{2.82 e+08}$ | $5.01 e+09_{5.01 e+09}^{2.59 e+08}$ | $3.97 e+08_{4.22 e+09}^{3.97 e+08}$ |
| 0559 | 17 | $11.49 \pm 0.05$ | $11.47 \pm 0.10$ | $2.36 e+10_{2.12 e+11}^{2.36 e+10}$ | $1.62 e+10_{1.90 e+10}^{1.13 e+10}$ | $7.40 e+09_{2.01 e+11}^{7.40 e+09}$ | $2.35 e+10_{2.48 e+11}^{2.35 e+10}$ |
| 0562 | 1 | $15.56 \pm 0.05$ | $15.60 \pm 0.10$ | $2.72 e+08_{2.72 e+08}^{2.09 e+08}$ | $7.60 e+07_{1.05 e+08}^{7.60 e+07}$ | $1.96 e+08_{1.96 e+08}^{1.06+08}$ | $1.04 e+08_{1.36 e+09}^{1.04 t+08}$ |
| 0566 | 11 | $15.97 \pm 0.05$ | $16.03 \pm 0.10$ | $1.79 e+08_{1.13 e+09}^{1.79 e+08}$ | $5.00 e+07_{6.01 e+07}^{5.007}$ | $1.29 e+08_{1.07 e+09}^{1.29 e+08}$ | $6.86 e+07_{3.75 e+08}^{6.86+07}$ |
| 0567 | 7 | $13.51 \pm 0.05$ | $13.52 \pm 0.10$ | $2.08 e+10_{2.08 e+10}^{2.59 e+09}$ | $1.11 e+09_{1.57 e+09}^{1.11 e+09}$ | $1.97 e+10_{1.97 e+10}^{1.033+09}$ | $1.56 e+09_{9.15 e+09}^{1.56 e+09}$ |

continued

| VCC | $\begin{gathered} \hline \hline \text { Type- } \\ \text { ID } \end{gathered}$ | mag $_{\mathrm{r}, \mathrm{obs}}$ <br> [mag] | $\operatorname{mag}_{\mathrm{r}, \bmod }$ [mag] | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \mathrm{M}_{\text {nonvis }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 0570 | 17 | $11.29 \pm 0.05$ | $11.26 \pm 0.10$ | $3.93 e+10_{4.79 e+10}^{3.93 e+10}$ | $2.70 e+10_{3.26 e+10}^{2.70 e+10}$ | $1.23 e+10_{1.53 e+10}^{1.23 e+10}$ | $3.93 e+10_{5.22 e+11}^{3.93 e+10}$ |
| 0574 | 25 | $16.48 \pm 0.09$ | $16.49 \pm 0.10$ | $3.25 e+08_{7.18 e+08}^{2.83+08}$ | $9.10 e+07_{1.25 e+08}^{99.10 e+07}$ | $2.34 e+08_{6.13 e+08}^{1.58+08}$ | $1.25 e+08_{1.21 e+09}^{1.25 e+08}$ |
| 0576 | 5 | $12.46 \pm 0.05$ | $12.44 \pm 0.10$ | $1.36 e+10_{1.22 e+10}^{1.36 e+11}$ | $9.00 e+09_{9}^{6.463 e+09}{ }^{\text {b }}$ | $4.56 e+09_{1.15 e+11}^{4.56+09}$ | $1.29 e+10_{8.28 e+10}^{1.29 e+10}$ |
| 0583 | 3 | $14.77 \pm 0.05$ | $14.74 \pm 0.10$ | $7.47 e+09_{7.47 e+09}^{9,23 e+08}$ | $3.98 e+08_{5.57 e+08}^{3.98+08}$ | $7.07 e+09_{7.07 e+09}^{3.66+08}$ | $5.60 e+08_{7.34 e+09}^{5.60 e+08}$ |
| 0584 | 3 | $15.56 \pm 0.07$ | $15.61 \pm 0.10$ | $3.77 e+08_{3.03 e+09}^{3.77 e+08}$ | $2.28 e+08_{2.86 e+08}^{1.53++08}$ | $1.50 e+08_{2.87 e+09}^{1.45 e+08}$ | $3.26 e+08_{3.90 e+09}^{3.26 e+08}$ |
| 0585 | 3 | $16.73 \pm 0.12$ | $16.76 \pm 0.10$ | $6.95 e+08_{6.96 e+08}^{1.47 e+08}$ | $3.71 e+07_{4.54 e+07}^{3.71 e+07}$ | $6.58 e+08_{6.59 e+08}^{1.01 e+08}$ | $5.22 e+07_{4.00 e+08}^{5.22 e+07}$ |
| 0593 | 23 | $14.43 \pm 0.05$ | $14.45 \pm 0.10$ | $7.98 e+09_{7.98 e+09}^{1.26 e+09}$ | $4.26 e+08_{5.56 e+08}^{4.26 e+08}$ | $7.55 e+09_{7.55 e+09}^{7.05+08}$ | $5.98 e+08_{3.87 e+09}^{5.98 e}$ |
| 0596 | 7 | $9.02 \pm 0.05$ | $9.03 \pm 0.10$ | $1.50 e+12_{1.50 e+12}^{1.85 e+11}$ | $8.01 e+10_{1.12 e+11}^{8.01 e+10}$ | $1.42 e+12_{1.42 e+12}^{7.36 e+10}$ | $1.13 e+11_{6.58 e+11}^{1.13 e+11}$ |
| 0618 | 3 | $16.24 \pm 0.06$ | $16.24 \pm 0.10$ | $9.21 e+08_{9.21 e+08}^{3.84+08}$ | $4.91 e+07_{5.61 e+07}^{4.92 e+07}$ | $8.72 e+08_{8.72 e+08}^{3.28+08}$ | $6.90 e+07_{3.650++07}^{6.902}$ |
| 0620 | 11 | $15.36 \pm 0.05$ | $15.37 \pm 0.10$ | $2.61 e+09_{2.612++09}^{1.03 e+09}$ | $1.39 e+08_{1.51++08}^{1.39 e+08}$ | $2.47 e+09_{2.47 e+09}^{8.80 e+08}$ | $1.96 e+08_{1.01 e+09}^{1.96+08}$ |
| 0630 | 9 | $12.23 \pm 0.05$ | $12.23 \pm 0.10$ | $9.62 e+09_{9.632 e+09}^{9.62 e+09}$ | $6.39 e+09_{6.39 e+09}^{6.39+09}$ | $3.24 e+09_{3.24 e+09}^{3.24 e+09}$ | $9.18 e+09^{9.76 e++10}$ |
| 0633 | 25 | $16.47 \pm 0.09$ | $16.42 \pm 0.10$ | $2.03 e+08_{1.64 e+09}^{2.02 e+08}$ | $1.35 e+08_{2.20 e+08}^{8.40 e+07}$ | $6.83 e+07_{1.55 e+09}^{6.83+07}$ | $1.94 e+08_{3.70 e+09}^{1.94+08}$ |
| 0641 | 1 | $15.59 \pm 0.05$ | $15.66 \pm 0.10$ | $2.04 e+09_{2.04 e+09}^{2.33 e+08}$ | $1.09 e+08_{1.56 e+08}^{1.03++08}$ | $1.93 e+09^{7.799 e+07}$ | $1.53 e+08_{3.04 e+09}^{1.53 e+08}$ |
| 0651 | 25 | $16.38 \pm 0.09$ | $16.24 \pm 0.10$ | $1.59 e+09_{1.59 e+09}^{2.12 e+08}$ | $8.48 e+07_{2.65 e+08}^{8.11 e+07}$ |  | $1.19 e+08_{3.96 e+09}^{1.19+08}$ |
| 0654 | 19 | $10.91 \pm 0.05$ | $10.89 \pm 0.10$ | $6.59 e+10_{8.03 e+10}^{6.13 e+10}$ | $4.52 e+10_{5.47 e+10}^{4.22 e+10}$ | $2.06 e+10_{2.56 e+10}^{1.922+10}$ | $6.59 e+10_{1.08 e+12}^{6599+10}$ |
| 0655 | 1 | $12.37 \pm 0.05$ | $12.39 \pm 0.10$ | $7.09 e+10_{7.09 e+10}^{8.51 e+09}$ | $3.78 e+09_{5.65 e+09}^{3.78+09}$ | $6.71 e+10_{6.71 e+10}^{2.86 e+09}$ | $5.32 e+09_{5.88 e+10}^{5.32 e+09}$ |
| 0656 | 5 | $11.99 \pm 0.05$ | $11.96 \pm 0.10$ | $2.15 e+10_{2.70 e+10}^{2.15 e+10}$ | $1.48 e+10_{1.84 e+10}^{1.48 e+10}$ | $6.74 e+09_{8.59 e+09}^{6.74 e+09}$ | $2.14 e+10_{2.92 e+11}^{2.14 e+10}$ |
| 0657 | 19 | $12.10 \pm 0.05$ | $12.11 \pm 0.10$ | $2.32 e+10_{2.83 e+10}^{2.16 e+10}$ | $1.60 e+10_{1.93 e+10}^{1.49 e+10}$ | $7.29 e+09^{9.704 e+09}$ | $2.32 e+10_{3.59 e+11}^{2.32 e+10}$ |
| 0664 | 7 | $13.10 \pm 0.05$ | $13.10 \pm 0.10$ | $1.63 e+10_{1.63 e+10}^{2.93 e+09}$ | $8.72 e+08_{8.72 e+08}^{8.22 e+08}$ | $1.55 e+10_{1.55 e+10}^{2.11 e+09}$ | $1.22 e+09_{4.80 e+09}^{1.22 e+09}$ |
| 0666 | 3 | $16.03 \pm 0.07$ | $16.01 \pm 0.10$ | $1.92 e+09_{1.92 e+09}^{3.022+08}$ | $1.02 e+08_{1.33 e+08}^{1.02 e+08}$ | $1.82 e+09_{1.82 e+09}^{1.69+08}$ | $1.44 e+08_{1.47 e+09}^{1.44 e+08}$ |

continued

| VCC | Type- <br> ID | mag $_{\mathrm{r}, \text { obs }}$ <br> [mag] | $\operatorname{mag}_{r, \bmod }$ <br> [mag] | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \hline \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \hline \mathrm{M}_{\text {nonvis }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 0667 | 7 | $13.20 \pm 0.05$ | $13.19 \pm 0.10$ | $4.01 e+09_{3.58 e+10}^{4.01 e+09}$ | $2.66 e+09_{2.666+09}^{1.91 e+09}$ | $1.35 e+09_{3.39 e+10}^{1.35 e+09}$ | $3.82 e+09_{1.42 e+10}^{3.82 e+09}$ |
| 0672 | 19 | $13.11 \pm 0.05$ | $13.12 \pm 0.10$ | $5.19 e+09_{4.48 e+10}^{5.00 e+09}$ | $3.56 e+09_{4.60 e+09}^{2.39 e+09}$ | $1.63 e+09_{4.24 e+10}^{1.63++9}$ | $5.16 e+09_{7.41++10}^{5.16 e+09}$ |
| 0679 | 11 | $15.17 \pm 0.05$ | $15.21 \pm 0.10$ | $3.96 e+09_{3.96 e+09}^{5.43 e+08}$ | $2.11 e+08_{3.28 e+08}^{2.11 e+08}$ | $3.75 e+09_{3.75 e+09}^{2.16 e+08}$ | $2.97 e+08_{3.25 e+09}^{2.97 e+08}$ |
| 0688 | 7 | $13.16 \pm 0.05$ | $13.14 \pm 0.10$ | $4.16 e+09_{3.64 e+109}^{4.166+10}$ | $2.76 e+09_{2.76 e+09}^{1.94+09}$ | $1.40 e+09_{3.44 e+109}^{1.46 e+10}$ | $3.97 e+09^{3.571 e+10}$ |
| 0692 | 7 | $12.21 \pm 0.05$ | $12.24 \pm 0.10$ | $5.56 e+10_{5.56 e+10}^{8.80+0}$ | $2.97 e+09_{3.88 e+09}^{2.97 e+09}$ | $5.27 e+10_{5.27 e+10}^{4.92 e+09}$ | $4.17 e+09_{3.76 e+10}^{4.17 e+09}$ |
| 0697 | 7 | $13.08 \pm 0.05$ | $13.11 \pm 0.10$ | $3.46 e+10_{3.46 e+10}^{4.27 e+09}$ | $1.85 e+09_{2.58 e+09}^{1.85 e+09}$ | $3.27 e+10_{3.28 e+10}^{1.70 e+9}$ | $2.59 e+09_{2.76 e+10}^{2.59 e+09}$ |
| 0699 | 7 | $12.99 \pm 0.05$ | $12.99 \pm 0.10$ | $2.23 e+10_{2.23 e+10}^{4.68+09}$ | $1.19 e+09_{1.45 c+09}^{1.19 e+09}$ | $2.11 e+10_{2.11 e+10}^{3.23 e+09}$ | $1.67 e+09_{9.11 l^{1.67 e+09}}^{2+09}$ |
| 0703 | 3 | $17.00 \pm 0.07$ | $16.94 \pm 0.10$ | $2.94 e+08_{7.43 e+08}^{1.34+08}$ | $4.30 e+07_{4.79 e+07}^{3.75 e+07}$ | $2.51 e+08_{7.033+07}^{9.64+07}$ | $6.05 e+07_{5.31++08}^{6.05 e+07}$ |
| 0705 | 25 | $16.09 \pm 0.07$ | $16.08 \pm 0.10$ | $1.05 e+09_{2.67 e+09}^{3.01 e+08}$ | $1.54 e+08_{2.84 e+08}^{1.33+08}$ | $8.98 e+08_{2.53 e+09}^{1.01 e+08}$ | $2.16 e+08_{5.10 e+09}^{2.16 e+08}$ |
| 0713 | 7 | $12.25 \pm 0.05$ | $12.23 \pm 0.10$ | $1.37 e+10_{1.78 e+10}^{1.37 e+10}$ | $9.39 e+09_{1.212 e+10}^{9.309+}$ | $4.29 e+09_{5.69 e+09}^{4+2.2 e+09}$ | $1.36 e+10_{1.72 e+11}^{1.36 e+10}$ |
| 0739 | 9 | $13.44 \pm 0.05$ | $13.47 \pm 0.10$ | $1.30 e+10_{1.30 e+10}^{2.98 e+09}$ | $6.92 e+08_{9.23 e+08}^{6.92 e+08}$ | $1.23 e+10_{1.23 e+10}^{2.06 e+09}$ | $9.73 e+08_{6.33 e+09}^{9.73 e+08}$ |
| 0740 | 11 | $15.40 \pm 0.05$ | $15.42 \pm 0.10$ | $2.30 e+09_{2.30 e+09}^{4.97 e+08}$ | $1.23 e+08_{1.54 e+08}^{1.23 e+08}$ | $2.18 e+09_{2.18 e+09}^{3.43 e+08}$ | $1.72 e+08_{1.33 e+09}^{1.72 e+08}$ |
| 0768 | 7 | $14.47 \pm 0.05$ | $14.50 \pm 0.10$ | $6.62 e+09_{6.62 e+09}^{1.05 e+09}$ | $3.53 e+08_{4.62 e+08}^{3.53 e+08}$ | $6.27 e+09_{6.27 e+09}^{5.86+08}$ | $4.96 e+08_{2.80 e+09}^{4.96 e+08}$ |
| 0780 | 25 | $16.65 \pm 0.10$ | $16.55 \pm 0.10$ | $2.63 e+08_{1.26 e+09}^{1.98 e+08}$ | $8.13 e+07_{8.72 e+07}^{6.402}$ | $1.81 e+08_{1.19 e+09}^{1.11 e+08}$ | $1.15 e+08_{1.16 e+09}^{1.15 e+08}$ |
| 0792 | 17 | $11.20 \pm 0.05$ | $11.17 \pm 0.10$ | $3.70 e+10_{4.39 e+10}^{3.72 e+10}$ | $2.54 e+10_{3.00 e+10}^{2.54 e+10}$ | $1.16 e+10_{1.40 e+10}^{1.16 e+10}$ | $3.70 e+10_{4.49 e+11}^{3.70 e+10}$ |
| 0793 | 3 | $0.00 \pm-99999.00$ | $198.00 \pm 99.00$ | $0.00 e+00_{0.00 e++00}^{0.000}$ | $0.00 e+00_{0.00 e++00}^{0.000+0}$ | $0.00 e+00_{0.00 e++00}^{0.000+0}$ | $6.95 e-310_{6.95 e-310}^{6.955}$ |
| 0801 | 15 | $11.86 \pm 0.05$ | $11.89 \pm 0.10$ | $7.44 e+10_{7.44 e+10}^{1.18 e+10}$ | $3.97 e+09_{5.19 e+09}^{3.97 e+09}$ | $7.04 e+10_{7.04 e+10}^{6.59 e+09}$ | $5.58 e+09_{4.515 e+10}^{5.5 e+09}$ |
| 0802 | 1 | $16.12 \pm 0.05$ | $16.16 \pm 0.10$ | $1.25 e+09_{1.25 e+09}^{1.72 e+08}$ | $6.67 e+07_{9.63 e+07}^{6.34+07}$ | $1.18 e+09_{1.18 e+09}^{8.73 e+07}$ | $9.38 e+07_{1.838++09}^{9.38 e+07}$ |
| 0809 | 7 | $13.80 \pm 0.05$ | $13.84 \pm 0.10$ | $2.08 e+09_{1.68 e+10}^{2.088}$ | $1.25 e+09_{1.25 e+09}^{8.98+08}$ | $8.24 e+08_{1.59 e+10}^{8.24+08}$ | $1.79 e+09_{1.26 e+10}^{1.79 e+09}$ |
| 0825 | 3 | $15.36 \pm 0.05$ | $15.38 \pm 0.10$ | $3.57 e+09_{3.57 e+09}^{4.67+08}$ | $1.91 e+08_{2.82 e+08}^{1.92 e+08}$ | $3.38 e+09_{3.38 e+09}^{1.85 e+08}$ | $2.68 e+08_{2.89 e+09}^{2.26 e+08}$ |

continued

| VCC | $\begin{gathered} \hline \hline \text { Type- } \\ \text { ID } \end{gathered}$ | $\begin{gathered} \mathrm{mag}_{\mathrm{r}, \mathrm{obs}} \\ {[\mathrm{mag}]} \end{gathered}$ | $\operatorname{mag}_{\mathrm{r}, \mathrm{mod}}$ [mag] | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \hline \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \mathrm{M}_{\mathrm{nonvis}} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 0826 | 3 | $14.69 \pm 0.06$ | $14.67 \pm 0.10$ | $1.86 e+09_{1.86 e+09}^{8.57 e+08}$ | $2.72 e+08_{3.11 e+08}^{2.40 e+08}$ | $1.59 e+09_{1.59 e+09}^{6.17 e+08}$ | $3.83 e+08^{3.253 e+09}{ }^{3.83}$ |
| 0827 | 7 | $12.64 \pm 0.05$ | $12.64 \pm 0.10$ | $7.32 e+09_{7.32 e+09}^{7.32+09}$ | $4.42 e+09_{4.42 e+09}^{4.42 e+09}$ | $2.90 e+09^{2} 2.90 e+0909$ | $6.32 e+09_{1.89 e+10}^{6.32 e+09}$ |
| 0836 | 17 | $10.76 \pm 0.05$ | $10.75 \pm 0.10$ | $4.26 e+10_{3.81 e+11}^{4.26 e+10}$ | $2.82 e+10_{2.84 e+10}^{2.03 e+10}$ | $1.43 e+10_{3.61 e+11}^{1.43 e+10}$ | $4.06 e+10_{2.91 e+11}^{4.06 e+10}$ |
| 0841 | 1 | $14.63 \pm 0.05$ | $14.65 \pm 0.10$ | $8.01 e+09_{8.01 e+09}^{9.89 e+08}$ | $4.27 e+08_{5.97 e+08}^{4.27 e+08}$ | $7.58 e+09_{7.58 e+09}^{3.92 e+08}$ | $6.00 e+08_{7.84 e+09}^{6.00 e+08}$ |
| 0849 | 5 | $12.56 \pm 0.05$ | $12.57 \pm 0.10$ | $3.68 e+10_{3.69 e++10}^{5.95 e+}$ | $1.97 e+09^{1.92 e++09}$ | $3.49 e+10_{3.49 e+10}^{3.332+99}$ | $2.76 e+09_{1.80 e+10}^{2.76 e+9}$ |
| 0851 | 7 | $13.13 \pm 0.05$ | $13.13 \pm 0.10$ | $4.20 e+09_{3.62 e+109}^{4.20+09}$ | $2.78 e+09^{1.79 e++09}$ | $1.41 e+09^{1.42 e++10}$ | $4.01 e+09_{2.52 e+10}^{4.01 e+09}$ |
| 0857 | 5 | $10.48 \pm 0.05$ | $10.45 \pm 0.10$ | $6.06 e+10_{5.52 e+111}^{6.06 e+10}$ | $4.16 e+10_{4.88 e+10}^{2.94 e+10}$ | $1.90 e+10_{5.22 e+11}^{1.90 e+10}$ | $6.03 e+10_{8.39 e+11}^{6.03 e+10}$ |
| 0859 | 7 | $13.23 \pm 0.05$ | $13.18 \pm 0.10$ | $6.69 e+09_{8.73 e+09}^{6.688+09}$ | $4.59 e+09_{5.95 e+09}^{4.59+09}$ | $2.10 e+09^{2} 2.78 e++09$ | $6.65 e+09_{1.06 e+11}^{665 e+09}$ |
| 0865 | 7 | $12.17 \pm 0.05$ | $12.16 \pm 0.10$ | $9.07 e+09_{1.98 e+10}^{8.92 e+09}$ | $2.54 e+09_{3.93 e+09}^{2.54 e+09}$ | $6.53 e+09_{1.69 e+10}^{4.99 e+09}$ | $3.49 e+09^{3.49 e+09}$ |
| 0869 | 25 | $14.71 \pm 0.06$ | $14.68 \pm 0.10$ | $1.01 e+09_{8.79 e+09}^{1.01 e+09}$ | $6.73 e+08_{6.74 e+08}^{4.69 e+08}$ | $3.41 e+08_{8.32 e+09}^{3.41 e+08}$ | $9.68 e+08_{9.25 e+09}^{99.68+08}$ |
| 0873 | 7 | $11.31 \pm 0.05$ | $11.29 \pm 0.10$ | $3.17 e+10_{4.15 e+10}^{3.17 e+10}$ | $2.18 e+10_{2.83 e+10}^{2.18 e+10}$ | $9.95 e+09^{99.94 e+09}$ | $3.16 e+10_{4.45 e+11}^{3.16 e+10}$ |
| 0874 | 7 | $11.94 \pm 0.05$ | $11.95 \pm 0.10$ | $1.36 e+10_{1.22 e+11}^{1.36+10}$ | $8.99 e+09_{9.06 e+09}^{6.48 e+09}$ | $4.56 e+09_{1.15 e+11}^{4.56+09}$ | $1.29 e+10_{7}^{1.29 e++10}$ |
| 0888 | 3 | $14.95 \pm 0.05$ | $14.97 \pm 0.10$ | $6.60 e+08_{4.09 e+09}^{6.60 e+08}$ | $2.91 e+08^{2.191 e+08}{ }^{2.18 e+08}$ | $3.69 e+08_{3.87 e+09}^{3.69+08}$ | $4.13 e+08_{2.60 e+09}^{4.133+08}$ |
| 0890 | 1 | $16.02 \pm 0.05$ | $16.03 \pm 0.10$ | $1.35 e+09_{1.35 e+09}^{2.82 e+08}$ | $7.18 e+07_{8.75 e+07}^{7.18 e+07}$ | $1.27 e+09_{1.27 e+09}^{1.95+08}$ | $1.01 e+08_{5.50 e+08}^{1.01 e+08}$ |
| 0905 | 7 | $12.74 \pm 0.05$ | $12.75 \pm 0.10$ | $3.36 e+10_{3.36 e+10}^{5.36+09}$ | $1.79 e+09_{2.34 e+09}^{1.79 e+09}$ | $3.18 e+10_{3.18 e+10}^{2.97 e+09}$ | $2.52 e+09_{1.96 e+10}^{2.52 e+09}$ |
| 0912 | 5 | $11.85 \pm 0.05$ | $11.85 \pm 0.10$ | $1.37 e+10_{1.21 e+11}^{1.37 e+10}$ | $9.10 e+09_{9.10 e+09}^{66.43 e+09}$ | $4.61 e+09_{1.14 e+11}^{4.61 e+09}$ | $1.31 e+10_{8.31 e+10}^{1.31 e+10}$ |
| 0938 | 7 | $12.22 \pm 0.05$ | $12.21 \pm 0.10$ | $9.82 e+09_{2.44 e+10}^{9.82 e+09}$ | $4.33 e+09_{4.33 e+09}^{3.13 e+09}$ | $5.49 e+09^{5.08 e++10}$ | $6.14 e+09_{4.84 e+10}^{6.14 e+09}$ |
| 0939 | 7 | $12.19 \pm 0.05$ | $12.22 \pm 0.10$ | $5.97 e+10_{5.97 e+10}^{99.44 e+9}$ | $3.19 e+09_{4.16 e+09}^{3.19 e+09}$ | $5.66 e+10_{5.66 e+10}^{5.28+09}$ | $4.48 e+09_{4.04 e+10}^{4.48+09}$ |
| 0945 | 11 | $14.86 \pm 0.05$ | $14.82 \pm 0.10$ | $3.34 e+09_{3.34 e+09}^{5.75+09}$ | $1.78 e+08_{1.78 e+08}^{\substack{4.6 e+08}}$ | $3.16 e+09_{3.16 e+09}^{4.14 e+080}$ | $2.50 e+08_{1.19 e+09}^{2.50++0}$ |
| 0950 | 11 | $14.73 \pm 0.05$ | $14.68 \pm 0.10$ | $7.70 e+08_{1.68 e+09}^{7.69+08}$ | $2.15 e+08_{3.12 e+08}^{2.15 e+08}$ | $5.55 e+08_{1.44 e+09}^{5.54 e+08}$ | $2.96 e+08_{2.46 e+09}^{2.96+08}$ |

continued

| VCC | Type- <br> ID | $\begin{gathered} \mathrm{mag}_{\mathrm{r}, \text { obs }} \\ {[\mathrm{mag}]} \end{gathered}$ | $\operatorname{mag}_{\mathrm{r}, \mathrm{mod}}$ [mag] | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \hline \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \hline \mathrm{M}_{\text {nonvis }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 0952 | 3 | $16.40 \pm 0.06$ | $16.48 \pm 0.10$ | $1.17 e+08_{7.39 e+08}^{1.17 e+08}$ | $3.28 e+07_{4.61 e+07}^{3.28 e+07}$ | $8.46 e+07_{7.00 e+08}^{8.44 e+07}$ | $4.51 e+07_{3.11 e+08}^{4.51 e+07}$ |
| 0957 | 7 | $12.00 \pm 0.05$ | $12.00 \pm 0.10$ | $6.69 e+10_{6.69 e+10}^{1.06 e+10}$ | $3.57 e+09_{4.67 e+09}^{3.57 e+09}$ | $6.33 e+10_{6.34 e+10}^{5.92 e+09}$ | $5.02 e+09_{3.44 e+10}^{5.02 e+9}$ |
| 0958 | 17 | $10.86 \pm 0.05$ | $10.82 \pm 0.10$ | $6.26 e+10_{7.63 e+10}^{6.25 e+10}$ | $4.30 e+10_{5.19 e+10}^{4.29 e+10}$ | $1.96 e+10_{2.43 e+10}^{1.96 e+10}$ | $6.26 e+10_{8.31 e+11}^{66.26+10}$ |
| 0963 | 3 | $16.41 \pm 0.06$ | $16.42 \pm 0.10$ | $2.24 e+08_{4.91 e+08}^{1.95 e+08}$ | $6.28 e+07_{8.59 e+07}^{6.28 e+07}$ | $1.62 e+08_{4.19 e+08}^{1.09++08}$ | $8.62 e+07_{8.09 e+08}^{8.62 e+07}$ |
| 0975 | 7 | $0.00 \pm-99999.00$ | $198.00 \pm 99.00$ | $0.00 e+00_{0.00 e+00}^{0.00 e+00}$ | $0.00 e+00_{0.00 e++00}^{0.00 e+00}$ | $0.00 e+00_{0.00 e+00}^{0.00 e+00}$ | $6.95 e-310_{6.95 e-310}^{6.95 e-310}$ |
| 0979 | 17 | $11.27 \pm 0.05$ | $11.25 \pm 0.10$ | $2.40 e+11_{2.40 e+11}^{2.67 e+10}$ | $1.28 e+10_{1.79 e+10}^{1.28 e+10}$ | $2.27 e+11_{2.27 e+11}^{9.00 e+09}$ | $1.80 e+10_{2.18}^{1.80 e++11}$ |
| 0980 | 7 | $13.69 \pm 0.05$ | $13.66 \pm 0.10$ | $9.68 e+09_{9.68 e+09}^{1.72 e+09}$ | $5.16 e+08_{5.17 e+08}^{4.81 e+08}$ | $9.17 e+09_{9.17 e+09}^{1.24 e+09}$ | $7.26 e+08_{3.50 e+09}^{7.26 e+08}$ |
| 0984 | 17 | $11.72 \pm 0.05$ | $11.71 \pm 0.10$ | $2.45 e+10_{2.99 e+10}^{2.45 e+10}$ | $1.68 e+10_{2.04 e+10}^{1.68 e+10}$ | $7.68 e+09_{9.53 e+09}^{7.68 e+09}$ | $2.45 e+10_{3.26 e+11}^{2.45 e+10}$ |
| 0989 | 7 | $16.04 \pm 0.06$ | $16.04 \pm 0.10$ | $1.33 e+09_{1.33 e+09}^{2.82 e+08}$ | $7.11 e+07_{8.73 e+07}^{7.11 e+07}$ | $1.26 e+09_{1.26 e+09}^{1.94+08}$ | $9.99 e+079.959++07$ |
| 0994 | 25 | $17.39 \pm 0.05$ | $17.39 \pm 0.10$ | $1.58 e+08_{2.088+08}^{1.57 e+08}$ | $1.08 e+08_{1.42 e+08}^{1.088+08}$ | $4.94 e+07_{6.65 e+07}^{4.94+07}$ | $1.57 e+08_{2.02 e+09}^{1.57 e+08}$ |
| 0995 | 7 | $14.64 \pm 0.05$ | $14.65 \pm 0.10$ | $4.30 e+09_{4.30 e+09}^{1.70 e+09}$ | $2.30 e+08_{2.48 e+08}^{2.29 e+08}$ | $4.07 e+09_{4.07 e+09}^{1.45+09}$ | $3.23 e+08_{1.67 e+09}^{3.23 e+08}$ |
| 1001 | 3 | $16.39 \pm 0.06$ | $16.39 \pm 0.10$ | $1.79 e+08_{4.46 e+08}^{1.79++8}$ | $7.88 e+07_{7.88 e+07}^{5.70 e+07}$ | $1.00 e+08_{3.81 e+08}^{1.00 e+08}$ | $1.12 e+08_{7.90 e+08}^{1.12 e+08}$ |
| 1011 | 9 | $14.32 \pm 0.05$ | $14.32 \pm 0.10$ | $9.26 e+09_{9.26 e+09}^{1.46+09}$ | $4.94 e+08_{6.44 e+08}^{4.94 e+08}$ | $8.76 e+09_{8.77 e+09}^{8.18+08}$ | $6.94 e+08_{3.91++09}^{6.94++08}$ |
| 1013 | 3 | $16.00 \pm 0.06$ | $15.96 \pm 0.10$ | $1.75 e+09_{1.76 e+09}^{2.78+08}$ | $9.36 e+07_{1.22 e+08}^{8.96 e+07}$ | $1.66 e+09_{1.66 e+09}^{1.55 e+08}$ | $1.32 e+08_{1.48 e+09}^{1.32 e+08}$ |
| 1017 | 3 | $14.49 \pm 0.05$ | $14.49 \pm 0.10$ | $1.20 e+09_{9.10 e+09}^{1.12 e+09}$ | $7.96 e+08_{1.29 e+09}^{4.62 e+08}$ | $4.04 e+08_{8.61 e+09}^{4.04+08}$ | $1.15 e+09_{2.11 e+10}^{1.15 e+09}$ |
| 1020 | 25 | $16.28 \pm 0.06$ | $16.27 \pm 0.10$ | $2.40 e+08_{2.12 e+09}^{2.40 e+08}$ | $1.59 e+08_{1.59 e+08}^{1.12 e+08}$ | $8.08 e+07_{2.00 e+09}^{8.08+07}$ | $2.29 e+08_{2.06 e+09}^{2.29 e+08}$ |
| 1021 | 3 | $14.64 \pm 0.05$ | $14.62 \pm 0.10$ | $1.13 e+09_{7.12 e+09}^{1.13 e+09}$ | $4.96 e+08_{4.96 e+08}^{3.80 e+08}$ | $6.29 e+08_{6.74 e+09}^{6.29+08}$ | $7.04 e+08_{5.64 e+09}^{7.04+08}$ |
| 1047 | 17 | $11.54 \pm 0.05$ | $11.54 \pm 0.10$ | $3.15 e+10_{3.85 e+10}^{3.15 e+10}$ | $2.17 e+10_{2.62 e+10}^{2.16 e+10}$ | $9.88 e+09_{1.23 e+10}^{9.88 e+09}$ | $3.15 e+10_{4.57 e+11}^{3.15 e+10}$ |
| 1060 | 11 | $15.34 \pm 0.05$ | $15.38 \pm 0.10$ | $3.37 e+08_{2.04 e+09}^{3.37 e+08}$ | $9.43 e+07_{1.09 e+08}^{9933+07}$ | $2.42 e+08_{1.93 e+09}^{2.42 e+08}$ | $1.29 e+08_{6.94 e+08}^{1.29 e+08}$ |
| 1086 | 23 | $12.54 \pm 0.05$ | $12.52 \pm 0.10$ | $1.00 e+10_{1.24 e+10}^{1.00 e+10}$ | $6.88 e+09_{8.44 e+09}^{6.88 e+09}$ | $3.14 e+09^{3.919 e+09}$ | $9.97 e+09_{1.25 e+11}^{9.97++9}$ |

continued

| VCC | $\begin{gathered} \hline \hline \text { Type- } \\ \text { ID } \end{gathered}$ | $\begin{gathered} \hline \mathrm{mag}_{\mathrm{r}, \mathrm{obs}} \\ {[\mathrm{mag}]} \end{gathered}$ | $\operatorname{mag}_{\mathrm{r}, \bmod }$ <br> [mag] | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \hline \mathrm{M}_{\text {nonvis }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 1091 | 5 | $13.68 \pm 0.05$ | $13.67 \pm 0.10$ | $9.60 e+09_{9.60 e+09}^{4.05 e+09}$ | $5.12 e+08_{5.92 e+08}^{5.12 e+08}$ | $9.09 e+09_{9.09 e+09}^{3.46 e+09}$ | $7.20 e+08_{3.83 e+09}^{7.20 e+08}$ |
| 1098 | 3 | $16.67 \pm 0.11$ | $16.55 \pm 0.10$ | $1.11 e+09_{1.11 e+09}^{1.59++08}$ | $5.93 e+07_{9.61 e+07}^{5.81 e+07}$ | $1.05 e+09_{1.05 e+09}^{6.31 e+07}$ | $8.34 e+07_{1.30 e+09}^{8.34 e+07}$ |
| 1102 | 3 | $16.93 \pm 0.06$ | $17.00 \pm 0.10$ | $6.19 e+08_{6.20 e+08}^{1.02 e+08}$ | $3.31 e+07_{4.47 e+07}^{3.30 e+07}$ | $5.86 e+08_{5.87 e+08}^{5.68+07}$ | $4.65 e+07_{4.81 e+08}^{4.65 e+07}$ |
| 1106 | 3 | $16.64 \pm 0.07$ | $16.67 \pm 0.10$ | $1.57 e+08_{1.27 e+09}^{1.57+08}$ | $9.46 e+07_{9.46 e+07}^{6.79+07}$ | $6.21 e+07_{1.21 e+09}^{6.21 e+07}$ | $1.35 e+08_{1.08 e+09}^{1.35+08}$ |
| 1110 | 17 | $9.70 \pm 0.05$ | $9.69 \pm 0.10$ | $1.59 e+11_{1.78 e+11}^{1.44 e+11}$ | $1.09 e+11_{1.22 e+11}^{99.87 e}$ | $5.02 e+10_{5.69 e+10}^{4.50 e+10}$ | $1.59 e+11_{1.78 e+12}^{1.59 e+11}$ |
| 1114 | 3 | $13.98 \pm 0.05$ | $13.99 \pm 0.10$ | $1.25 e+10_{1.25 e+10}^{1.25 e+10}$ | $6.69 e+08_{6.69 e+08}^{6.6 .62 e+08}$ | $1.19 e+10_{1.19 e+10}^{1.19 e+10}$ | $9.40 e+08_{2.82 e+09}^{9.40+0}$ |
| 1118 | 7 | $12.41 \pm 0.05$ | $12.43 \pm 0.10$ | $6.77 e+10_{6.77 e+10}^{8.22 e+09}$ | $3.61 e+09_{5.45 e+09}^{3.61 e+09}$ | $6.41 e+10_{6.41 e+10}^{2.77 e+09}$ | $5.08 e+09_{5.64 e+10}^{5.082+09}$ |
| 1121 | 3 | $16.22 \pm 0.08$ | $16.18 \pm 0.10$ | $3.10 e+08_{2.61 e+09}^{2.90+08}$ | $2.13 e+08_{3.07 e+08}^{1.39+08}$ | $9.70 e+07_{2.47 e+09}^{9.70 e+07}$ | $3.08 e+08_{7.12 e+09}^{3.08 e+08}$ |
| 1126 | 7 | $11.91 \pm 0.05$ | $11.88 \pm 0.10$ | $1.62 e+10_{1.89 e+10}^{1.62 e+10}$ | $1.11 e+10_{1.30 e+10}^{1.11 e+10}$ | $5.07 e+09_{5.94 e+09}^{5.07 e+09}$ | $1.61 e+10_{9.91 e+10}^{1.610}$ |
| 1128 | 3 | $16.07 \pm 0.09$ | $16.06 \pm 0.10$ | $7.07 e+08_{1.76 e+09}^{3.25 e+08}$ | $1.03 e+08_{1.15 e+08}^{99.11 e+07}$ | $6.04 e+08_{1.67 e+09}^{2.34 e+08}$ | $1.45 e+08_{1.28 e+09}^{1.45 e}$ |
| 1141 | 1 | $15.59 \pm 0.05$ | $15.59 \pm 0.10$ | $3.11 e+09_{3.11 e+09}^{3.86 e+08}$ | $1.66 e+08_{2.33 e+08}^{1.66 e+08}$ | $2.94 e+09^{1.934 e+09}$ | $2.33 e+08_{3.85 e+09}^{2.33+08}$ |
| 1145 | 5 | $10.47 \pm 0.05$ | $10.45 \pm 0.10$ | $7.46 e+10_{8.74 e+10}^{7.45 e+10}$ | $5.12 e+10_{5.96 e+10}^{5.12 e+10}$ | $2.34 e+10_{2.78 e+10}^{2.34 e+10}$ | $7.45 e+10_{8.99 e+11}^{7.45 e+10}$ |
| 1156 | 7 | $13.97 \pm 0.05$ | $13.99 \pm 0.10$ | $7.92 e+09_{7.92 e+09}^{3.14 e+09}$ | $4.23 e+08_{4.59 e+08}^{4.23 e+08}$ | $7.50 e+09_{7}^{22.588 e+09}$ | $5.94 e+08_{3.07 e+09}^{5.94 e+08}$ |
| 1158 | 17 | $10.89 \pm 0.05$ | $10.90 \pm 0.10$ | $7.10 e+10_{7.90 e+10}^{6.03 e+10}$ | $4.86 e+10_{5.38 e+10}^{4.14 e+10}$ | $2.24 e+10_{2.52 e+10}^{1.89 e+10}$ | $7.10 e+10_{1.01 e+12}^{7.10 e+10}$ |
| 1165 | 25 | $17.40 \pm 0.11$ | $17.31 \pm 0.10$ | $1.19 e+08_{6.53 e+08}^{1.03 e+08}$ | $3.35 e+07_{4.53 e+07}^{3.35 e+07}$ | $8.60 e+07_{6.18 e+08}^{5.75+07}$ | $4.59 e+07_{5.99 e+07}^{4.59 e+07}$ |
| 1166 | 3 | $16.32 \pm 0.08$ | $16.32 \pm 0.10$ | $1.56 e+09_{1.56 e+09}^{1.99 e+08}$ | $8.32 e+07_{1.99 e+08}^{7.94+07}$ | $1.48 e+09_{1.48 e+09}^{7.66+07}$ | $1.17 e+08_{2.47 e+09}^{1.17 e+08}$ |
| 1168 | 3 | $16.51 \pm 0.05$ | $16.28 \pm 0.10$ | $1.01 e+09_{1.01 e+09}^{1.91 e+08}$ | $5.41 e+07_{2.31 e+08}^{5.35+07}$ | $9.59 e+08_{9.59 e+08}^{7.86 e+07}$ | $7.60 e+07_{3.02 e+09}^{7.60 e+07}$ |
| 1169 | 3 | $17.48 \pm 0.08$ | $17.50 \pm 0.10$ | $4.95 e+07_{2.87 e+08}^{4.94+07}$ | $1.38 e+07_{1.53 e+07}^{1.38+07}$ | $3.56 e+07_{2.72 e+08}^{3.56 e+07}$ | $1.90 e+07_{1.00 e+08}^{1.90 e+07}$ |
| 1179 | 29 | $14.75 \pm 0.05$ | $14.75 \pm 0.10$ | $5.62 e+09_{5.62 e+09}^{2.22 e+09}$ | $3.00 e+08_{3.24 e+08}^{3.00 e+08}$ | $5.32 e+09_{5.32 e+09}^{1.89 e+09}$ | $4.21 e+08_{2.18 e+09}^{4.21 e+08}$ |
| 1186 | 25 | $17.16 \pm 0.12$ | $17.11 \pm 0.10$ | $1.34 e+08_{2.93 e+08}^{1.15 e+08}$ | $3.77 e+07_{5.08 e+07}^{3.76 e+07}$ | $9.68 e+07_{2.51 e+08}^{6.44 e+07}$ | $5.17 e+07_{5.54 e+08}^{5.17 e+07}$ |

continued

| VCC | $\begin{gathered} \hline \hline \text { Type- } \\ \text { ID } \end{gathered}$ | mag $_{\mathrm{r}, \mathrm{obs}}$ <br> [mag] | $\operatorname{mag}_{\mathrm{r}, \bmod }$ [mag] | $\begin{gathered} \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \mathrm{M}_{\text {nonvis }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 1188 | 15 | $12.80 \pm 0.05$ | $12.81 \pm 0.10$ | $6.12 e+09_{5.48 e+10}^{6.12 e+09}$ | $4.06 e+09_{4.08 e+09}^{2.92 e+09}$ | $2.06 e+09_{5.19 e+10}^{2.06 e+09}$ | $5.84 e+09_{5.71 e+10}^{5.84 e+09}$ |
| 1189 | 7 | $13.21 \pm 0.05$ | $13.20 \pm 0.10$ | $4.98 e+09_{9}^{3.766 e+09}$ | $1.54 e+09_{1.66 e+09}^{1.37 e+09}$ | $3.43 e+09^{2.100 e+09}$ | $2.18 e+09_{1.51 e+10}^{2.18 e+09}$ |
| 1190 | 17 | $10.95 \pm 0.05$ | $10.94 \pm 0.10$ | $5.88 e+10_{7.17 e+10}^{5.88 e+10}$ | $4.04 e+10_{4.89 e+10}^{4.04 e+10}$ | $1.84 e+10_{2.29 e+10}^{1.84 e+10}$ | $5.88 e+10_{7.82 e+11}^{5.88 e+10}$ |
| 1193 | 7 | $13.52 \pm 0.05$ | $13.53 \pm 0.10$ | $1.93 e+10_{1.93 e+10}^{1.93 e+10}$ | $1.03 e+09_{1.33 e+09}^{1.03 e+09}$ | $1.83 e+10_{1.83 e+10}^{1.83 e+10}$ | $1.45 e+09_{4.35 e+09}^{1.45 e+09}$ |
| 1200 | 3 | $14.70 \pm 0.05$ | $14.72 \pm 0.10$ | $3.67 e+09^{1.557++09}$ | $1.96 e+08_{2.27 e+08}^{1.96 e+08}$ | $3.48 e+09^{1.3338++09}$ | $2.75 e+08_{1.15 e+09}^{2.75 e+08}$ |
| 1205 | 7 | $12.25 \pm 0.05$ | $12.25 \pm 0.10$ | $5.44 e+10_{5.44 e+10}^{8.61 e+09}$ | $2.90 e+09_{3.79 e+09}^{2.90 e+09}$ | $5.15 e+10_{5.15 e+10}^{4.81 e+09}$ | $4.08 e+09^{4.768 e++10}$ |
| 1208 | 11 | $15.13 \pm 0.05$ | $15.16 \pm 0.10$ | $2.92 e+09^{6} 9.292 e+098$ | $1.56 e+08_{1.95 e+08}^{1.56+08}$ | $2.76 e+09^{4} .734 e+088$ | $2.19 e+08_{1.41 e+09}^{2.19 e+08}$ |
| 1217 | 11 | $14.02 \pm 0.05$ | $14.04 \pm 0.10$ | $9.21 e+09_{9.21 e+09}^{1.933+09}$ | $4.91 e+08_{5.98 e+08}^{4.91 e+08}$ | $8.72 e+09_{8.72 e+09}^{1.33 e+09}$ | $6.91 e+08_{5.26 e+09}^{6.91 e+08}$ |
| 1227 | 25 | $16.52 \pm 0.11$ | $16.66 \pm 0.10$ | $5.24 e+08_{5.24 e+08}^{3.922+08}$ | $3.57 e+08_{3.57 e+08}^{2.53 e+08}$ | $1.67 e+08_{1.67 e+08}^{1.24+08}$ | $5.24 e+08_{8.23 e+09}^{5.24+08}$ |
| 1237 | 29 | $15.22 \pm 0.05$ | $15.24 \pm 0.10$ | $3.71 e+09^{4.976 e+08}$ | $1.98 e+08_{2.58 e+08}^{1.87 e+08}$ | $3.51 e+09^{2.552 e+08}$ | $2.78 e+08_{4.66 e+09}^{2.78+08}$ |
| 1257 | 3 | $15.96 \pm 0.05$ | $15.94 \pm 0.10$ | $2.53 e+08_{1.39 e+09}^{2.53 e+08}$ | $7.07 e+07_{8.04 e+07}^{7.07 e+07}$ | $1.82 e+08_{1.31 e+09}^{1.82 e+08}$ | $9.71 e+07_{6.22 e+08}^{9.71 e+07}$ |
| 1266 | 9 | $14.91 \pm 0.05$ | $14.92 \pm 0.10$ | $5.23 e+08_{5.23 e+08}^{5.23+08}$ | $1.46 e+08_{1.46 e+08}^{1.46 e+08}$ | $3.76 e+08_{3.77 e+08}^{3.76 e+08}$ | $2.01 e+08_{6.03 e+08}^{2.02 e+08}$ |
| 1273 | 3 | $14.45 \pm 0.05$ | $14.46 \pm 0.10$ | $1.23 e+09_{1.02 e+10}^{1.23 e+09}$ | $8.18 e+08_{8.18 e+08}^{5.43 e+08}$ | $4.15 e+08_{9.64 e+09}^{4.15 e+08}$ | $1.18 e+09_{9.88 e+09}^{1.18 e+09}$ |
| 1287 | 3 | $25.67 \pm 5.10$ | $25.79 \pm 0.10$ | $2.23 e+04_{1.36 e+05}^{2.23+04}$ | $6.23 e+03_{1.79 e+04}^{66.23 e+03}$ | $1.61 e+04_{1.28 e+05}^{1.18 e+04}$ | $8.56 e+03_{1.88 e+05}^{8.56 e+03}$ |
| 1290 | 5 | $12.25 \pm 0.05$ | $12.25 \pm 0.10$ | $6.65 e+10_{6.65 e+10}^{8.38+09}$ | $3.55 e+09_{5.06 e+09}^{3.55 e+09}$ | $6.29 e+10_{6.30 e+10}^{3.32 e+09}$ | $4.99 e+09^{4.969 e++10}$ |
| 1294 | 19 | $14.69 \pm 0.05$ | $14.70 \pm 0.10$ | $9.93 e+08_{8.70 e+09}^{9.933+08}$ | $6.59 e+08_{6.59 e+08}^{4.64 e+08}$ | $3.34 e+08_{8.24 e+09}^{3.34+08}$ | $9.48 e+08_{9.12 e+09}^{99.48+08}$ |
| 1313 | 1 | $16.64 \pm 0.05$ | $16.77 \pm 0.10$ | $5.58 e+07_{9}^{4.34 e+07}$ | $1.87 e+07_{3.24 e+07}^{1.87}$ | $3.71 e+07_{6.73 e+07}^{1.36+07}$ | $2.61 e+07_{5.35 e+08}^{2.61 e+07}$ |
| 1326 | 17 | $12.26 \pm 0.05$ | $12.27 \pm 0.10$ | $1.09 e+10_{9.80 e+10}^{1.09 e+10}$ | $7.26 e+09_{7.91 e+09}^{5.23 e+09}$ | $3.68 e+09_{9.27 e+10}^{3.61 e+09}$ | $1.04 e+10_{1.34 e+11}^{1.04 e+10}$ |
| 1330 | 17 | $11.28 \pm 0.05$ | $11.25 \pm 0.10$ | $4.29 e+10_{4.29 e+10}^{3.51 e+10}$ | $2.92 e+10_{2.92 e+10}^{2.41 e+10}$ | $1.37 e+10_{1.37 e+10}^{1.10 e+10}$ | $4.29 e+10_{4.82 e+11}^{4.29 e+10}$ |
| 1331 | 25 | $16.57 \pm 0.11$ | $16.50 \pm 0.10$ | $2.65 e+08_{1.48 e+09}^{4.231+0+08}$ | $7.42 e+07_{1.02 e+08}^{7.42 e+07}$ | $1.91 e+08_{1.40 e+09}^{1.29+08}$ | $1.02 e+08_{1.33 e+09}^{1.02 e+08}$ |

continued

| VCC | $\begin{gathered} \hline \hline \text { Type- } \\ \text { ID } \end{gathered}$ | $\begin{gathered} \text { mag }_{\mathrm{r}, \mathrm{bb}} \\ {[\mathrm{mag}]} \end{gathered}$ | $\begin{gathered} \mathrm{mag}_{\mathrm{r}, \mathrm{mod}} \\ {[\mathrm{mag}]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \mathrm{M}_{\text {nonvis }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 1336 | 25 | $16.84 \pm 0.10$ | $16.88 \pm 0.10$ | $6.45 e+08_{6.46 e+08}^{1.13 e+08}$ | $3.44 e+07_{1.45 e+08}^{3.22 e+07}$ | $6.11 e+08_{6.11 e+08}^{4.51 e+07}$ | $4.84 e+07_{1.98 e+09}^{4.84 e+07}$ |
| 1356 | 11 | $14.75 \pm 0.05$ | $14.77 \pm 0.10$ | $3.61 e+09_{3.61 e+09}^{1.47 e+09}$ | $1.93 e+08_{2.15 e+08}^{1.93 e+08}$ | $3.42 e+09_{3.42 e+09}^{1.26 e+09}$ | $2.71 e+08_{1.12 e+09}^{2.71 e+08}$ |
| 1358 | 30 | $14.46 \pm 0.05$ | $14.54 \pm 0.10$ | $8.51 e+09_{8.51 e+09}^{1.05 e+09}$ | $4.54 e+08_{6.34 e+08}^{4.54+08}$ | $8.05 e+09_{8.06 e+09}^{4.17 e+08}$ | $6.38 e+08_{8.33 e+09}^{6.38+08}$ |
| 1361 | 25 | $16.77 \pm 0.11$ | $16.76 \pm 0.10$ | $1.77 e+09_{1.77 e+09}^{1.93 e+08}$ | $9.43 e+07_{1.75 e+08}^{99.43 e+07}$ | $1.67 e+09_{1.67 e+09}^{6.06+07}$ | $1.33 e+08_{3.89 e+09}^{1.33 e+08}$ |
| 1374 | 1 | $14.10 \pm 0.05$ | $14.10 \pm 0.10$ | $6.48 e+09_{6.48 e+09}^{6.48 e+09}$ | $3.46 e+08_{3.46 e+08}^{3.46+08}$ | $6.13 e+09_{6.13 e+09}^{6.13 e+09}$ | $4.86 e+08_{1.46 e+09}^{8.86+08}$ |
| 1375 | 7 | $11.40 \pm 0.05$ | $11.40 \pm 0.10$ | $1.04 e+11_{1.04 e+11}^{1.04 t+11}$ | $5.55 e+09_{5.55 e+09}^{5.54+09}$ | $9.84 e+10_{9.84 e+10}^{9.84 e+10}$ | $7.79 e+09_{2.34 e+10}^{7.79 e}$ |
| 1377 | 3 | $14.70 \pm 0.06$ | $14.71 \pm 0.10$ | $7.34 e+08_{4.03 e+09}^{7.33++08}$ | $2.05 e+08_{2.36 e+08}^{2.05 e+08}$ | $5.28 e+08_{3.82 e+09}^{5.27 e+08}$ | $2.82 e+08_{2.11 e+09}^{2.82 e+08}$ |
| 1379 | 7 | $11.93 \pm 0.05$ | $11.94 \pm 0.10$ | $7.72 e+10_{7}^{1.73 e+10} 10$ | $4.12 e+09_{5.39 e+09}^{4.12 e+09}$ | $7.31 e+10_{7.32 e+10}^{6.84+09}$ | $5.79 e+0{ }_{3}^{5.759++09}$ |
| 1393 | 7 | $12.72 \pm 0.05$ | $12.74 \pm 0.10$ | $3.76 e+10_{3.76 e+10}^{5.95 e+09}$ | $2.01 e+09_{2.62 e+09}^{2.01 e+09}$ | $3.56 e+10_{3.56 e+10}^{3.33+09}$ | $2.82 e+0{ }_{1.83 e+10}^{2.82 e+09}$ |
| 1397 | 3 | $17.72 \pm 0.09$ | $17.68 \pm 0.10$ | $3.71 e+08_{3.71 e+07}^{5.65 e+07}$ | $1.98 e+07_{3.41 e+07}^{1.94+07}$ | $3.51 e+08_{3.51++07}^{2.24 e+07}$ | $2.78 e+07_{4.39 e+08}^{7.78+07}$ |
| 1403 | 3 | $16.25 \pm 0.07$ | $16.25 \pm 0.10$ | $3.56 e+08_{7.92 e+08}^{3.14 e+08}$ | $9.96 e+07_{1.38 e+08}^{9.96+07}$ | $2.57 e+08_{6.76 e+08}^{1.75 e+08}$ | $1.37 e+08_{1.49 e+09}^{1.37 e+08}$ |
| 1408 | 25 | $16.99 \pm 0.11$ | $16.95 \pm 0.10$ | $2.12 e+08_{1.18 e+09}^{1.35 e+08}$ | $5.95 e+07_{1.28 e+07}^{5.95 e+07}$ | $1.53 e+08_{1.12 e+09}^{4.53 e+09}$ | $8.16 e+07_{2.07 e+09}^{8.16 e+07}$ |
| 1410 | 11 | $13.77 \pm 0.05$ | $13.78 \pm 0.10$ | $1.23 e+10_{1.23 e+109}^{1.95 e+09}$ | $6.54 e+08_{8.59 e+08}^{6.54+08}$ | $1.16 e+10_{1.16 e+10}^{1.09+09}$ | $9.19 e+08_{7.18 e+09}^{9.19 e+08}$ |
| 1411 | 1 | $14.90 \pm 0.05$ | $14.90 \pm 0.10$ | $4.99 e+09_{4.99 e+09}^{7.87 e+08}$ | $2.66 e+08_{3.47 e+08}^{2.66 e+08}$ | $4.72 e+09_{4.72 e+08}^{4.46 e+08}$ | $3.74 e+08_{4.67 e+09}^{3.74 e+08}$ |
| 1412 | 17 | $10.71 \pm 0.05$ | $10.69 \pm 0.10$ | $7.29 e+10_{8.55 e+10}^{7.29 e+10}$ | $5.01 e+10_{5.83 e+10}^{5.010+10}$ | $2.29 e+10_{2.72 e+10}^{2.28 e+10}$ | $7.29 e+10_{8.80 e+11}^{7.29 e+10}$ |
| 1413 | 25 | $17.20 \pm 0.10$ | $17.16 \pm 0.10$ | $2.30 e+08_{1.29 e+09}^{1.44 e+08}$ | $6.44 e+07_{9.92 e+07}^{6.44 e+07}$ | $1.66 e+08_{1.22 e+09}^{4.53 e+07}$ | $8.84 e+07_{1.85 e+09}^{8.84 e+07}$ |
| 1419 | 23 | $12.54 \pm 0.05$ | $12.56 \pm 0.10$ | $8.44 e+09_{7.55 e+10}^{88.44+9}$ | $5.60 e+09_{6.09 e+09}^{4.03 e+09}$ | $2.84 e+09_{7.15 e+10}^{2.78 e+09}$ | $8.05 e+00_{9.62 e+10}^{8.05 e+09}$ |
| 1426 | 3 | $14.81 \pm 0.05$ | $14.78 \pm 0.10$ | $1.83 e+09_{1.05 e+10}^{1.19 e+09}$ | $5.11 e+08_{7.88 e+08}^{5.11 e+08}$ | $1.32 e+09_{9.90 e+09}^{4.00 e+08}$ | $7.02 e+08_{1.26 e+10}^{7.02 e+08}$ |
| 1427 | 29 | $14.64 \pm 0.05$ | $14.69 \pm 0.10$ | $6.61 e+09_{6.61+09}^{8.84 e+08}$ | $3.53 e+08=0.53+0808$ | $6.26 e+09_{6.26 e+09}^{3}$ | $4.95 e+08_{5.99 e+09}^{4.955+08}$ |
| 1435 | 3 | $13.49 \pm 0.05$ | $13.49 \pm 0.10$ | $2.32 e+10_{2.32 e+109}^{2.87 e+10}$ | $1.24 e+09_{1.73 e+09}^{1.24 e+09}$ | $2.20 e+10_{2.20 e+10}^{1.14 e+09}$ | $1.74 e+0{ }_{2} 9.727 e+109$ |

continued

| VCC [1] | TypeID [2] | $\begin{gathered} \hline \hline \mathrm{mag}_{\mathrm{r}, \mathrm{obs}} \\ {[\mathrm{mag}]} \\ {[3]} \end{gathered}$ | $\begin{gathered} \mathrm{mag}_{\mathrm{r}, \mathrm{mod}} \\ {[\mathrm{mag}]} \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \\ {[5]} \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \\ {[6]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{M}_{\mathrm{gas}} \\ {\left[\mathrm{M}_{\odot}\right]} \\ {[7]} \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{M}_{\text {nonvis }} \\ {\left[\mathrm{M}_{\odot}\right]} \\ {[8]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1437 | 1 | $14.38 \pm 0.05$ | $14.40 \pm 0.10$ | $8.91 e+09_{8.91 e+09}^{1.15 e+09}$ | $4.75 e+08_{6.97 e+08}^{4.75 e+08}$ | $8.43 e+09_{8.43 e+09}^{4.58++08}$ | $6.68 e+08_{1.03 e+10}^{6.68 e+08}$ |
| 1442 | 9 | $13.85 \pm 0.05$ | $13.85 \pm 0.10$ | $1.22 e+10_{1.22 e+10}^{1.93 e+09}$ | $6.52 e+08_{8.52 e+08}^{6.52 e+08}$ | $1.16 e+10_{1.16 e+10}^{1.08+09}$ | $9.16 e+08_{5.94 e+09}^{99.16 e+08}$ |
| 1448 | 25 | $13.13 \pm 0.05$ | $13.14 \pm 0.10$ | $2.58 e+10_{2.58}^{4.09++09}$ | $1.38 e+09_{1.80 e+09}^{1.38+09}$ | $2.44 e+10_{2.44 e+10}^{2.29 e+09}$ | $1.94 e+09_{1.98 e+10}^{1.94 e+09}$ |
| 1450 | 7 | $12.44 \pm 0.05$ | $12.47 \pm 0.10$ | $3.60 e+10_{3.60 e+10}^{7.54 e+09}$ | $1.92 e+09^{1.34 e+09} 1.09$ | $3.41 e+10_{3.41 e+10}^{5.21 e+09}$ | $2.70 e+09_{1.47 e+10}^{2.70 e+09}$ |
| 1455 | 3 | $15.89 \pm 0.05$ | $15.89 \pm 0.10$ | $2.13 e+09_{2.13 e+09}^{3.36 e+08}$ | $1.14 e+08_{1.48 e+08}^{1.14 e+08}$ | $2.02 e+09_{2.02 e+09}^{1.88 e+08}$ | $1.60 e+08_{1.44 e+09}^{1.60 e+08}$ |
| 1459 | 1 | $15.88 \pm 0.05$ | $15.87 \pm 0.10$ | $2.61 e+09^{3} 2.622 e+08$ | $1.39 e+08_{1.94 e+08}^{1.33+08}$ | $2.47 e+09^{1.47 e++09}$ | $1.96 e+08_{4.09 e+09}^{1.96 e+08}$ |
| 1465 | 3 | $15.05 \pm 0.06$ | $14.74 \pm 0.10$ | $5.82 e+08_{3.67 e+09}^{5.81 e+08}$ | $1.63 e+08_{9.73 e+08}^{1.63+08}$ | $4.19 e+08_{3.48 e+09}^{3.31 e+08}$ | $2.23 e+08_{1.23 e+10}^{2.23 e+08}$ |
| 1468 | 3 | $15.46 \pm 0.05$ | $15.50 \pm 0.10$ | $2.98 e+08_{1.83 e+09}^{2.98+08}$ | $8.36 e+07_{9,76 e+07}^{8.36+07}$ | $2.15 e+08_{1.73 e+09}^{2.15 e+08}$ | $1.15 e+08_{6.19 e+08}^{1.15+08}$ |
| 1507 | 11 | $14.81 \pm 0.05$ | $14.83 \pm 0.10$ | $5.93 e+09_{5.93 e+09}^{7.76 e+08}$ | $3.16 e+08_{4.69 e+08}^{3.16 e+08}$ | $5.61 e+09_{5.61 e+09}^{3.08 e+08}$ | $4.45 e+08_{3.85 e+09}^{4.45 e}$ |
| 1508 | 7 | $11.72 \pm 0.05$ | $11.72 \pm 0.10$ | $7.26 e+10_{7.26 e+10}^{7.26 e+10}$ | $3.87 e+09_{3.87 e+09}^{3.87 e+09}$ | $6.88 e+10_{6.88 e+10}^{6.87 e+10}$ | $5.44 e+09^{5.454 e+09}$ |
| 1515 | 3 | $17.18 \pm 0.09$ | $17.20 \pm 0.10$ | $3.73 e+08_{3.73 e+08}^{5.95+07}$ | $1.99 e+07_{1.01 e+08}^{1.67+07}$ | $3.53 e+08_{3.53 e+08}^{3.42 e+07}$ | $2.79 e+07_{1.13 e+09}^{2.79 e+07}$ |
| 1516 | 7 | $12.10 \pm 0.05$ | $12.09 \pm 0.10$ | $8.42 e+10_{8.42 e+10}^{1.04 e+10}$ | $4.49 e+09_{6.27 e+09}^{4.49+09}$ | $7.97 e+10_{7.97 e+10}^{4.12 e+09}$ | $6.31 e+09_{5.88 e+10}^{6.31 e+09}$ |
| 1524 | 9 | $12.64 \pm 0.05$ | $12.60 \pm 0.10$ | $2.83 e+10_{2.83 e+10}^{5.13 e+09}$ | $1.51 e+09_{1.63 e+09}^{1.44 e+09}$ | $2.68 e+10_{2.68 e+10}^{3.69+09}$ | $2.12 e+09_{1.49 e+10}^{2.122+09}$ |
| 1529 | 9 | $13.93 \pm 0.05$ | $13.94 \pm 0.10$ | $1.22 e+10_{1.23 e+10}^{1.94 e+09}$ | $6.54 e+08_{8.54 e+08}^{6.53 e+08}$ | $1.16 e+10_{1.16 e+10}^{1.08+09}$ | $9.19 e+08_{8.29 e+09}^{9.19 e+08}$ |
| 1532 | 7 | $12.80 \pm 0.05$ | $12.83 \pm 0.10$ | $3.47 e+10_{3.48 e+10}^{5.50 e+09}$ | $1.85 e+09^{1.425 e+09}$ | $3.29 e+10_{3.29 e+10}^{3.08 e+09}$ | $2.61 e+09^{2} 2.392 e+10$ |
| 1540 | 5 | $10.17 \pm 0.05$ | $10.15 \pm 0.10$ | $1.17 e+11_{1.04 e+12}^{1.13 e+11}$ | $7.74 e+10_{8.27 e+10}^{5.55 e+10}$ | $3.93 e+10_{9.85 e+11}^{3.53 e+10}$ | $1.11 e+11_{1.06 e+12}^{1.11 e+11}$ |
| 1552 | 17 | $11.25 \pm 0.05$ | $11.24 \pm 0.10$ | $3.43 e+10_{4.25 e+10}^{3.42 e+10}$ | $2.35 e+10_{2.90 e+10}^{2.35 e+10}$ | $1.07 e+10_{1.36 e+10}^{1.07 e+10}$ | $3.43 e+10_{4.62 e+11}^{3.43 e+10}$ |
| 1554 | 11 | $11.57 \pm 0.05$ | $11.60 \pm 0.10$ | $1.18 e+10_{6.70 e+10}^{1.18}$ | $3.30 e+09^{3.350 e+09}$ | $8.49 e+09_{6.35 e+10}^{88.4 e+09}$ | $4.53 e+09_{2.36 e+10}^{4.53}$ |
| 1555 | 7 | $9.69 \pm 0.05$ | $9.72 \pm 0.10$ | $7.39 e+11_{7.39 e+11}^{19.13 e+10}$ | $3.94 e+10_{5.51 e+10}^{3.94 e+10}$ | $6.99 e+11_{7.00 e+11}^{3.62 e+10}$ | $5.54 e+10_{5.89 e+11}^{5.54 e+10}$ |
| 1557 | 7 | $13.39 \pm 0.05$ | $13.41 \pm 0.10$ | $2.45 e+10_{2.45 e+10}^{3.03 e+09}$ | $1.31 e+09_{1.83 e+09}^{1.31 e+09}$ | $2.32 e+10_{2.32 e+10}^{1.20 e+09}$ | $1.84 e+09_{1.27 e+10}^{1.84 e+09}$ |

continued

| VCC | $\begin{gathered} \hline \text { Type- } \\ \text { ID } \end{gathered}$ | $\begin{gathered} \hline \mathrm{mag}_{\mathrm{r}, \text { obs }} \\ {[\mathrm{mag}]} \end{gathered}$ | $\begin{gathered} \mathrm{mag}_{\mathrm{r}, \mathrm{mod}} \\ {[\mathrm{mag}]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \mathrm{M}_{\mathrm{nonvis}} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 1562 | 7 | $10.16 \pm 0.05$ | $10.16 \pm 0.10$ | $5.58 e+11_{5.58 e+11}^{6.90 e+10}$ | $2.98 e+10_{4.16 e+10}^{2.98 e+10}$ | $5.28 e+11_{5.28 e+11}^{2.73 e+10}$ | $4.18 e+10_{1.85 e+11}^{4.18 e+10}$ |
| 1566 | 9 | $14.48 \pm 0.05$ | $14.48 \pm 0.10$ | $1.46 e+09_{3.21 e+09}^{1.29 e+09}$ | $4.09 e+08_{5.68 e+08}^{4.09+08}$ | $1.05 e+09_{2.74 e+09}^{7.21 e+08}$ | $5.62 e+08_{5.45 e+09}^{5.62 e+08}$ |
| 1569 | 7 | $14.88 \pm 0.05$ | $14.92 \pm 0.10$ | $4.60 e+09_{4.60 e+09}^{7.25 e+08}$ | $2.45 e+08_{3.19 e+08}^{2.45 e+08}$ | $4.35 e+09_{4.35 e+09}^{4.05 e+08}$ | $3.45 e+08_{3.53 e+09}^{3.45+08}$ |
| 1572 | 1 | $15.35 \pm 0.05$ | $15.39 \pm 0.10$ | $5.27 e+08_{6.90 e+08}^{3.99 e+08}$ | $1.87 e+08_{2.76 e+08}^{1.87+08}$ | $3.40 e+08_{4.76 e+08}^{1.33+08}$ | $2.59 e+08_{6.15 e+09}^{2.59 e+08}$ |
| 1574 | 23 | $16.45 \pm 0.05$ | $16.46 \pm 0.10$ | $2.51 e+08_{1.59 e+09}^{1.96 e+08}$ | $1.11 e+08_{1.19 e+08}^{8.4 e+07}$ | $1.40 e+08_{1.51 e+09}^{7.79 e+07}$ | $1.57 e+08_{1.34 e+09}^{1.57 e+08}$ |
| 1575 | 11 | $12.76 \pm 0.05$ | $12.76 \pm 0.10$ | $5.95 e+09_{5.95 e+09}^{5.95++09}$ | $3.95 e+09_{3.95 e+09}^{3.95 e+09}$ | $2.00 e+09^{2.000 e++09}$ | $5.68 e+09_{1.70 e+10}^{5.68 e+09}$ |
| 1581 | 11 | $13.93 \pm 0.05$ | $13.91 \pm 0.10$ | $1.24 e+10_{1.24 e+10}^{1.97 e+9}$ | $6.63 e+08_{8.66 e+08}^{6.63 e+08}$ | $1.18 e+10_{1.18 e+10}^{1.10 e+09}$ | $9.32 e+08_{7.27 e+09}^{9.32 e+08}$ |
| 1582 | 25 | $16.14 \pm 0.07$ | $16.14 \pm 0.10$ | $1.68 e+09_{1.68 e+09}^{2.66+08}$ | $8.97 e+07_{1.17 e+08}^{8.97 e+07}$ | $1.59 e+09_{1.59 e+09}^{1.49 e+08}$ | $1.26 e+08_{1.30 e+09}^{1.26 e+08}$ |
| 1585 | 3 | $14.66 \pm 0.05$ | $14.66 \pm 0.10$ | $1.66 e+09_{4.14 e+09}^{1.66 e+09}$ | $2.42 e+08_{2.42 e+08}^{2.21 e+08}$ | $1.42 e+09^{1.422 e+09}$ | $3.41 e+08_{1.64 e+09}^{3.41 e+08}$ |
| 1588 | 7 | $11.54 \pm 0.05$ | $11.54 \pm 0.10$ | $1.83 e+10_{1.57 e+11}^{1.83 e+10}$ | $1.21 e+10_{1.21 e+10}^{8.37 e+09}$ | $6.14 e+09_{1.49 e+11}^{6.14 e+09}$ | $1.74 e+10_{1.22 e+11}^{1.74 e+10}$ |
| 1596 | 3 | $17.14 \pm 0.07$ | $17.16 \pm 0.10$ | $4.18 e+08_{4.18 e+08}^{7.72 e+07}$ | $2.23 e+07_{3.12 e+07}^{2.16 e+07}$ | $3.96 e+08_{3.96 e+08}^{5.56+07}$ | $3.14 e+07_{2.81 e+08}^{3.14 e+07}$ |
| 1605 | 9 | $16.61 \pm 0.06$ | $16.63 \pm 0.10$ | $8.03 e+08_{8.04 e+08}^{1.42 e+08}$ | $4.29 e+07_{6.26 e+07}^{4.29+07}$ | $7.61 e+08_{7.66 e+07}^{7.95+08}$ | $6.03 e+07_{5.50 e+08}^{6.03 e+07}$ |
| 1615 | 5 | $9.72 \pm 0.05$ | $9.69 \pm 0.10$ | $1.43 e+11_{1.77 e+11}^{1.43 e+11}$ | $9.79 e+10_{1.21 e+11}^{99.79+10}$ | $4.47 e+10_{5.64 e+10}^{4.47 e+10}$ | $1.43 e+11_{2.09 e+12}^{1.43 e+11}$ |
| 1623 | 29 | $15.99 \pm 0.05$ | $16.05 \pm 0.10$ | $1.67 e+09_{1.67 e+09}^{1.92 e+08}$ | $8.92 e+07_{1.51 e+08}^{8.46 e+07}$ | $1.58 e+09_{1.58 e+09}^{6.43 e+07}$ | $1.25 e+08_{3.36 e+09}^{1.25 e+08}$ |
| 1624 | 7 | $12.78 \pm 0.05$ | $12.78 \pm 0.10$ | $7.46 e+09_{6.68 e+10}^{7.29 e+9}$ | $4.95 e+09_{5.01++09}^{3.56 e+09}$ | $2.51 e+09_{6.32 e+10}^{2.288+09}$ | $7.12 e+09_{4.09 e+10}^{7.12 e+09}$ |
| 1644 | 11 | $16.80 \pm 0.06$ | $16.69 \pm 0.10$ | $6.57 e+08_{6.57 e+08}^{1.22 e+08}$ | $3.50 e+07_{4.80 e+07}^{3.41 e+07}$ | $6.22 e+08_{6.22 e+08}^{8.78+07}$ | $4.92 e+07_{4.17 e+08}^{4.92 e+07}$ |
| 1645 | 25 | $17.57 \pm 0.14$ | $17.61 \pm 0.10$ | $6.95 e+07_{4.45 e+08}^{6.04++07}$ | $3.06 e+07_{3.65 e+07}^{2.27 e+07}$ | $3.89 e+07_{4.21 e+08}^{2.40+07}$ | $4.35 e+07_{5.18 e+08}^{4.35+07}$ |
| 1654 | 3 | $15.48 \pm 0.05$ | $15.49 \pm 0.10$ | $2.25 e+09_{2.25 e+09}^{4.74 e+08}$ | $1.20 e+08_{1.47 e+08}^{1.20 e+08}$ | $2.13 e+09_{2.13 e+09}^{3.27 e+08}$ | $1.69 e+08_{1.29 e+09}^{1.69+08}$ |
| 1656 | 25 | $15.88 \pm 0.06$ | $15.92 \pm 0.10$ | $2.92 e+08_{2.37 e+09}^{2.292+08}$ | $1.76 e+08_{2.16 e+08}^{1.27 e+08}$ | $1.16 e+08_{2.25 e+09}^{1.09 e+08}$ | $2.52 e+08_{2.71 e+09}^{2.52 e+08}$ |
| 1675 | 15 | $13.72 \pm 0.05$ | $13.72 \pm 0.10$ | $2.18 e+09_{1.76 e+10}^{2.18 e+9}$ | $1.32 e+09_{1.32 e+09}^{994 l e+08}$ | $8.64 e+08_{1.67 e+10}^{8.64 e+08}$ | $1.88 e+09_{1.34 e+10}^{1.88 e+09}$ |

continued

| VCC [1] | TypeID [2] | mag $_{\text {r,obs }}$ <br> [mag] <br> [3] | $\operatorname{mag}_{\mathrm{r}, \bmod }$ <br> [mag] <br> [4] | $\begin{gathered} \hline \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \\ {[5]} \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \\ {[6]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{M}_{\mathrm{gas}} \\ {\left[\mathrm{M}_{\odot}\right]} \\ {[7]} \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{M}_{\text {nonvis }} \\ {\left[\mathrm{M}_{\odot}\right]} \\ {[8]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1678 | 9 | $13.55 \pm 0.05$ | $13.59 \pm 0.10$ | $1.10 e+10_{1.10 e+10}^{4.39 e+09}$ | $5.89 e+08_{6.42 e+08}^{5.89 e+08}$ | $1.04 e+10_{1.04 e+10}^{3.75+0}$ | $8.27 e+00_{4.29 e+09}^{8.27 e+08}$ |
| 1685 | 9 | $15.42 \pm 0.05$ | $15.44 \pm 0.10$ | $1.88 e+09_{1.88 e+09}^{3.39 e+08}$ | $1.00 e+08_{1.16 e+08}^{99.49+07}$ | $1.78 e+09_{1.78 e+09}^{2.44 e+08}$ | $1.41 e+08_{8.81 e+08}^{1.41 e+08}$ |
| 1686 | 11 | $12.61 \pm 0.05$ | $12.60 \pm 0.10$ | $2.83 e+10_{2.83 e+10}^{6.661+09}$ | $1.51 e+09_{2.055+09}^{1.51 e+09}$ | $2.68 e+10_{2.68 e+10}^{4.56 e+10}$ | $2.12 e+09_{1.68 e+10}^{2.12 e+09}$ |
| 1690 | 17 | $9.27 \pm 0.05$ | $9.25 \pm 0.10$ | $1.62 e+11_{1.45 e+12}^{1.62 e+11}$ | $1.07 e+11_{1.08 e+11}^{7.73 e+10}$ | $5.44 e+10_{1.37 e+12}^{5.44+10}$ | $1.54 e+11_{1.64 e+12}^{1.54 e+11}$ |
| 1696 | 7 | $10.81 \pm 0.05$ | $10.78 \pm 0.10$ | $3.73 e+10_{3.35 e+11}^{3.73 e+10}$ | $2.48 e+10_{2.50 e+10}^{1.79 e+10}$ | $1.26 e+10_{3.17 e+11}^{1.26 e+10}$ | $3.56 e+10_{3.16 e+11}^{3.56 e+10}$ |
| 1699 | 11 | $13.73 \pm 0.05$ | $13.74 \pm 0.10$ | $9.84 e+09^{3.888 e+09}$ | $5.25 e+08_{5.65 e+08}^{5.25+08}$ | $9.31 e+09_{9.31 e+09}^{3.31 e+09}$ | $7.37 e+08_{3.81 e+09}^{7.37 e+08}$ |
| 1713 | 29 | $15.43 \pm 0.05$ | $15.40 \pm 0.10$ | $2.39 e+09^{4} 4.14 e++08$ | $1.28 e+08_{1.56 e+08}^{1.28+08}$ | $2.27 e+09_{2.27 e+09}^{2.68 e+08}$ | $1.79 e+08_{2.20 e+09}^{1.79 e+08}$ |
| 1725 | 11 | $13.76 \pm 0.05$ | $13.77 \pm 0.10$ | $1.11 e+10_{1.11 e+10}^{1.11 e+10}$ | $5.93 e+08_{5.93 e+08}^{5.93+08}$ | $1.05 e+10_{1.05 e+10}^{1.05 e+10}$ | $8.33 e+08_{2.50 e+09}^{8.33 e+08}$ |
| 1726 | 9 | $14.63 \pm 0.05$ | $14.64 \pm 0.10$ | $6.73 e+08_{6}^{6.772 e+08}$ | $1.88 e+08_{1.88 e+08}^{1.88 e+08}$ | $4.85 e+08_{4.85 e+08}^{4.8 t+08}$ | $2.58 e+08_{7.75 e+08}^{2.588+08}$ |
| 1727 | 17 | $9.18 \pm 0.05$ | $9.15 \pm 0.10$ | $2.35 e+11_{2.70 e+11}^{2.18 e+11}$ | $1.62 e+11_{1.85 e+11}^{1.50 e+11}$ | $7.37 e+10_{8.57 e+10}^{6.84+10}$ | $2.35 e+11_{2.74 e+12}^{2.35 e+11}$ |
| 1728 | 3 | $15.91 \pm 0.05$ | $15.96 \pm 0.10$ | $1.95 e+09^{2.74 e+08}$ | $1.04 e+08_{1.65 e+08}^{1.04+08}$ | $1.84 e+09_{1.85 e+09}^{1.09 e+08}$ | $1.46 e+08_{1.79 e+09}^{1.46 e+08}$ |
| 1730 | 7 | $11.49 \pm 0.05$ | $11.47 \pm 0.10$ | $2.78 e+10_{3.45 e+10}^{2.78 e+10}$ | $1.91 e+10_{2.35 e+10}^{1.91 e+10}$ | $8.72 e+09_{1.10 e+10}^{88.72 e+9}$ | $2.78 e+10_{3.42 e+11}^{2.78 e+10}$ |
| 1744 | 1 | $16.33 \pm 0.05$ | $16.40 \pm 0.10$ | $1.26 e+08_{8.01 e+08}^{99.79+07}$ | $4.43 e+07_{6.76 e+07}^{3.58+07}$ | $8.14 e+07_{7.58 e+08}^{3.29++7}$ | $6.17 e+07_{1.2 l e+09}^{6.17 e+07}$ |
| 1756 | 25 | $17.41 \pm 0.11$ | $17.33 \pm 0.10$ | $1.19 e+08_{2.61 e+08}^{1.03 e+08}$ | $3.33 e+07_{4.53 e+07}^{3.33 e+07}$ | $8.57 e+07_{2.23 e+07}^{5.75 e+08}$ | $4.58 e+07_{4.93 e+08}^{4.58+07}$ |
| 1757 | 17 | $12.77 \pm 0.05$ | $12.79 \pm 0.10$ | $6.10 e+09_{5}^{6.146 e++10}$ | $4.05 e+09_{4.85 e+09}^{2.91 e+09}$ | $2.05 e+09_{5.17 e+10}^{2.05 e+09}$ | $5.82 e+09_{7.59 e+10}^{5.82 e+09}$ |
| 1758 | 7 | $14.09 \pm 0.05$ | $14.11 \pm 0.10$ | $1.18 e+10_{1.18 e+10}^{1.51 e+09}$ | $6.28 e+08_{9,14 e+08}^{6.28 e+08}$ | $1.11 e+10_{1.11 e+10}^{6.01 e+08}$ | $8.82 e+08_{5.26 e+09}^{8.82 e+08}$ |
| 1760 | 17 | $11.36 \pm 0.05$ | $11.31 \pm 0.10$ | $3.80 e+10_{4.63 e+10}^{3.79 e+10}$ | $2.61 e+10_{3.15 e+10}^{2.61 e+10}$ | $1.19 e+10_{1.48 e+10}^{1.19 e+10}$ | $3.80 e+10_{5.06 e+11}^{3.80 e+10}$ |
| 1776 | 25 | $16.78 \pm 0.17$ | $16.66 \pm 0.10$ | $1.68 e+09_{1.68 e+09}^{1.91 e+08}$ | $8.98 e+07_{1.83 e+08}^{8.46 e+07}$ | $1.59 e+09_{1.59 e+09}^{6.44 e+07}$ | $1.26 e+08_{3.40 e+09}^{1.26 e+08}$ |
| 1778 | 15 | $13.45 \pm 0.05$ | $13.47 \pm 0.10$ | $2.00 e+10_{2.00 e+10}^{2.7 l e+09}$ | $1.07 e+09_{1.63 e+09}^{1.07 e+09}$ | $1.90 e+10_{1.90 e++10}^{1.07}$ | $1.50 e+09_{1.64 e+10}^{1.50 e+0}$ |
| 1784 | 3 | $15.87 \pm 0.07$ | $15.89 \pm 0.10$ | $4.04 e+08_{3.45 e+09}^{3.83+08}$ | $2.77 e+08_{4.00 e+08}^{1.84 e+08}$ | $1.27 e+08_{3.26 e+09}^{1.27 e+08}$ | $4.02 e+08_{8.93 e+09}^{4.02 e+08}$ |

continued

| VCC | $\begin{gathered} \hline \hline \text { Type- } \\ \text { ID } \end{gathered}$ | $\begin{gathered} \mathrm{mag}_{\mathrm{r}, \mathrm{obs}} \\ {[\mathrm{mag}]} \end{gathered}$ | $\begin{gathered} \text { mag }_{\mathrm{r}, \mathrm{mod}} \\ {[\mathrm{mag}]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \mathrm{M}_{\text {nonvis }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 1789 | 3 | $14.51 \pm 0.05$ | $14.53 \pm 0.10$ | $7.88 e+09_{7.89 e+09}^{1.02 e+09}$ | $4.21 e+08_{6.15 e+08}^{4.21 e+08}$ | $7.46 e+09_{7.46 e+09}^{4.04+08}$ | $5.91 e+08_{5.59 e+09}^{5.91 e+08}$ |
| 1791 | 11 | $13.89 \pm 0.05$ | $13.90 \pm 0.10$ | $1.26 e+09_{1.26 e+09}^{1.26 e+09}$ | $3.53 e+08_{3.553 e+08}^{3.53+08}$ | $9.07 e+08_{9.09 e+08}^{9.07 e+08}$ | $4.85 e+08_{1.45 e+09}^{4.85+08}$ |
| 1804 | 29 | $15.05 \pm 0.05$ | $15.07 \pm 0.10$ | $5.72 e+09_{5.72 e+09}^{7.06 e+08}$ | $3.05 e+08_{4.26 e+08}^{3.05 e+08}$ | $5.41 e+09_{5.41 e+09}^{2.80 e+08}$ | $4.29 e+08_{3.43 e+09}^{4.29+08}$ |
| 1811 | 7 | $12.13 \pm 0.05$ | $12.15 \pm 0.10$ | $5.77 e+10_{5.77 e+10}^{99.12 e+09}$ | $3.08 e+09_{4.02 e+09}^{3.088+09}$ | $5.46 e+10_{5.46 e+10}^{5.10 e+9}$ | $4.32 e+09^{4} .832 e++10$ |
| 1813 | 17 | $10.03 \pm 0.05$ | $10.02 \pm 0.10$ | $1.29 e+11_{1.51 e+11}^{1.29 e+11}$ | $8.85 e+10_{1.03 e+11}^{8.84 e+10}$ | $4.04 e+10_{4.80 e+10}^{4.04 e+10}$ | $1.29 e+11_{1.55 e+12}^{1.29 e+11}$ |
| 1816 | 3 | $15.49 \pm 0.06$ | $15.42 \pm 0.10$ | $1.99 e+09_{1.99 e+09}^{33.73 e+08}$ | $1.06 e+08_{1.92 e+08}^{1.05 e+08}$ | $1.89 e+09_{1.89 e+09}^{2.44 e+08}$ | $1.50 e+08_{1.78 e+09}^{1.50 e+08}$ |
| 1821 | 11 | $16.88 \pm 0.06$ | $16.85 \pm 0.10$ | $1.99 e+08_{9.58 e+08}^{1.50+08}$ | $6.17 e+07_{6.63 e+07}^{5.11 e+07}$ | $1.37 e+08_{9.07 e+08}^{8.41 e+07}$ | $8.72 e+07_{7}^{8.72 e++07}$ |
| 1822 | 3 | $15.86 \pm 0.05$ | $15.84 \pm 0.10$ | $1.76 e+09_{1.76 e+09}^{2.92 e+08}$ | $9.36 e+07_{1.29 e+08}^{99.36+07}$ | $1.66 e+09_{1.66 e+09}^{1.63++08}$ | $1.32 e+08_{1.18 e+09}^{1.32+08}$ |
| 1825 | 25 | $15.29 \pm 0.06$ | $15.31 \pm 0.10$ | $5.65 e+08_{4.99 e+09}^{5.64+08}$ | $3.75 e+08_{4.72 e+08}^{2.62 e+08}$ | $1.90 e+08_{4.65 e+09}^{1.90 e+08}$ | $5.39 e+08_{7.43 e+09}^{5.39+08}$ |
| 1834 | 19 | $12.20 \pm 0.05$ | $12.17 \pm 0.10$ | $1.24 e+10_{1.67 e+10}^{1.24+10}$ | $8.51 e+09_{1.14 e+10}^{8.51 e+09}$ | $3.88 e+09_{5.29 e+09}^{3.88 e+09}$ | $1.23 e+10_{1.88 e+11}^{1.23 e+10}$ |
| 1837 | 25 | $15.97 \pm 0.06$ | $16.01 \pm 0.10$ | $2.79 e+08_{2.27 e+09}^{2.79 e+08}$ | $1.68 e+08_{1.688+08}^{1.21 e+08}$ | $1.11 e+08_{2.15 e+09}^{1.11 e+08}$ | $2.41 e+08_{1.92 e+09}^{2.41 e+08}$ |
| 1855 | 19 | $14.72 \pm 0.05$ | $14.73 \pm 0.10$ | $9.27 e+08_{7.50 e+09}^{99.27 e+08}$ | $5.59 e+08_{5.59 e+08}^{4.00 e+08}$ | $3.67 e+08_{7.10 e+09}^{3.67 e+08}$ | $8.00 e+08_{4.26 e+09}^{8.00 e+08}$ |
| 1859 | 17 | $11.66 \pm 0.05$ | $11.64 \pm 0.10$ | $2.02 e+10_{1.77 e+11}^{1.94+10}$ | $1.39 e+10_{1.79 e+10}^{99.28 e+09}$ | $6.33 e+09_{1.65 e+11}^{6.33 e+09}$ | $2.01 e+10_{3.23 e+11}^{2.01 e+10}$ |
| 1860 | 25 | $16.22 \pm 0.06$ | $16.34 \pm 0.10$ | $1.11 e+09_{1.11 e+09}^{1.87+08}$ | $5.92 e+07_{2.42 e+08}^{5.92 e+07}$ | $1.05 e+09_{1.05 c+09}^{7.49 e+07}$ | $8.32 e+07_{3.34 e+09}^{8.32 e+07}$ |
| 1868 | 7 | $12.57 \pm 0.05$ | $12.52 \pm 0.10$ | $1.32 e+10_{1.32 e+10}^{1.01 e+10}$ | $9.00 e+09_{9.00 e+09}^{6.96 e+09}$ | $4.21 e+09_{4.21+e+09}^{3.18 e+09}$ | $1.32 e+10_{1.48 e+11}^{1.32 e+10}$ |
| 1873 | 11 | $16.27 \pm 0.06$ | $16.26 \pm 0.10$ | $4.30 e+08_{1.09 e+09}^{1.97 e+08}$ | $6.28 e+07_{6.29 e+07}^{5.52 e+07}$ | $3.67 e+08_{1.03 e+09}^{1.42 e+08}$ | $8.85 e+07_{5.80 e+08}^{8.85 e+07}$ |
| 1883 | 19 | $11.14 \pm 0.05$ | $11.12 \pm 0.10$ | $3.85 e+10_{4.78 e+10}^{3.85 e+10}$ | $2.64 e+10_{3.26 e+10}^{2.24 e+10}$ | $1.21 e+10_{1.52 e+10}^{1.2 l e+10}$ | $3.85 e+10_{4.73 e+11}^{3.85 e+10}$ |
| 1884 | 25 | $15.83 \pm 0.10$ | $15.64 \pm 0.10$ | $1.58 e+09_{1.58 e+09}^{2.52 e+08}$ | $8.40 e+07_{4.62 e+08}^{7.06+07}$ | $1.49 e+09_{1.49 e+09}^{1.43 e+08}$ | $1.18 e+08_{6.01 e+09}^{1.18 e+08}$ |
| 1885 | 3 | $15.59 \pm 0.06$ | $15.59 \pm 0.10$ | $5.77 e+08_{1.45 e+09}^{5}$ | $2.54 e+00_{2.54 e+08}^{4.82 e+08}$ | $3.23 e+08_{1.24 e+09}^{3.23 e+08}$ | $3.61 e+08_{2.86 e+09}^{3.610+08}$ |
| 1898 | 25 | $14.99 \pm 0.05$ | $14.99 \pm 0.10$ | $8.10 e+08_{7.24 e+09}^{8.10+0)}$ | $5.38 e+08_{5.39 e+08}^{3.86 e+08}$ | $2.73 e+08_{6.85 e+09}^{2.73 e+08}$ | $7.73 e+08_{5.53 e+09}^{7.733+08}$ |

continued

| VCC | $\begin{gathered} \hline \hline \text { Type- } \\ \text { ID } \end{gathered}$ | $\begin{gathered} \text { mag }_{\mathrm{r}, \mathrm{bs}} \\ {[\mathrm{mag}]} \end{gathered}$ | $\begin{gathered} \mathrm{mag}_{\mathrm{r}, \mathrm{mod}} \\ {[\mathrm{mag}]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \mathrm{M}_{\text {nonvis }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 1900 | 25 | $15.48 \pm 0.06$ | $15.54 \pm 0.10$ | $3.53 e+09_{3.53 e+09}^{4.35+08}$ | $1.88 e+00_{3.09 e+08}^{1.76 e+08}$ | $3.34 e+09_{3.34 e+09}^{1.56+08}$ | $2.65 e+08_{4.58 e+09}^{2.65 e+08}$ |
| 1902 | 19 | $12.34 \pm 0.05$ | $12.35 \pm 0.10$ | $1.05 e+10_{1.55 e+10}^{1.04+10}$ | $7.23 e+09_{1.04 e+10}^{6.93 e+09}$ | $3.30 e+09_{4.88 e+09}^{3.30 e+09}$ | $1.05 e+10_{2.05 e+11}^{1.05 e+10}$ |
| 1905 | 25 | $16.74 \pm 0.17$ | $16.76 \pm 0.10$ | $3.59 e+08_{1.98 e+09}^{2.26 e+08}$ | $1.00 e+08_{2.02 e+08}^{1.00 e+08}$ | $2.58 e+08_{1.87 e+09}^{7}$ | $1.38 e+08_{3.87 e+09}^{2.158 e+08}$ |
| 1906 | 19 | $15.00 \pm 0.05$ | $15.00 \pm 0.10$ | $9.14 e+08_{7.77 e+09}^{8.71 e+08}$ | $6.28 e+08_{9.05 s e+08}^{4.14 e+08}$ | $2.86 e+08_{7.36 e+09}^{2.86+08}$ | $9.09 e+08_{1.80 e+10}^{99.09 e+08}$ |
| 1918 | 3 | $16.09 \pm 0.05$ | $16.11 \pm 0.10$ | $1.03 e+09_{1.03 e+09}^{1.03 z+09}$ | $5.51 e+07_{5.51 e+07}^{5.51 e+07}$ | $9.78 e+08_{9.78 e+08}^{99.78 e+08}$ | $7.75 e+07_{2.32 e+08}^{7.75 e+07}$ |
| 1920 | 19 | $14.68 \pm 0.05$ | $14.71 \pm 0.10$ | $9.84 e+08_{8.63 e+09}^{9.84 e+08}$ | $6.53 e+08_{6.53 e+08}^{4.60 e+08}$ | $3.31 e+08_{8.17 e+09}^{3.31 e+08}$ | $9.39 e+08_{5.95 e+09}^{9.39 e+08}$ |
| 1923 | 11 | $12.06 \pm 0.05$ | $12.07 \pm 0.10$ | $9.99 e+09_{8.08 e+10}^{9.99 e+09}$ | $6.03 e+09_{6.03 e+09}^{4.31 e+09}$ | $3.96 e+09_{7.65 e+10}^{3.96+09}$ | $8.62 e+09_{5.40 e+10}^{8.62 e+09}$ |
| 1929 | 7 | $12.88 \pm 0.05$ | $12.89 \pm 0.10$ | $3.31 e+10_{3.31 e+10}^{5.24+09}$ | $1.77 e+09_{2.31 e+09}^{1.77 e+09}$ | $3.13 e+10_{3.13 e+10}^{2.93 e+10}$ | $2.48 e+09_{1.67 e+10}^{2.48 e+09}$ |
| 1931 | 3 | $15.08 \pm 0.05$ | $15.09 \pm 0.10$ | $2.74 e+09_{2.74 e+09}^{4.97 e+08}$ | $1.46 e+08_{1.62 e+08}^{1.39+08}$ | $2.60 e+09_{2.60 e+09}^{3.58 e+08}$ | $2.06 e+08_{1.26 e+09}^{1.206 e+08}$ |
| 1932 | 7 | $12.13 \pm 0.05$ | $12.10 \pm 0.10$ | $1.11 e+10_{1.00 e+11}^{1.11 e+10}$ | $7.40 e+09_{7}^{5.362 e++09}$ | $3.75 e+09_{9.47 e+10}^{3.75 e+09}$ | $1.06 e+10_{7.63 e+10}^{1.106 e+10}$ |
| 1933 | 17 | $15.98 \pm 0.07$ | $16.00 \pm 0.10$ | $1.32 e+09_{1.32 e+09}^{2.43 e+08}$ | $7.05 e+07_{9.12 e+07}^{6.81 e+07}$ | $1.25 e+09_{1.25 e+09}^{1.75+08}$ | $9.90 e+07_{8.959+09}^{9.90 e+08}$ |
| 1943 | 5 | $11.29 \pm 0.05$ | $11.27 \pm 0.10$ | $2.33 e+10_{2.07 e+11}^{2.33 e+10}$ | $1.55 e+10_{1.55 e+10}^{1.10 e+10}$ | $7.85 e+09_{1.96 e+11}^{7.85 e+9}$ | $2.23 e+10_{1.20 e+11}^{2.23 e+10}$ |
| 1944 | 31 | $16.19 \pm 0.05$ | $16.16 \pm 0.10$ | $3.67 e+08_{4.81 e+08}^{3.66+08}$ | $2.52 e+08_{3.28 e+08}^{2.48+08}$ | $1.15 e+08_{1.54 e+08}^{1.15 e+08}$ | $3.64 e+08_{5.49 e+09}^{3.64 e+08}$ |
| 1952 | 3 | $15.82 \pm 0.05$ | $15.82 \pm 0.10$ | $2.58 e+08_{2.58 e+08}^{2.588+08}$ | $7.22 e+07_{7}^{7.22 e+027}$ | $1.86 e+08_{1.86 e+08}^{1.86+08}$ | $9.91 e+07.9 .91+e+07$ |
| 1955 | 30 | $12.74 \pm 0.05$ | $12.79 \pm 0.10$ | $4.56 e+10_{4.56 e+10}^{5.63 e+09}$ | $2.43 e+09^{2} .488 e+09$ | $4.32 e+10_{4.32 e+109}^{1.97 e+09}$ | $3.42 e+09_{4.75 e+10}^{3.42 e+09}$ |
| 1960 | 29 | $17.39 \pm 0.06$ | $17.37 \pm 0.10$ | $4.52 e+08_{4.52 e+07}^{7.20 e+07}$ | $2.41 e+07_{3.17 e+07}^{2}$ | $4.28 e+08_{4.28 e+07}^{4.032+07}$ | $3.39 e+07_{3.48 e+08}^{3.39 e+07}$ |
| 1965 | 3 | $15.88 \pm 0.06$ | $15.90 \pm 0.10$ | $2.11 e+09_{2.112 e+09}^{3.32 e+08}$ | $1.13 e+08_{1.46 e+08}^{1.13 e+08}$ | $2.00 e+09_{2.00 e+09}^{1.86 e+08}$ | $1.58 e+08_{1.45 e+09}^{1.58 e+08}$ |
| 1970 | 3 | $0.00 \pm-99999.00$ | $198.00 \pm 99.00$ | $0.00 e+00_{0.00 e+00}^{0.00 e+00}$ | $0.00 e+00_{0.00 e e+00}^{0.00 e+00}$ | $0.00 e+00_{0.00 e++00}^{0.000+0}$ | $6.95 e-310_{6.95 e-310}^{6.95 e-310}$ |
| 1987 | 7 | $10.23 \pm 0.05$ | $10.24 \pm 0.10$ | $4.49 e+11_{4.49 e+11}^{5.55 e+10}$ | $2.40 e+10_{3.35 e+10}^{2.40 e+10}$ | $4.25 e+11_{4.26 e+11}^{2.20 e+10}$ | $3.37 e+10_{1.97 e+11}^{3.37 e+10}$ |
| 1992 | 3 | $15.28 \pm 0.05$ | $15.30 \pm 0.10$ | $2.15 e+09_{2.15 e+09}^{3.65 e+08}$ | $1.14 e+08_{1.73 e+08}^{1.02 e+08}$ | $2.03 e+09_{2.03 e+09}^{2.63 e+08}$ | $1.61 e+08_{1.38 e+09}^{1.61 e+08}$ |

continued

| VCC | $\begin{gathered} \hline \hline \text { Type- } \\ \text { ID } \end{gathered}$ | mag $_{\mathrm{r}, \mathrm{obs}}$ <br> [mag] | $\operatorname{mag}_{r, \bmod }$ [mag] | $\begin{gathered} \hline \mathrm{M}_{\mathrm{tot}} \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \mathrm{M}_{\text {stellar }} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ | $\begin{aligned} & \mathrm{M}_{\mathrm{gas}} \\ & {\left[\mathrm{M}_{\odot}\right]} \end{aligned}$ | $\begin{gathered} \mathrm{M}_{\mathrm{nonvis}} \\ {\left[\mathrm{M}_{\odot}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] |
| 1994 | 25 | $16.44 \pm 0.07$ | $16.37 \pm 0.10$ | $4.19 e+08_{3.55 c+09}^{3.90 e+08}$ | $2.88 e+08_{3.47 e+08}^{1.89+08}$ | $1.31 e+08_{3.36 e+09}^{1.22 e+08}$ | $4.19 e+08_{8.46 e+09}^{4.19 e+08}$ |
| 1999 | 17 | $11.93 \pm 0.05$ | $11.91 \pm 0.10$ | $2.17 e+10_{2.41 e+10}^{1.98 e+10}$ | $1.48 e+10_{1.64 e+10}^{1.36 e+10}$ | $6.85 e+09_{7.69 e+09}^{6.19 e+09}$ | $2.17 e+10_{2.89 e+11}^{2.17 e+10}$ |
| 2006 | 9 | $13.06 \pm 0.05$ | $13.11 \pm 0.10$ | $3.34 e+10_{3.34 e+10}^{4.12 e+09}$ | $1.78 e+09_{2.49 e+09}^{1.78+09}$ | $3.16 e+10_{3,16 e+10}^{1.64+09}$ | $2.50 e+09^{2.597 e+10}$ |
| 2007 | 29 | $15.20 \pm 0.05$ | $15.22 \pm 0.10$ | $3.80 e+09_{3.81 e+09}^{6.01 e+08}$ | $2.03 e+08_{2.65 e+08}^{2.03 e+08}$ | $3.60 e+09_{3.60 e+09}^{3.36+08}$ | $2.85 e+08_{2.92 e+09}^{2.85 e+08}$ |
| 2015 | 1 | $15.60 \pm 0.05$ | $15.58 \pm 0.10$ | $2.49 e+09_{2.49 e+09}^{3.93 e+08}$ | $1.33 e+08_{1.73 e+08}^{1.33+08}$ | $2.36 e+09_{2.36 e+09}^{2.20 e+08}$ | $1.87 e+08_{8.06 e+08}^{1.87 e+08}$ |
| 2023 | 7 | $13.32 \pm 0.05$ | $13.35 \pm 0.10$ | $1.76 e+10_{1.76 e+10}^{2.89+09}$ | $9.41 e+08_{1.27 e+09}^{9.40 e+08}$ | $1.67 e+10_{1.67 e+10}^{1.61 e+09}$ | $1.32 e+09_{1.04 e+10}^{1.32 e+09}$ |
| 2033 | 1 | $14.54 \pm 0.05$ | $14.54 \pm 0.10$ | $6.57 e+09_{6.588+09}^{1.04+09}$ | $3.51 e+08_{4.58 e+08}^{3.51 e+08}$ | $6.22 e+09_{6.23 e+09}^{5.82 e+08}$ | $4.93 e+08_{6.17 e+09}^{4.93 e+08}$ |
| 2034 | 3 | $15.10 \pm 0.05$ | $15.09 \pm 0.10$ | $3.90 e+09_{3.90 e+09}^{6.17 e+08}$ | $2.08 e+08_{2.72 e+08}^{2.08+08}$ | $3.69 e+09_{3.69 e+09}^{3.45 e+08}$ | $2.92 e+08_{2.68 e+09}^{2.92 e+08}$ |
| 2037 | 3 | $14.83 \pm 0.05$ | $14.86 \pm 0.10$ | $4.57 e+09_{4.57 e+09}^{7.29 e+08}$ | $2.44 e+08_{4.54 e+08}^{2.36+08}$ | $4.33 e+09_{4.33 e+09}^{2.99++08}$ | $3.43 e+08_{5.49 e+09}^{3.43 e+08}$ |
| 2058 | 7 | $10.59 \pm 0.05$ | $10.57 \pm 0.10$ | $4.73 e+10_{4.25 e+11}^{4.73 e+10}$ | $3.14 e+10_{3.17 e+10}^{2.26 e+10}$ | $1.59 e+10_{4.02 e+11}^{1.59 e+10}$ | $4.52 e+10_{2.44 e+11}^{4.52 e+10}$ |
| 2066 | 15 | $11.25 \pm 0.05$ | $11.25 \pm 0.10$ | $2.38 e+10_{2.09 e+11}^{2.38 e+10}$ | $1.58 e+10_{1.58 e+10}^{1.11 e+10}$ | $8.00 e+09_{1.98 e+11}^{8.00 e+09}$ | $2.27 e+10_{1.78 e+11}^{2.27 e+10}$ |
| 2070 | 17 | $10.28 \pm 0.05$ | $10.25 \pm 0.10$ | $9.54 e+10_{1.16 e+10}^{99.54+11}$ | $6.55 e+10_{7.92 e+10}^{6.55 e+10}$ | $2.99 e+10_{3.71 e+10}^{2.99 e+10}$ | $9.54 e+10_{1.16 e+12}^{9.54 e+10}$ |
| 2089 | 31 | $15.04 \pm 0.05$ | $15.08 \pm 0.10$ | $1.88 e+09_{2.09 e+09}^{1.71 e+09}$ | $1.29 e+09_{1.42 e+09}^{1.18 e+09}$ | $5.94 e+08_{6.67 e+08}^{5.37+08}$ | $1.88 e+09^{1.380 e+10} 10$ |
| 2090 | 25 | $15.00 \pm 0.05$ | $14.99 \pm 0.10$ | $5.38 e+09_{5.39 e+09}^{6.71 e+08}$ | $2.87 e+08_{4.05 e+08}^{2.87 e+08}$ | $5.10 e+09_{5.10 e+09}^{2.66+08}$ | $4.04 e+08_{4.4 l e+09}^{4.04+08}$ |
| 2094 | 3 | $17.68 \pm 0.12$ | $17.59 \pm 0.10$ | $9.06 e+07_{5.04 e+07}^{7.88 e+08}$ | $2.53 e+07_{3.47 e+07}^{2.53 e+07}$ | $6.53 e+07_{4.78 e+08}^{4.41 e+07}$ | $3.48 e+07_{4.14 e+08}^{3.48+07}$ |

Results from GALEV/GAZELLE runs (Part I).

Table 6.5: Results from GALEV/GAZELLE runs (Part II).
Upper and lower values correspond to the minimum and maximum values of the GAZELLE output.

| VCC | Type- | $\mathrm{Z}_{\text {gas }}$ | SFR | $\tau$ | $\mathrm{E}(\mathrm{B}-\mathrm{V})$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | ID |  | $\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]$ | $[\mathrm{Gyr}]$ | $[\mathrm{mag}]$ |
| $[1]$ | $[2]$ | $[3]$ | $[4]$ | $[5]$ | $[6]$ |
| 0001 | 1 | $0.030_{0.030}^{0.001}$ | $0.022_{0.159}^{0.022}$ | $5_{25}^{5}$ | $0.0025_{0.2804}^{0.0025}$ |
| 0004 | 3 | $0.001_{0.007}^{0.001}$ | $0.018_{0.033}^{0.017}$ | $25_{25}^{4}$ | $0.0153_{0.0648}^{0.0000}$ |
| 0010 | 1 | $0.001_{0.025}^{0.001}$ | $0.097_{0.097}^{0.032}$ | $25_{25}^{6}$ | $0.1877_{0.1800}^{0.0000}$ |
| 0015 | 7 | $0.001_{0.025}^{0.001}$ | $0.144_{0.144}^{0.046}$ | $25_{25}^{6}$ | $0.2002_{0.2002}^{0.0012}$ |
| 0017 | 3 | $0.004_{0.007}^{0.001}$ | $0.048_{0.091}^{0.048}$ | $15_{25}^{4}$ | $0.0093_{0.0737}^{0.0093}$ |
| 0020 | 24 | $0.007_{0.014}^{0.004}$ | $0.020_{0.020}^{0.007}$ | $4_{15}^{4}$ | $0.2008_{0.2009}^{0.0506}$ |
| 0022 | 1 | $0.001_{0.014}^{0.001}$ | $0.050_{0.050}^{0.031}$ | $25_{25}^{8}$ | $0.1473_{0.2276}^{0.0386}$ |
| 0024 | 1 | $0.001_{0.014}^{0.001}$ | $0.100_{0.100}^{0.063}$ | $25_{25}^{8}$ | $0.1434_{0.1434}^{0.0356}$ |
| 0025 | 7 | $0.001_{0.001}^{0.001}$ | $1.352_{1.353}^{1.352}$ | $25_{25}^{25}$ | $0.0998_{0.0998}^{0.0998}$ |
| 0026 | 3 | $0.007_{0.009}^{0.001}$ | $0.013_{0.013}^{0.007}$ | $4_{25}^{4}$ | $0.0552_{0.0552}^{0.0000}$ |
| 0030 | 3 | $0.007_{0.009}^{0.001}$ | $0.042_{0.042}^{0.022}$ | $4_{25}^{4}$ | $0.0587_{0.0588}^{0.0000}$ |
| 0031 | 24 | $0.001_{0.014}^{0.001}$ | $0.065_{0.065}^{0.042}$ | $25_{25}^{8}$ | $0.1027_{0.1300}^{0.0000}$ |
| 0034 | 7 | $0.001_{0.009}^{0.001}$ | $0.123_{0.123}^{0.099}$ | $25_{25}^{10}$ | $0.0621_{0.0621}^{0.0000}$ |
| 0041 | 3 | $0.001_{0.030}^{0.001}$ | $0.027_{0.049}^{0.006}$ | $25_{25}^{4}$ | $0.1219_{0.1600}^{0.0000}$ |
| 0048 | 9 | $0.001_{0.009}^{0.001}$ | $0.161_{0.161}^{0.137}$ | $25_{25}^{10}$ | $0.0431_{0.0431}^{0.0000}$ |

continued

| VCC | TypeID | $\mathrm{Z}_{\text {gas }}$ | $\begin{gathered} \mathrm{SFR} \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\tau$ [Gyr] | $\begin{gathered} \mathrm{E}(\mathrm{~B}-\mathrm{V}) \\ {[\mathrm{mag}]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] |
| 0052 | 3 | $0.001_{0.014}^{0.001}$ | $0.014_{0.014}^{0.009}$ | $25^{8}$ | $0.1370_{0.1370}^{0.0254}$ |
| 0058 | 5 | $0.001_{0.025}^{0.001}$ | $1.038_{1.038}^{0.339}$ | $25^{6}{ }^{6}$ | $0.1895_{0.1895}^{0.0000}$ |
| 0066 | 7 | $0.001_{0.014}^{0.001}$ | $2.6922_{2.693}^{1.734}$ | $25_{25}^{8}$ | $0.0984_{0.0984}^{0.0000}$ |
| 0067 | 7 | $0.001_{0.009}^{0.001}$ | $0.302_{0.302}^{2.240}$ | $25_{25}^{10}$ | $0.0948_{0.0948}^{0.0301}$ |
| 0073 | 5 | $0.030_{0.033}^{0.001}$ | $0.372_{2.688}^{0.069}$ | $5_{25}^{4}$ | $0.1210_{0.3980}^{0.0388}$ |
| 0074 | 1 | $0.001_{0.014}^{0.001}$ | $0.044_{0.044}^{0.028}$ | $25^{8}$ | $0.1746_{0.1747}^{0.0661}$ |
| 0081 | 7 | $0.001_{0.009}^{0.001}$ | $0.110_{0.110}^{0.094}$ | $25_{25}^{10}$ | $0.0417_{0.0417}^{0.0000}$ |
| 0083 | 3 | $0.001_{0.014}^{0.001}$ | $0.079_{0.079}^{0.050}$ | $25_{25}^{8}$ | $0.1327_{0.1328}^{0.0245}$ |
| 0085 | 3 | $0.009_{0.014}^{0.001}$ | $0.011_{0.014}^{0.010}$ | $10_{25}^{8}$ | $0.00988_{0.0764}^{0.0000}$ |
| 0087 | 11 | $0.001_{0.009}^{0.001}$ | $0.131_{0.131}^{0.111}$ | $25_{25}^{10}$ | $0.0444_{0.0445}^{0.0000}$ |
| 0089 | 7 | $0.001_{0.025}^{0.001}$ | $2.904_{2.904}^{0.924}$ | $25_{25}^{6}$ | $0.2145_{0.2145}^{0.0158}$ |
| 0092 | 5 | $0.033_{0.033}^{0.030}$ | $0.626_{3.382}^{0.000}$ | 45 | $0.0158_{0.0983}^{0.0000}$ |
| 0093 | 25 | $0.001_{0.014}^{0.001}$ | $0.018_{0.033}^{0.014}$ | $25_{25}^{4}$ | $0.0362_{0.0822}^{0.0000}$ |
| 0094 | 19 | $0.033_{0.033}^{0.032}$ | $0.001_{0.001}^{0.000}$ | $3{ }_{3}^{0}$ | $0.0340_{0.0340}^{0.0276}$ |
| 0097 | 7 | $0.030_{0.033}^{0.001}$ | $0.3433_{2.485}^{0.069}$ | $5_{25}^{4}$ | $0.0541_{0.3319}^{0.0000}$ |
| 0099 | 17 | $0.030_{0.033}^{0.001}$ | $0.041_{0.296}^{0.008}$ | $5_{25}^{4}$ | $0.0544_{0.3313}^{0.0000}$ |
| 0104 | 25 | $0.001_{0.033}^{0.001}$ | $0.009_{0.013}^{0.001}$ | $25_{25}^{4}$ | $0.0000_{0.0000}^{0.0000}$ |
| 0105 | 9 | $0.001_{0.004}^{0.001}$ | $0.502_{0.502}^{0.472}$ | $25_{25}^{15}$ | $0.1414_{0.1414}^{0.1193}$ |
| 0113 | 24 | $0.007_{0.009}^{0.004}$ | $0.028_{0.028}^{0.014}$ | $4_{15}^{4}$ | $0.0826_{0.0826}^{0.0000}$ |
| 0114 | 3 | $0.001_{0.030}^{0.001}$ | $0.047_{0.083}^{0.008}$ | $25_{25}^{4}$ | $0.19933_{0.2399}^{0.0000}$ |

continued

| VCC | TypeID | $\mathrm{Z}_{\text {gas }}$ | $\begin{gathered} \mathrm{SFR} \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\begin{gathered} \tau \\ {[\mathrm{Gyr}]} \end{gathered}$ | E(B-V) [mag] |
| :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] |
| 0117 | 3 | $0.001_{0.007}^{0.001}$ | $0.023_{0.038}^{0.023}$ | $25_{25}^{4}$ | $0.0000_{0.0161}^{0.0000}$ |
| 0119 | 7 | $0.001_{0.007}^{0.001}$ | $0.0633_{0.107}^{0.063}$ | $25_{25}^{4}$ | $0.0000_{0.0281}^{0.0000}$ |
| 0120 | 7 | $0.001_{0.001}^{0.001}$ | $0.830_{0.830}^{0.830}$ | $25_{25}^{25}$ | $0.1382_{0.1382}^{0.1382}$ |
| 0122 | 19 | $0.033_{0.033}^{0.033}$ | $0.001_{0.038}^{0.000}$ | $3{ }_{4}^{0}$ | $0.0000_{0.0216}^{0.0000}$ |
| 0126 | 9 | $0.001_{0.014}^{0.001}$ | $0.212_{0.212}^{0.133}$ | $25_{25}^{8}$ | $0.1267_{0.1267}^{0.0179}$ |
| 0128 | 25 | $0.001_{0.025}^{0.001}$ | $0.173_{0.173}^{0.055}$ | $25_{25}^{6}$ | $0.2771_{0.2771}^{0.0780}$ |
| 0130 | 1 | $0.001_{0.010}^{0.001}$ | $0.016_{0.028}^{0.011}$ | $25_{25}^{4}$ | $0.0145_{0.1023}^{0.0000}$ |
| 0131 | 7 | $0.030_{0.030}^{0.030}$ | $0.085_{0.085}^{0.085}$ | 55 | $0.0210_{0.0211}^{0.0210}$ |
| 0132 | 3 | $0.001_{0.009}^{0.001}$ | $0.027_{0.049}^{0.024}$ | $25_{25}^{4}$ | $0.0357_{0.0803}^{0.0000}$ |
| 0135 | 23 | $0.030_{0.033}^{0.001}$ | $0.070_{0.511}^{0.014}$ | $5_{25}^{4}$ | $0.0505_{0.3280}^{0.0000}$ |
| 0137 | 25 | $0.025_{0.025}^{0.001}$ | $0.012_{0.039}^{0.012}$ | $6{ }_{25}^{6}$ | $0.0309_{0.2300}^{0.0309}$ |
| 0143 | 7 | $0.001_{0.025}^{0.001}$ | $0.134_{0.134}^{0.043}$ | $25_{25}^{6}$ | $0.2152_{0.2152}^{0.0163}$ |
| 0144 | 1 | $0.010_{0.016}^{0.009}$ | $0.032_{0.074}^{0.028}$ | $10_{10}^{8}$ | $0.0064_{0.0700}^{0.0000}$ |
| 0145 | 7 | $0.001_{0.025}^{0.001}$ | $1.851_{1.851}^{0.589}$ | $25_{25}^{6}$ | $0.2299_{0.2299}^{0.0313}$ |
| 0152 | 7 | $0.033^{0.033}$ | $0.057{ }_{0.305}^{0.001}$ | $4{ }_{5}^{3}$ | $0.0244_{0.1068}^{0.0000}$ |
| 0155 | 25 | $0.001_{0.030}^{0.001}$ | $0.196_{0.196}^{0.031}$ | $25_{25}^{5}$ | $0.2314_{0.2314}^{0.0000}$ |
| 0157 | 7 | $0.030_{0.030}^{0.001}$ | $0.8966_{6.347}^{0.895}$ | $5_{25}^{5}$ | $0.0000_{0.2700}^{0.0000}$ |
| 0159 | 3 | $0.007_{0.007}^{0.004}$ | $0.078_{0.078}^{0.043}$ | $4_{15}^{4}$ | $0.0559_{0.0559}^{0.0000}$ |
| 0162 | 9 | $0.001_{0.004}^{0.001}$ | $0.146_{0.146}^{0.138}$ | $25_{25}^{15}$ | $0.0639_{0.0639}^{0.0421}$ |
| 0166 | 19 | $0.033_{0.033}^{0.032}$ | $0.002_{0.002}^{0.000}$ | $3{ }_{3}^{0}$ | $0.0532_{0.0532}^{0.0465}$ |

continued

| VCC | TypeID | $Z_{\text {gas }}$ | $\begin{gathered} \mathrm{SFR} \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\tau$ <br> [Gyr] | $\begin{gathered} \mathrm{E}(\mathrm{~B}-\mathrm{V}) \\ {[\mathrm{mag}]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] |
| 0167 | 5 | $0.0333_{0.033}^{0.032}$ | $0.020_{0.020}^{0.000}$ | $3_{3}^{0}$ | $0.1329_{0.1329}^{0.1265}$ |
| 0168 | 3 | $0.001_{0.025}^{0.001}$ | $0.033_{0.033}^{0.010}$ | $25_{25}^{6}$ | $0.2199_{0.2199}^{0.0205}$ |
| 0169 | 3 | $0.007_{0.007}^{0.001}$ | $0.030_{0.030}^{0.019}$ | $4_{25}^{4}$ | $0.0000_{0.0000}^{0.0000}$ |
| 0172 | 1 | $0.001_{0.010}^{0.001}$ | $0.121_{0.214}^{0.082}$ | $25_{25}^{4}$ | $0.0421_{0.1281}^{0.0000}$ |
| 0181 | 25 | $0.001_{0.014}^{0.001}$ | $0.008_{0.008}^{0.005}$ | $25_{25}^{8}$ | $0.1088_{0.1088}^{0.0000}$ |
| 0187 | 7 | $0.001_{0.025}^{0.001}$ | $1.060_{1.060}^{0.338}$ | $25_{25}^{6}$ | $0.29744_{0.2974}^{0.0992}$ |
| 0190 | 25 | $0.004_{0.030}^{0.001}$ | $0.026_{0.050}^{0.005}$ | $15_{25}^{4}$ | $0.15977_{0.2250}^{0.0000}$ |
| 0199 | 17 | $0.0333_{0.033}^{0.032}$ | $0.003_{0.003}^{0.000}$ | $3_{3}^{0}$ | $0.0934_{0.0935}^{0.0871}$ |
| 0207 | 1 | $0.010_{0.010}^{0.009}$ | $0.007{ }_{0.008}^{0.006}$ | $10_{10}^{10}$ | $0.0000_{0.0782}^{0.0000}$ |
| 0213 | 23 | $0.001_{0.025}^{0.001}$ | $0.420_{0.420}^{0.134}$ | $25_{25}^{6}$ | $0.2214_{0.2214}^{0.0227}$ |
| 0217 | 3 | $0.007_{0.009}^{0.001}$ | $0.073_{0.073}^{0.042}$ | $4_{25}^{4}$ | $0.0304_{0.0304}^{0.0000}$ |
| 0221 | 7 | $0.001_{0.004}^{0.001}$ | $0.867_{0.867}^{0.814}$ | $25_{25}^{15}$ | $0.1492_{0.1492}^{0.1271}$ |
| 0222 | 17 | $0.0333_{0.033}^{0.032}$ | $0.004_{0.004}^{0.000}$ | $3_{3}^{0}$ | $0.0926_{0.0926}^{0.0863}$ |
| 0223 | 1 | $0.001_{0.014}^{0.001}$ | $0.041_{0.041}^{0.026}$ | $25_{25}^{8}$ | $0.1260_{0.1261}^{0.0179}$ |
| 0226 | 7 | $0.030_{0.030}^{0.001}$ | $0.6611_{4.797}^{0.661}$ | $5_{25}^{5}$ | $0.0584_{0.3362}^{0.0584}$ |
| 0234 | 17 | $0.0333_{0.033}^{0.032}$ | $0.000_{0.002}^{0.000}$ | 23 | $0.0330_{0.0361}^{0.0297}$ |
| 0237 | 29 | $0.007_{0.007}^{0.004}$ | $0.025_{0.025}^{0.014}$ | $4_{15}^{4}$ | $0.0367_{0.0367}^{0.0000}$ |
| 0241 | 9 | $0.007_{0.007}^{0.001}$ | $0.187_{0.187}^{0.121}$ | $4_{25}^{4}$ | $0.0000_{0.0000}^{0.0000}$ |
| 0247 | 3 | $0.004_{0.033}^{0.001}$ | $0.023_{0.043}^{0.000}$ | $15_{25}^{2}$ | $0.2507_{0.7669}^{0.0000}$ |
| 0252 | 25 | $0.030_{0}^{0.0033}$ | $0.001_{0.015}^{0.000}$ | $5_{25}^{4}$ | $0.0000_{0.3004}^{0.0000}$ |

continued

|  | VCC | $\begin{gathered} \hline \hline \text { Type- } \\ \text { ID } \end{gathered}$ | $\mathrm{Z}_{\text {gas }}$ | $\begin{gathered} \text { SFR } \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\begin{gathered} \tau \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{E}(\mathrm{~B}-\mathrm{V}) \\ {[\mathrm{mag}]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | [1] | [2] | [3] | [4] | [5] | [6] |
|  | 0260 | 3 | $0.001_{0.014}^{0.001}$ | $0.028_{0.028}^{0.018}$ | $25_{25}^{8}$ | $0.0950_{0.0950}^{0.0000}$ |
|  | 0267 | 5 | $0.001_{0.014}^{0.001}$ | $0.567_{0.567}^{0.355}$ | $25_{25}^{8}$ | $0.1198_{0.1199}^{0.0119}$ |
|  | 0274 | 1 | $0.001_{0.030}^{0.001}$ | $0.012_{0.022}^{0.004}$ | $25_{25}^{4}$ | $0.0138_{0.1014}^{0.0000}$ |
|  | 0275 | 25 | $0.001_{0.014}^{10.001}$ | $0.252_{0.252}^{0.158}$ | $25_{25}^{8}$ | $0.1569_{0.1569}^{0.0491}$ |
|  | 0282 | 25 | $0.001_{0.014}^{0.001}$ | $0.015_{0.015}^{0.009}$ | $25_{25}^{8}$ | $0.1304_{0.1305}^{0.0194}$ |
|  | 0286 | 3 | $0.001_{0.014}^{0.001}$ | $0.035_{0.035}^{0.022}$ | $25_{25}^{8}$ | $0.1230_{0.1231}^{0.0149}$ |
|  | 0289 | 7 | $0.001_{0.004}^{10.001}$ | $0.190_{0.190}^{0.179}$ | $25_{25}^{15}$ | $0.0877_{0.0878}^{0.0658}$ |
|  | 0293 | 25 | $0.014_{0.025}^{40.001}$ | $0.060_{0.168}^{0.031}$ | $8{ }_{25}^{4}$ | $0.1153_{0.2621}^{0.0283}$ |
| N | 0304 | 25 | $0.001_{0.014}^{10.001}$ | $0.035_{0.035}^{0.021}$ | $25_{25}^{8}$ | $0.1364_{0.1364}^{0.0255}$ |
|  | 0307 | 7 | $0.001_{0.014}^{10.001}$ | $15.659_{15.609}^{9.989}$ | $25_{25}^{8}$ | $0.1533_{0.1540}^{0.0465}$ |
|  | 0309 | 3 | $0.001_{0.009}^{0.001}$ | $0.032_{0.032}^{0.028}$ | $25_{25}^{10}$ | $0.0370_{0.0370}^{0.0000}$ |
|  | 0312 | 19 | $0.033_{0.033}^{0.032}$ | $0.001{ }_{0.001}^{0.000}$ | $3{ }_{3}^{0}$ | $0.0754_{0.0754}^{0.0691}$ |
|  | 0318 | 7 | $0.004_{0.004}^{0.001}$ | $0.188_{0.200}^{0.188}$ | $15_{25}^{15}$ | $0.0027_{0.0245}^{0.0027}$ |
|  | 0320 | 3 | $0.007_{0.014}^{0.001}$ | $0.084_{0.084}^{0.030}$ | $4{ }_{25}^{4}$ | $0.2452_{0.2452}^{0.0950}$ |
|  | 0322 | 3 | $0.001_{0.014}^{10.001}$ | $0.167_{0.167}^{0.104}$ | $25_{25}^{8}$ | $0.1432_{0.1432}^{0.0349}$ |
|  | 0323 | 17 | $0.033_{0.033}^{0.032}$ | $0.021_{0.021}^{0.000}$ | $4{ }_{4}^{0}$ | $0.0558_{0.0559}^{0.0173}$ |
|  | 0324 | 1 | $0.001_{0.004}^{10.001}$ | $0.228_{0.228}^{0.215}$ | $25_{25}^{15}$ | $0.0436_{0.0436}^{0.0217}$ |
|  | 0328 | 3 | $0.001_{0.009}^{0.001}$ | $0.031{ }_{0.031}^{0.026}$ | $25_{25}^{10}$ | $0.0488_{0.0489}^{0.0000}$ |
|  | 0329 | 3 | $0.001_{0.007}^{0.001}$ | $0.006_{0.010}^{0.006}$ | $25_{25}^{4}$ | $0.0000_{0.0359}^{0.0000}$ |
|  | 0334 | 1 | $0.001_{0.001}^{10.001}$ | $0.063_{0.063}^{0.063}$ | $25_{25}^{25}$ | $0.0542_{0.0542}^{0.0542}$ |

continued

| VCC | TypeID | $Z_{\text {gas }}$ | $\begin{gathered} \mathrm{SFR} \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\tau$ <br> [Gyr] | $\begin{gathered} \mathrm{E}(\mathrm{~B}-\mathrm{V}) \\ {[\mathrm{mag}]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] |
| 0340 | 1 | $0.001_{0.001}^{0.001}$ | $0.186_{0.186}^{0.186}$ | $25_{25}^{25}$ | $0.1233_{0.1233}^{0.123}$ |
| 0341 | 17 | $0.0333_{0.033}^{0.032}$ | $0.004_{0.004}^{0.000}$ | $3_{3}^{0}$ | $0.0754_{0.0754}^{0.069}$ |
| 0342 | 19 | $0.033_{0.033}^{0.032}$ | $0.001_{0.001}^{0.000}$ | $3_{3}^{0}$ | $0.0478_{0.0478}^{0.0414}$ |
| 0343 | 9 | $0.001_{0.014}^{0.001}$ | $0.115_{0.115}^{0.072}$ | $25_{25}^{8}$ | $0.1379_{0.1380}^{0.0296}$ |
| 0350 | 3 | $0.001_{0.014}^{0.001}$ | $0.030_{0.030}^{0.019}$ | $25_{25}^{8}$ | $0.1001_{0.1002}^{0.0000}$ |
| 0354 | 25 | $0.030_{0.033}^{0.001}$ | $0.022_{0.157}^{0.000}$ | $5_{25}^{2}$ | $0.0148_{0.2888}^{0.0000}$ |
| 0358 | 17 | $0.033_{0.033}^{0.032}$ | $0.001_{0.057}^{0.000}$ | $3_{4}^{0}$ | $0.0550_{0.0866}^{0.0486}$ |
| 0364 | 3 | $0.007_{0.025}^{0.001}$ | $0.052_{0.052}^{0.010}$ | $4_{25}^{4}$ | $0.2704_{0.2704}^{0.0378}$ |
| 0366 | 19 | $0.030_{0.033}^{0.001}$ | $0.101_{1.293}^{0.023}$ | $5_{25}^{4}$ | $0.0114_{0.3312}^{0.0000}$ |
| 0367 | 3 | $0.007_{0.033}^{0.001}$ | $0.277_{0.277}^{0.000}$ | $4_{25}^{3}$ | $0.3199_{0.3199}^{0.0000}$ |
| 0371 | 19 | $0.033_{0.033}^{0.032}$ | $0.001_{0.001}^{0.000}$ | $3_{3}^{0}$ | $0.0916_{0.0916}^{0.0852}$ |
| 0373 | 19 | $0.033_{0.033}^{0.033}$ | $0.000_{0.356}^{0.000}$ | $2{ }_{5}^{0}$ | $0.0300_{0.1443}^{0.0266}$ |
| 0375 | 19 | $0.033_{0.033}^{0.032}$ | $0.002_{0.002}^{0.000}$ | $3{ }_{3}^{0}$ | $0.0385_{0.0385}^{0.0319}$ |
| 0379 | 3 | $1.000_{1.000}^{1.000}$ | $0.000_{0.000}^{0.000}$ | $0_{0}^{0}$ | $0.0000_{0.0000}^{0.0000}$ |
| 0381 | 3 | $0.001_{0.025}^{0.001}$ | $0.029_{0.052}^{0.013}$ | $25_{25}^{4}$ | $0.0908_{0.1325}^{0.0000}$ |
| 0382 | 7 | $0.001_{0.001}^{0.001}$ | $2.195_{2.196}^{2.195}$ | $25_{25}^{25}$ | $0.1403_{0.1403}^{0.1403}$ |
| 0386 | 17 | $0.033_{0.033}^{0.032}$ | $0.000_{0.001}^{0.000}$ | 23 | $0.0221_{0.0252}^{0.0188}$ |
| 0393 | 7 | $0.001_{0.025}^{0.001}$ | $0.854_{0.854}^{0.277}$ | $25_{25}^{6}$ | $0.1920_{0.1920}^{0.0000}$ |
| 0404 | 7 | $0.030_{0.033}^{0.001}$ | $0.065_{0.468}^{0.012}$ | $5_{25}^{4}$ | $0.0772_{0.3545}^{0.00000}$ |
| 0408 | 19 | $0.0333_{0.033}^{0.032}$ | $0.006_{0.006}^{0.000}$ | $3{ }_{3}^{0}$ | $0.0992_{0.0992}^{0.0925}$ |

continued

| VCC | $\begin{gathered} \hline \hline \text { Type- } \\ \text { ID } \end{gathered}$ | $\mathrm{Z}_{\text {gas }}$ | $\begin{gathered} \mathrm{SFR} \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\begin{gathered} \tau \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \mathrm{E}(\mathrm{~B}-\mathrm{V}) \\ {[\mathrm{mag}]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] |
| 0409 | 25 | $0.001_{0.014}^{0.001}$ | $0.008_{0.008}^{0.005}$ | $25^{8}$ | $0.1203_{0.1204}^{0.0086}$ |
| 0410 | 1 | $0.010_{0.017}^{0.001}$ | $0.006_{0.018}^{0.005}$ | $10_{25}^{4}$ | $0.0000_{0.1181}^{0.0000}$ |
| 0413 | 3 | $0.001_{0.025}^{10.001}$ | $0.016_{0.016}^{0.006}$ | $25_{25}^{6}$ | $0.16400_{0.1640}^{0.0000}$ |
| 0414 | 25 | $0.030_{0.030}^{0.001}$ | $0.004_{0.029}^{0.004}$ | $5{ }_{25}^{5}$ | $0.0000_{0.2735}^{0.0000}$ |
| 0415 | 9 | $0.001_{0.014}^{0.001}$ | $0.153_{0.153}^{0.096}$ | $25^{8}$ | $0.1681_{0.1681}^{0.0600}$ |
| 0423 | 3 | $0.007_{0.033}^{7.001}$ | $0.302_{0.302}^{20.000}$ | $4{ }_{25}^{0.5}$ | $0.3796_{0.3797}^{0.0000}$ |
| 0425 | 3 | $0.033_{0.033}^{0.001}$ | $0.000_{0.046}^{0.000}$ | $3{ }_{25}^{0}$ | $0.2543^{0.24539}$ |
| 0428 | 1 | $0.009_{0.031}^{0.009}$ | $0.009_{0.024}^{\text {9.004 }}$ | $10_{10}^{6}$ | $0.0000_{0.0000}^{0.0000}$ |
| 0429 | 29 | $0.007_{0.330}^{70.004}$ | $0.071_{0.071}^{10.012}$ | $4{ }_{15}^{4}$ | $0.2101_{0.2101}^{0.0602}$ |
| 0446 | 29 | $0.014_{0.014}^{0.001}$ | $0.052_{0.084}^{0.052}$ | $8{ }_{25}^{8}$ | $0.0540_{0.1626}^{0.0540}$ |
| 0448 | 3 | $0.001_{0.009}^{0.001}$ | $0.019_{0.019}^{0.016}$ | $25_{25}^{10}$ | $0.0587_{0.0587}^{0.0000}$ |
| 0449 | 5 | $0.014_{0.025}^{0.009}$ | $0.6788_{0.861}^{0.347}$ | $8{ }_{10}^{6}$ | $0.3230_{0.3659}^{0.2338}$ |
| 0450 | 19 | $0.030_{0.030}^{0.001}$ | $0.028_{0.199}^{0.028}$ | $5{ }_{25}^{5}$ | $0.0482_{0.3243}^{0.0482}$ |
| 0453 | 11 | $0.001_{0.025}^{10.001}$ | $0.053_{0.053}^{0.017}$ | $25_{25}^{6}$ | $0.1900_{0.1900}^{0.0000}$ |
| 0459 | 1 | $0.001_{0.001}^{0.001}$ | $0.114_{0.114}^{0.114}$ | $25_{25}^{25}$ | $0.0352_{0.0353}^{0.0352}$ |
| 0460 | 17 | $0.033_{0.033}^{30.032}$ | $0.010_{0.484}^{0.000}$ | $3{ }_{4}^{0}$ | $0.0116_{0.0439}^{0.0053}$ |
| 0465 | 7 | $0.001_{0.004}^{0.001}$ | $1.234_{1.234}^{1.161}$ | $25_{25}^{15}$ | $0.0462_{0.0462}^{0.0244}$ |
| 0472 | 25 | $0.033_{0.033}^{0.001}$ | $0.001_{0.041}^{0.000}$ | $4{ }_{25}^{0}$ | $0.0000_{0.2936}^{0.0000}$ |
| 0476 | 3 | $0.001_{0.030}^{10.001}$ | $0.011_{0.020}^{10.002}$ | $25_{25}^{4}$ | $0.1284_{0.1732}^{0.0000}$ |
| 0477 | 3 | $0.007_{0.014}^{7.001}$ | $0.038_{0.039}^{0.013}$ | $4{ }_{25}^{4}$ | $0.1607_{0.1607}^{0.0056}$ |

continued

| VCC | TypeID | $\mathrm{Z}_{\text {gas }}$ | $\begin{gathered} \mathrm{SFR} \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\begin{gathered} \tau \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \mathrm{E}(\mathrm{~B}-\mathrm{V}) \\ {[\mathrm{mag}]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] |
| 0479 | 3 | $0.001_{0.014}^{0.001}$ | $0.038_{0.038}^{0.024}$ | $25_{25}^{8}$ | $0.1234_{0.1234}^{0.0140}$ |
| 0483 | 7 | $1.000_{1.000}^{1.000}$ | $0.000_{0.000}^{0.000}$ | $0_{0}^{0}$ | $0.0000_{0.0000}^{0.0000}$ |
| 0488 | 3 | $0.001_{0.025}^{0.001}$ | $0.018_{0.018}^{0.007}$ | $25_{25}^{6}$ | $0.1519_{0.1519}^{0.0000}$ |
| 0491 | 7 | $0.001_{0.007}^{0.001}$ | $0.774_{1.284}^{0.774}$ | $25_{25}^{4}$ | $0.0000_{0.0222}^{0.0000}$ |
| 0494 | 25 | $0.001_{0.033}^{0.001}$ | $0.115_{0.206}^{0.000}$ | $25_{25}^{3}$ | $0.21899_{0.2591}^{0.0000}$ |
| 0497 | 7 | $1.000_{1.000}^{1.000}$ | $0.000_{0.000}^{0.000}$ | $0_{0}^{0}$ | $0.0000_{0.0000}^{0.0000}$ |
| 0509 | 9 | $0.009_{0.009}^{0.001}$ | $0.0699_{0.087}^{0.069}$ | $10_{25}^{10}$ | $0.0105_{0.0749}^{0.0105}$ |
| 0512 | 11 | $0.001_{0.014}^{0.001}$ | $0.069_{0.069}^{0.046}$ | $25_{25}^{8}$ | $0.0875_{0.0875}^{0.0000}$ |
| 0513 | 1 | $0.001_{0.030}^{0.001}$ | $0.174_{0.174}^{0.026}$ | $25_{25}^{5}$ | $0.2480_{0.2801}^{0.0000}$ |
| 0514 | 7 | $0.009_{0.014}^{0.004}$ | $0.173_{0.384}^{0.136}$ | $10_{15}^{4}$ | $0.0890_{0.1948}^{0.0459}$ |
| 0520 | 3 | $0.001_{0.009}^{0.001}$ | $0.005_{0.009}^{0.004}$ | $25_{25}^{4}$ | $0.0703_{0.1141}^{0.0032}$ |
| 0522 | 17 | $0.033_{0.033}^{0.030}$ | $0.057_{0.298}^{0.000}$ | $4_{5}^{0}$ | $0.0000_{0.0714}^{0.0000}$ |
| 0524 | 5 | $0.033_{0.033}^{0.032}$ | $0.003_{0.003}^{0.000}$ | $3{ }_{3}^{0}$ | $0.0881_{0.0882}^{0.0816}$ |
| 0530 | 3 | $0.009_{0.014}^{0.001}$ | $0.084_{0.106}^{0.067}$ | $10_{25}^{8}$ | $0.03977_{0.1057}^{0.0000}$ |
| 0534 | 17 | $0.033_{0.033}^{0.032}$ | $0.000_{0.001}^{0.000}$ | 20 | $0.0045_{0.0076}^{0.0011}$ |
| 0546 | 25 | $0.001_{0.025}^{0.001}$ | $0.107_{0.107}^{0.034}$ | $25_{25}^{6}$ | $0.1987_{0.1988}^{0.0000}$ |
| 0559 | 17 | $0.033_{0.033}^{0.001}$ | $0.114_{4.309}^{0.002}$ | $4_{25}^{3}$ | $0.0000_{0.3477}^{0.0000}$ |
| 0562 | 1 | $0.007_{0.017}^{0.007}$ | $0.054_{0.054}^{0.017}$ | $4{ }_{10}^{4}$ | $0.0091_{0.1299}^{0.0000}$ |
| 0566 | 11 | $0.007{ }_{0.007}^{0.001}$ | $0.035_{0.035}^{0.023}$ | $4{ }_{25}^{4}$ | $0.0000_{0.0000}^{0.0000}$ |
| 0567 | 7 | $0.001_{0.025}^{0.001}$ | $0.421_{0.421}^{0.136}$ | $25_{25}^{6}$ | $0.1950_{0.1950}^{0.0000}$ |

continued

| VCC | TypeID | $\mathrm{Z}_{\text {gas }}$ | $\begin{gathered} \hline \text { SFR } \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\begin{gathered} \tau \\ {[\mathrm{Gyr}]} \end{gathered}$ | E(B-V) <br> [mag] |
| :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] |
| 0570 | 17 | $0.033_{0.033}^{0.032}$ | $0.004_{0.004}^{0.000}$ | $3{ }_{3}^{0}$ | $0.0463_{0.0463}^{0.0399}$ |
| 0574 | 25 | $0.007_{0.014}^{0.004}$ | $0.065_{0.065}^{0.023}$ | $4{ }_{15}^{4}$ | $0.3145_{0.3145}^{0.1654}$ |
| 0576 | 5 | $0.030_{0.030}^{0.001}$ | $0.340_{2.454}^{0.340}$ | $5{ }_{25}^{5}$ | $0.1628_{0.4393}^{0.1628}$ |
| 0583 | 3 | $0.001_{0.025}^{0.0001}$ | $0.151_{0.151}^{0.048}$ | $25_{25}^{6}$ | $0.2247_{0.2247}^{0.0261}$ |
| 0584 | 3 | $0.025_{0.030}^{0.001}$ | $0.020_{0.108}^{0.011}$ | $6_{25}^{4}$ | $0.0000_{0.2331}^{0.0000}$ |
| 0585 | 3 | $0.001_{0.009}^{0.001}$ | $0.014_{0.014}^{0.011}$ | $25_{25}^{10}$ | $0.0668_{0.0668}^{0.0013}$ |
| 0593 | 23 | $0.001_{0.014}^{0.001}$ | $0.162_{0.162}^{0.101}$ | $25^{8}$ | $0.1630_{0.1630}^{0.0549}$ |
| 0596 | 7 | $0.001_{0.025}^{0.001}$ | $30.433_{30.437}^{9.969}$ | $25^{6}$ | $0.2376_{0.2377}^{0.0392}$ |
| 0618 | 3 | $0.001_{0.004}^{10.001}$ | $0.019_{0.019}^{0.019}$ | $25_{25}^{15}$ | $0.0059_{0.0059}^{0.0000}$ |
| 0620 | 11 | $0.001_{0.004}^{0.001}$ | $0.053_{0.053}^{0.050}$ | $25_{25}^{15}$ | $0.0795_{0.0795}^{0.0574}$ |
| 0630 | 9 | $0.030_{0.030}^{0.030}$ | $0.241_{0.241}^{0.241}$ | 55 | $0.0000_{0.0000}^{0.0000}$ |
| 0633 | 25 | $0.030_{0.033}^{0.001}$ | $0.005_{0.060}^{0.000}$ | $5_{25}^{2}$ | $0.0000_{0.2812}^{0.0000}$ |
| 0641 | 1 | $0.001_{0.030}^{0.001}$ | $0.041_{0.073}^{0.013}$ | $25_{25}^{4}$ | $0.0873_{0.1282}^{0.0000}$ |
| 0651 | 25 | $0.001_{0.033}^{10.001}$ | $0.032_{0.058}^{0.000}$ | $25_{25}^{2}$ | $0.1744_{0.2135}^{0.0000}$ |
| 0654 | 19 | $0.033_{0.033}^{0.032}$ | $0.006_{0.297}^{0.000}$ | $3{ }_{4}^{0}$ | $0.1012_{0.1324}^{0.0948}$ |
| 0655 | 1 | $0.001_{0.030}^{0.001}$ | $1.438_{1.438}^{0.213}$ | $25^{5} 5$ | $0.2526_{0.2527}^{0.0000}$ |
| 0656 | 5 | $0.033_{0.033}^{0.032}$ | $0.104_{0.104}^{0.000}$ | $4_{4}^{0}$ | $0.1082_{0.1084}^{0.0700}$ |
| 0657 | 19 | $0.033_{0.033}^{0.032}$ | $0.002_{0.105}^{0.000}$ | $3{ }_{4}^{0}$ | $0.1255_{0.1566}^{0.1191}$ |
| 0664 | 7 | $0.001_{0.007}^{0.001}$ | $0.331_{0.583}^{0.331}$ | $25_{25}^{4}$ | $0.0000_{0.0405}^{0.0000}$ |
| 0666 | 3 | $0.001_{0.014}^{0.001}$ | $0.039_{0.039}^{0.024}$ | $25_{25}^{8}$ | $0.1674_{0.1674}^{0.0563}$ |

continued

| VCC | Type- <br> ID | $Z_{\text {gas }}$ | SFR $\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]$ | $\begin{gathered} \tau \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{E}(\mathrm{~B}-\mathrm{V}) \\ {[\mathrm{mag}]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] |
| 0667 | 7 | $0.030_{0.030}^{0.001}$ | $0.101_{0.726}^{0.100}$ | $5_{25}^{5}$ | $0.0000_{0.2766}^{0.0000}$ |
| 0672 | 19 | $0.033_{0.033}^{0.001}$ | $0.025_{0.908}^{0.000}$ | $4_{25}^{2}$ | $0.0000_{0.3326}^{0.0000}$ |
| 0679 | 11 | $0.001_{0.025}^{0.001}$ | $0.080_{0.080}^{0.028}$ | $25_{25}^{6}$ | $0.1628_{0.1628}^{0.000}$ |
| 0688 | 7 | $0.030_{0.030}^{0.001}$ | $0.104_{0.737}^{0.104}$ | $5_{25}^{5}$ | $0.0000_{0.2693}^{0.0000}$ |
| 0692 | 7 | $0.001_{0.014}^{0.001}$ | $1.128_{1.128}^{0.706}$ | $25_{25}^{8}$ | $0.13511_{0.1352}^{0.0272}$ |
| 0697 | 7 | $0.001_{0.025}^{0.001}$ | $0.702_{0.702}^{0.223}$ | $25_{25}^{6}$ | $0.23466_{0.2346}^{0.0360}$ |
| 0699 | 7 | $0.001_{0.009}^{0.001}$ | $0.452_{0.452}^{0.360}$ | $25_{25}^{10}$ | $0.0648_{0.0648}^{0.0006}$ |
| 0703 | 3 | $0.004_{0.009}^{0.001}$ | $0.014_{0.027}^{0.012}$ | $15_{25}^{4}$ | $0.1129_{0.1735}^{0.0666}$ |
| 0705 | 25 | $0.004_{0.033}^{0.001}$ | $0.0511_{0.094}^{0.000}$ | $15_{25}^{3}$ | $0.26211_{0.3204}^{0.0000}$ |
| 0713 | 7 | $0.033_{0.033}^{0.032}$ | $0.066_{0.066}^{0.000}$ | $4_{4}^{0.5}$ | $0.0470_{0.0471}^{0.0083}$ |
| 0739 | 9 | $0.001_{0.009}^{0.001}$ | $0.2633_{0.263}^{0.230}$ | $25_{25}^{10}$ | $0.03488_{0.0349}^{0.0000}$ |
| 0740 | 11 | $0.001_{0.009}^{0.001}$ | $0.0477_{0.047}^{0.038}$ | $25_{25}^{10}$ | $0.0548_{0.0548}^{0.0000}$ |
| 0768 | 7 | $0.001_{0.014}^{0.001}$ | $0.134_{0.134}^{0.084}$ | $25_{25}^{8}$ | $0.11855_{0.1185}^{0.0108}$ |
| 0780 | 25 | $0.009_{0.014}^{0.001}$ | $0.020_{0.045}^{0.016}$ | $10_{25}^{4}$ | $0.11988_{0.2276}^{0.0760}$ |
| 0792 | 17 | $0.033_{0.033}^{0.032}$ | $0.003_{0.003}^{0.000}$ | $3_{3}^{0}$ | $0.0017_{0.0017}^{0.0000}$ |
| 0793 | 3 | $1.000_{1.000}^{1.000}$ | $0.000_{0.000}^{0.000}$ | $0_{0}^{0}$ | $0.0000_{0.0000}^{0.0000}$ |
| 0801 | 15 | $0.001_{0.014}^{0.001}$ | $1.509_{1.509}^{0.945}$ | $25_{25}^{8}$ | $0.1239_{0.1240}^{0.0165}$ |
| 0802 | 1 | $0.001_{0.017}^{0.001}$ | $0.025_{0.045}^{0.014}$ | $25_{25}^{4}$ | $0.0769_{0.1592}^{0.0000}$ |
| 0809 | 7 | $0.025_{0.025}^{0.001}$ | $0.109_{0.341}^{0.109}$ | $6_{25}^{6}$ | $0.0233_{0.2220}^{0.023}$ |
| 0825 | 3 | $0.001_{0.025}^{0.001}$ | $0.072_{0.072}^{0.024}$ | $25_{25}^{6}$ | $0.1799_{0.1800}^{0.0000}$ |

continued

| VCC | TypeID | $\mathrm{Z}_{\text {gas }}$ | $\begin{gathered} \mathrm{SFR} \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\begin{gathered} \tau \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{E}(\mathrm{~B}-\mathrm{V}) \\ {[\mathrm{mag}]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] |
| 0826 | 3 | $0.004_{0.009}^{0.004}$ | $0.090_{0.170}^{0.077}$ | $15_{15}^{4}$ | $0.0381_{0.1021}^{0.0000}$ |
| 0827 | 7 | $0.025_{0.025}^{0.025}$ | $0.383_{0.383}^{0.382}$ | 66 | $0.0709_{0.0709}^{0.0708}$ |
| 0836 | 17 | $0.030_{0.030}^{0.001}$ | $1.066_{7.734}^{1.066}$ | $5_{25}^{5}$ | $0.03711_{0.3149}^{0.0371}$ |
| 0841 | 1 | $0.001_{0.025}^{0.001}$ | $0.162_{0.162}^{0.052}$ | $25_{25}^{6}$ | $0.2210_{0.2210}^{0.0222}$ |
| 0849 | 5 | $0.001_{0.014}^{0.001}$ | $0.747_{0.747}^{0.478}$ | $25_{25}^{8}$ | $0.1011_{0.1011}^{0.0000}$ |
| 0851 | 7 | $0.030_{0.030}^{0.001}$ | $0.105_{0.733}^{0.105}$ | $5{ }_{25}^{5}$ | $0.0000_{0.2641}^{0.0000}$ |
| 0857 | 5 | $0.033_{0.033}^{0.001}$ | $0.294_{11.186}^{0.006}$ | $4_{25}^{3}$ | $0.0000_{0.3517}^{0.0000}$ |
| 0859 | 7 | $0.033_{0.033}^{0.033}$ | $0.032_{0.174}^{0.000}$ | 45 | $0.0956_{0.1766}^{0.0572}$ |
| 0865 | 7 | $0.007_{0.014}^{0.004}$ | $1.801_{1.801}^{0.716}$ | $4{ }_{15}^{4}$ | $0.1120_{0.1120}^{0.0000}$ |
| 0869 | 25 | $0.030_{0.030}^{0.001}$ | $0.025_{0.178}^{0.025}$ | $5{ }_{25}^{5}$ | $0.0000_{0.2649}^{0.0000}$ |
| 0873 | 7 | $0.033_{0.033}^{0.032}$ | $0.153_{0.154}^{0.000}$ | $4{ }_{4}^{0}$ | $0.0373_{0.0374}^{0.0000}$ |
| 0874 | 7 | $0.030_{0.030}^{0.001}$ | $0.339_{2.465}^{0.339}$ | $5_{25}^{5}$ | $0.0246_{0.3027}^{0.0246}$ |
| 0888 | 3 | $0.014_{0.014}^{0.001}$ | $0.053_{0.083}^{0.053}$ | $8{ }_{25}^{8}$ | $0.0000_{0.1015}^{0.0000}$ |
| 0890 | 1 | $0.001_{0.009}^{0.001}$ | $0.027_{0.027}^{0.022}$ | $25_{25}^{10}$ | $0.0644_{0.0644}^{0.0000}$ |
| 0905 | 7 | $0.001_{0.014}^{0.001}$ | $0.681_{0.681}^{0.426}$ | $25_{25}^{8}$ | $0.1244_{0.1244}^{0.0168}$ |
| 0912 | 5 | $0.030_{0.030}^{0.001}$ | $0.344_{2.446}^{0.344}$ | $5{ }_{25}^{5}$ | $0.0000_{0.2712}^{0.0000}$ |
| 0938 | 7 | $0.014_{0.014}^{0.004}$ | $0.788_{2.217}^{0.788}$ | $8{ }_{15}^{4}$ | $0.0476_{0.1969}^{0.0476}$ |
| 0939 | 7 | $0.001_{0.014}^{0.001}$ | $1.212_{1.212}^{0.758}$ | $25_{25}^{8}$ | $0.14988_{0.1498}^{0.0418}$ |
| 0945 | 11 | $0.001_{0.007}^{0.001}$ | $0.068{ }_{0.114}^{0.068}$ | $25_{25}^{4}$ | $0.0000_{0.0266}^{0.0000}$ |
| 0950 | 11 | $0.007_{0.009}^{0.004}$ | $0.153_{0.153}^{0.078}$ | $4_{15}^{4}$ | $0.0670_{0.0670}^{0.0000}$ |

continued

| VCC | TypeID | $\mathrm{Z}_{\text {gas }}$ | $\begin{gathered} \hline \mathrm{SFR} \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\begin{gathered} \tau \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \mathrm{E}(\mathrm{~B}-\mathrm{V}) \\ {[\mathrm{mag}]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] |
| 0952 | 3 | $0.007{ }_{0.007}^{0.001}$ | $0.023_{0.023}^{0.015}$ | $4{ }_{25}^{4}$ | $0.0000_{0.0000}^{0.0000}$ |
| 0957 | 7 | $0.001_{0.014}^{0.001}$ | $1.357_{1.357}^{0.850}$ | $25_{25}^{8}$ | $0.1242_{0.1243}^{0.0166}$ |
| 0958 | 17 | $0.033_{0.033}^{0.032}$ | $0.006_{0.006}^{0.000}$ | $3{ }_{3}^{0}$ | $0.0665_{0.0665}^{0.0601}$ |
| 0963 | 3 | $0.007{ }_{0.014}^{0.004}$ | $0.045_{0.045}^{0.016}$ | $4_{15}^{4}$ | $0.1825_{0.1825}^{0.0316}$ |
| 0975 | 7 | $1.000_{1.000}^{1.000}$ | $0.000_{0.000}^{0.000}$ | $0_{0}^{0}$ | $0.0000_{0.0000}^{0.0000}$ |
| 0979 | 17 | $0.001_{0.030}^{0.001}$ | $4.8633_{4.863}^{0.670}$ | $25_{25}^{5}$ | $0.3029_{0.3029}^{0.0249}$ |
| 0980 | 7 | $0.001_{0.007}^{0.001}$ | $0.196_{0.341}^{0.196}$ | $25_{25}^{4}$ | $0.0000_{0.0362}^{0.0000}$ |
| 0984 | 17 | $0.033_{0.033}^{0.032}$ | $0.002_{0.002}^{0.000}$ | $33_{3}^{0}$ | $0.0301_{0.0302}^{0.0237}$ |
| 0989 | 7 | $0.001_{0.009}^{0.001}$ | $0.027_{0.027}^{0.022}$ | $25_{25}^{10}$ | $0.0635_{0}^{0.00635}$ |
| 0994 | 25 | $0.033_{0.033}^{0.032}$ | $0.001_{0.001}^{0.000}$ | $4_{4}^{0}$ | $0.1360_{0.1361}^{0.1027}$ |
| 0995 | 7 | $0.001_{0.004}^{0.001}$ | $0.087_{0.087}^{0.082}$ | $25_{25}^{15}$ | $0.0285_{0.0285}^{0.0069}$ |
| 1001 | 3 | $0.014_{0.014}^{0.004}$ | $0.014_{0.040}^{0.014}$ | $8{ }_{15}^{4}$ | $0.0000_{0.1483}^{0.0000}$ |
| 1011 | 9 | $0.001_{0.014}^{0.001}$ | $0.188_{0.188}^{0.117}$ | $25_{25}^{8}$ | $0.1719_{0.1719}^{0.0634}$ |
| 1013 | 3 | $0.001_{0.014}^{0.001}$ | $0.036_{0.064}^{0.022}$ | $25_{25}^{4}$ | $0.12399_{0.1661}^{00.0149}$ |
| 1017 | 3 | $0.030_{0.033}^{0.001}$ | $0.030_{0.327}^{0.000}$ | $5_{25}^{2}$ | $0.0000_{0.2598}^{0.0000}$ |
| 1020 | 25 | $0.030_{0.030}^{0.001}$ | $0.006_{0.043}^{0.006}$ | $5_{25}^{5}$ | $0.0078_{0.2789}^{0.0078}$ |
| 1021 | 3 | $0.014_{0.014}^{0.001}$ | $0.090_{0.144}^{0.090}$ | $8_{25}^{8}$ | $0.0641_{0.1724}^{0.0641}$ |
| 1047 | 17 | $0.033_{0.033}^{0.032}$ | $0.003_{0.003}^{0.000}$ | $3{ }_{3}^{0}$ | $0.0590_{0.0590}^{0.0527}$ |
| 1060 | 11 | $0.007{ }_{0.007}^{0.001}$ | $0.0677_{0.067}^{0.041}$ | $4{ }_{25}^{4}$ | $0.0122_{0.0122}^{00.0000}$ |
| 1086 | 23 | $0.033_{0.033}^{0.033}$ | $0.048_{0.049}^{0.000}$ | $4_{4}^{0}$ | $0.0313_{0.0314}^{0.0000}$ |

continued

| VCC | TypeID | $\mathrm{Z}_{\text {gas }}$ | $\begin{gathered} \mathrm{SFR} \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\begin{gathered} \tau \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{E}(\mathrm{~B}-\mathrm{V}) \\ {[\mathrm{mag}]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] |
| 1091 | 5 | $0.001_{0.004}^{0.001}$ | $0.195_{0.196}^{0.195}$ | $25_{25}^{15}$ | $0.0000_{0.0000}^{0.0000}$ |
| 1098 | 3 | $0.001_{0.025}^{0.001}$ | $0.023^{0.041}$ | $25_{25}^{4}$ | $0.1527_{0.1991}^{0.0000}$ |
| 1102 | 3 | $0.001_{0.014}^{0.001}$ | $0.013^{0.013}$ | $25_{25}^{8}$ | $0.0980_{0.0981}^{0.0000}$ |
| 1106 | 3 | $0.025_{0.025}^{0.001}$ | $0.008_{0.026}^{0.008}$ | $6{ }_{25}^{6}$ | $0.0314_{0.2313}^{0.0314}$ |
| 1110 | 17 | $0.033_{0.033}^{0.032}$ | $0.000_{0.013}^{0.000}$ | $2{ }_{3}^{0}$ | $0.0000_{0.0000}^{0.0000}$ |
| 1114 | 3 | $0.001_{0.001}^{0.001}$ | $0.254_{0.254}^{0.254}$ | $25_{25}^{25}$ | $0.1705_{0.1706}^{0.1705}$ |
| 1118 | 7 | $0.001_{0.030}^{0.001}$ | $1.373_{1.374}^{0.206}$ | $25_{25}^{5}$ | $0.2485_{0.2486}^{0.0000}$ |
| 1121 | 3 | $0.033_{0.033}^{0.001}$ | $0.001_{0.053}^{0.000}$ | $4_{25}^{0}$ | $0.0000_{0.3204}^{0.0000}$ |
| 1126 | 7 | $0.033_{0.033}^{0.030}$ | $0.078_{0.420}^{0.002}$ | 45 | $0.0000_{0.0805}^{0.0000}$ |
| 1128 | 3 | $0.004_{0.009}^{0.001}$ | $0.034_{0.065}^{0.029}$ | $15_{25}^{4}$ | $0.1327_{0.1958}^{0.0861}$ |
| 1141 | 1 | $0.001_{0.025}^{0.001}$ | $0.063_{0.063}^{0.020}$ | $25_{25}^{6}$ | $0.19688_{0.2764}^{0.0000}$ |
| 1145 | 5 | $0.033_{0.033}^{0.032}$ | $0.0077_{0.007}^{0.000}$ | $3{ }_{3}^{0}$ | $0.0154_{0.0154}^{0.0089}$ |
| 1156 | 7 | $0.001_{0.004}^{0.001}$ | $0.161_{0.161}^{0.152}$ | $25_{25}^{15}$ | $0.0318_{0.0319}^{0.0110}$ |
| 1158 | 17 | $0.033_{0.033}^{0.032}$ | $0.000_{0.291}^{0.000}$ | $2_{4}^{0}$ | $0.0967{ }_{0.1309}^{0.0932}$ |
| 1165 | 25 | $0.007_{0.014}^{0.001}$ | $0.024_{0.024}^{0.008}$ | $4{ }_{25}^{4}$ | $0.23955_{0.2395}^{0.0864}$ |
| 1166 | 3 | $0.001_{0.033}^{0.001}$ | $0.032_{0.056}^{0.001}$ | $25_{25}^{4}$ | $0.1892_{0.2277}^{0.0000}$ |
| 1168 | 3 | $0.001_{0.033}^{0.001}$ | $0.021_{0.038}^{0.000}$ | $25_{25}^{3}$ | $0.0477_{0.0989}^{0.0000}$ |
| 1169 | 3 | $0.007{ }_{0.007}^{0.001}$ | $0.010_{0.010}^{0.006}$ | $4{ }_{25}^{4}$ | $0.0233_{0.0233}^{0.0000}$ |
| 1179 | 29 | $0.001_{0.004}^{0.001}$ | $0.114_{0.114}^{0.107}$ | $25_{25}^{15}$ | $0.1412_{0.1412}^{0.1193}$ |
| 1186 | 25 | $0.007_{0.014}^{0.004}$ | $0.027_{0.027}^{0.009}$ | $4_{15}^{4}$ | $0.2186_{0.2186}^{0.0639}$ |

continued

| VCC | TypeID | $\mathrm{Z}_{\text {gas }}$ | $\begin{gathered} \mathrm{SFR} \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\begin{gathered} \tau \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{E}(\mathrm{~B}-\mathrm{V}) \\ {[\mathrm{mag}]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] |
| 1188 | 15 | $0.030_{0.030}^{0.001}$ | $0.153_{1.112}^{0.153}$ | $5_{25}^{5}$ | $0.0236_{0.3014}^{0.0236}$ |
| 1189 | 7 | $0.009_{0.014}^{0.004}$ | $0.3833_{0.453}^{0.302}$ | $10_{15}^{8}$ | $0.0745_{0.1170}^{0.0313}$ |
| 1190 | 17 | $0.033_{0.033}^{0.032}$ | $0.005_{0.005}^{0.000}$ | $3{ }_{3}^{0}$ | $0.0801_{0.0801}^{0.0737}$ |
| 1193 | 7 | $0.001_{0.001}^{0.001}$ | $0.392_{0.392}^{0.392}$ | $25_{25}^{25}$ | $0.1748_{0.1749}^{0.1748}$ |
| 1200 | 3 | $0.001_{0.004}^{0.001}$ | $0.074_{0.075}^{0.074}$ | $25_{25}^{15}$ | $0.0000_{0.0000}^{0.0000}$ |
| 1205 | 7 | $0.001_{0.014}^{0.001}$ | $1.103_{1.103}^{0.691}$ | $25_{25}^{8}$ | $0.1302_{0.1303}^{0.0227}$ |
| 1208 | 11 | $0.001_{0.009}^{0.001}$ | $0.059_{0.059}^{0.048}$ | $25_{25}^{10}$ | $0.05599_{0.0559}^{0.0000}$ |
| 1217 | 11 | $0.001_{0.009}^{0.001}$ | $0.187_{0.187}^{0.149}$ | $25_{25}^{10}$ | $0.0910_{0.0910}^{0.0261}$ |
| 1227 | 25 | $0.032_{0.033}^{0.025}$ | $0.000_{0.022}^{0.000}$ | $0.5{ }_{6}^{0}$ | $0.1815_{0.3403}^{0.1815}$ |
| 1237 | 29 | $0.001_{0.017}^{0.001}$ | $0.075_{0.133}^{0.041}$ | $25_{25}^{4}$ | $0.1516_{0.2326}^{0.0432}$ |
| 1257 | 3 | $0.007_{0.007}^{0.001}$ | $0.050_{0.050}^{0.027}$ | $4{ }_{25}^{4}$ | $0.0811_{0.0811}^{0.0171}$ |
| 1266 | 9 | $0.007{ }_{0.007}^{0.007}$ | $0.104_{0.104}^{0.104}$ | $4{ }_{4}^{4}$ | $0.0174_{0.0175}^{0.0174}$ |
| 1273 | 3 | $0.030_{0.030}^{0.001}$ | $0.0311_{0.207}^{0.031}$ | $5{ }_{25}$ | $0.0000_{0.2493}^{0.0000}$ |
| 1287 | 3 | $0.007_{0.025}^{0.001}$ | $0.000_{0.000}^{0.000}$ | $4{ }_{25}^{4}$ | $0.0000_{0.0000}^{0.0000}$ |
| 1290 | 5 | $0.001_{0.025}^{0.001}$ | $1.348_{1.349}^{0.438}$ | $25_{25}^{6}$ | $0.1919_{0.1919}^{0.0000}$ |
| 1294 | 19 | $0.030_{0.030}^{0.001}$ | $0.025_{0.176}^{0.025}$ | $5{ }_{25}$ | $0.0000_{0.2696}^{0.0000}$ |
| 1313 | 1 | $0.009_{0.031}^{0.007}$ | $0.009_{0.024}^{0.004}$ | $10_{10}^{4}$ | $0.0000_{0.0000}^{0.0000}$ |
| 1326 | 17 | $0.030_{0.033}^{0.001}$ | $0.274_{1.987}^{0.056}$ | $5{ }_{25}^{4}$ | $0.0506_{0.3280}^{0.0000}$ |
| 1330 | 17 | $0.032_{0.033}^{0.033}$ | $0.000_{0.003}^{0.000}$ | $0.5{ }_{3}^{0}$ | $0.0054_{0.0116}^{0.0053}$ |
| 1331 | 25 | $0.007_{0.014}^{0.001}$ | $0.053_{0.053}^{0.019}$ | $4{ }_{25}^{4}$ | $0.2558_{0.2558}^{0.1056}$ |

continued

| VCC | $\begin{gathered} \hline \hline \text { Type- } \\ \text { ID } \end{gathered}$ | $\mathrm{Z}_{\text {gas }}$ | $\begin{gathered} \hline \hline \text { SFR } \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\begin{gathered} \tau \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{E}(\mathrm{~B}-\mathrm{V}) \\ {[\mathrm{mag}]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] |
| 1336 | 25 | $0.001_{0.033}^{0.001}$ | $0.013_{0.023}^{0.000}$ | $25_{25}^{2}$ | $0.0778_{0.1094}^{0.0000}$ |
| 1356 | 11 | $0.001_{0.004}^{0.001}$ | $0.073_{0.073}^{0.071}$ | $25_{25}^{15}$ | $0.0111_{0.0112}^{0.0000}$ |
| 1358 | 30 | $0.001_{0.025}^{0.001}$ | $0.173_{0.173}^{0.055}$ | $25_{25}^{6}$ | $0.2089_{0.2090}^{0.0099}$ |
| 1361 | 25 | $0.001_{0.033}^{10.001}$ | $0.036_{0.036}^{0.000}$ | $25_{25}^{0}$ | $0.3526_{0.3526}^{0.0000}$ |
| 1374 | 1 | $0.001_{0.001}^{0.001}$ | $0.131_{0.131}^{0.131}$ | $25_{25}^{25}$ | $0.0006_{0.0006}^{0.0006}$ |
| 1375 | 7 | $0.001_{0.001}^{10.001}$ | $2.108_{2.108}^{2.108}$ | $25_{25}^{25}$ | $0.0889_{0.0889}^{0.0089}$ |
| 1377 | 3 | $0.007_{0.007}^{0.001}$ | $0.145_{0.146}^{0.078}$ | $4{ }_{25}^{4}$ | $0.0601_{0.0601}^{0.0000}$ |
| 1379 | 7 | $0.001_{0.014}^{0.001}$ | $1.567_{1.567}^{0.981}$ | $25_{25}^{8}$ | $0.1498_{0.1498}^{0.0421}$ |
| 1393 | 7 | $0.001_{0.014}^{10.001}$ | $0.763_{0.763}^{0.477}$ | $25_{25}^{8}$ | $0.1553_{0.1553}^{0.0474}$ |
| 1397 | 3 | $0.001_{0.025}^{10.001}$ | $0.008_{0.014}^{0.003}$ | $25_{25}^{4}$ | $0.1341_{0.1820}^{0.0000}$ |
| 1403 | 3 | $0.007_{0.014}^{70.004}$ | $0.071_{0.071}^{0.025}$ | $4{ }_{15}^{4}$ | $0.2758_{0.2758}^{0.1288}$ |
| 1408 | 25 | $0.007_{0.033}^{70.001}$ | $0.042_{0.042}^{0.000}$ | $4{ }_{25}^{3}$ | $0.3153_{0.3154}^{0.0000}$ |
| 1410 | 11 | $0.001_{0.014}^{0.001}$ | $0.249_{0.249}^{0.156}$ | $25_{25}^{8}$ | $0.1064_{0.1064}^{0.0000}$ |
| 1411 | 1 | $0.001_{0.014}^{10.001}$ | $0.101_{0.101}^{0.063}$ | $25_{25}^{8}$ | $0.1459_{0.1792}^{0.0370}$ |
| 1412 | 17 | $0.033_{0.033}^{30.032}$ | $0.007_{0.007}^{0.000}$ | $3{ }_{3}^{0}$ | $0.0769_{0.0769}^{0.0706}$ |
| 1413 | 25 | $0.007_{0.033}^{70.001}$ | $0.046_{0.046}^{0.001}$ | $4{ }_{25}^{4}$ | $0.3971_{0.3971}^{0.00199}$ |
| 1419 | 23 | $0.030_{0.033}^{0.001}$ | $0.211_{1.531}^{0.043}$ | $5{ }_{25}^{4}$ | $0.0512_{0.3285}^{0.0000}$ |
| 1426 | 3 | $0.007_{0.030}^{0.001}$ | $0.363_{0.363}^{0.030}$ | $4{ }_{25}^{4}$ | $0.3598_{0.3598}^{0.0577}$ |
| 1427 | 29 | $0.001_{0.025}^{10.001}$ | $0.134_{0.134}^{0.046}$ | $25_{25}^{6}$ | $0.1718_{0.1718}^{0.0000}$ |
| 1435 | 3 | $0.001_{0.025}^{0.001}$ | $0.470_{0.470}^{0.150}$ | $25_{25}^{6}$ | $0.2196_{0.2197}^{0.0210}$ |

continued

| VCC | TypeID | $\mathrm{Z}_{\text {gas }}$ | $\begin{gathered} \mathrm{SFR} \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\tau$ [Gyr] | $\mathrm{E}(\mathrm{B}-\mathrm{V})$ [mag] |
| :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] |
| 1437 | 1 | $0.001_{0.025}^{0.001}$ | $0.181_{0.181}^{0.060}$ | $25^{6}$ | $0.1824_{0.2152}^{0.0000}$ |
| 1442 | 9 | $0.001_{0.014}^{0.001}$ | $0.248_{0.248}^{0.155}$ | $25^{8}$ | $0.1251_{0.1251}^{0.0174}$ |
| 1448 | 25 | $0.001_{0.014}^{0.001}$ | $0.523_{0.523}^{0.328}$ | $25_{25}^{8}$ | $0.1540_{0.1540}^{0.0464}$ |
| 1450 | 7 | $0.001_{0.009}^{0.001}$ | $0.730_{0.730}^{0.581}$ | $25_{25}^{10}$ | $0.0650_{0.0650}^{0.0009}$ |
| 1455 | 3 | $0.001_{0.014}^{0.001}$ | $0.043_{0.043}^{0.027}$ | $25_{25}^{8}$ | $0.1636_{0.1636}^{0.0538}$ |
| 1459 | 1 | $0.001_{0.025}^{0.001}$ | $0.053_{0.094}^{0.017}$ | $25_{25}^{4}$ | $0.2226_{0.3020}^{0.0235}$ |
| 1465 | 3 | $0.007_{0.033}^{0.001}$ | $0.115_{0.115}^{0.000}$ | $4_{25}^{3}$ | $0.0000_{0.0000}^{0.0000}$ |
| 1468 | 3 | $0.007_{0.007}^{0.0001}$ | $0.059_{0.059}^{0.037}$ | $4{ }_{25}^{4}$ | $0.0088_{0.0089}^{0.0000}$ |
| 1507 | 11 | $0.001_{0.025}^{0.001}$ | $0.120_{0.120}^{0.041}$ | $25_{25}^{6}$ | $0.1790_{0.1790}^{0.0000}$ |
| 1508 | 7 | $0.001_{0.001}^{0.001}$ | $1.473_{1.473}^{1.473}$ | $25_{25}^{25}$ | $0.0697_{0.0697}^{0.0696}$ |
| 1515 | 3 | $0.001_{0.033}^{0.001}$ | $0.008_{0.012}^{0.000}$ | $25^{3}$ | $0.0000_{0.0000}^{0.0000}$ |
| 1516 | 7 | $0.001_{0.025}^{0.001}$ | $1.707_{1.707}^{0.543}$ | $25^{6}$ | $0.2188_{0.2188}^{0.0201}$ |
| 1524 | 9 | $0.001_{0.007}^{0.001}$ | $0.574_{1.019}^{0.541}$ | $25_{25}^{4}$ | $0.0283_{0.0718}^{0.0068}$ |
| 1529 | 9 | $0.001_{0.014}^{0.001}$ | $0.248_{0.248}^{0.156}$ | $25_{25}^{8}$ | $0.1498_{0.1499}^{0.0418}$ |
| 1532 | 7 | $0.001_{0.014}^{0.001}$ | $0.705_{0.705}^{0.441}$ | $25_{25}^{8}$ | $0.1572_{0.1573}^{0.0496}$ |
| 1540 | 5 | $0.030_{0.033}^{0.001}$ | $2.922_{21.110}^{0.011}$ | $5_{25}^{3}$ | $0.1770_{0.4533}^{0.0632}$ |
| 1552 | 17 | $0.033_{0.033}^{0.032}$ | $0.003_{0.003}^{0.000}$ | $3{ }_{3}^{0}$ | $0.0000_{0.0000}^{0.0000}$ |
| 1554 | 11 | $0.007_{0.007}^{0.001}$ | $2.341_{2.341}^{1.359}$ | $4{ }_{25}^{4}$ | $0.0339_{0.0340}^{0.0000}$ |
| 1555 | 7 | $0.001_{0.025}^{0.001}$ | $14.984_{14.986}^{4.770}$ | $25_{25}^{6}$ | $0.2150_{0.2150}^{0.0165}$ |
| 1557 | 7 | $0.001_{0.025}^{0.001}$ | $0.497_{0.497}^{0.158}$ | $25_{25}^{6}$ | $0.2121_{0.2121}^{0.0135}$ |

continued

| VCC | TypeID | $\mathrm{Z}_{\text {gas }}$ | $\begin{gathered} \mathrm{SFR} \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\begin{gathered} \tau \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{E}(\mathrm{~B}-\mathrm{V}) \\ {[\mathrm{mag}]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] |
| 1562 | 7 | $0.001_{0.025}^{0.001}$ | $11.318_{11.318}^{3.602}$ | $25_{25}^{6}$ | $0.2543_{0.2543}^{0.0557}$ |
| 1566 | 9 | $0.007_{0.014}^{0.004}$ | $0.290_{0.290}^{0.103}$ | $4{ }_{15}^{4}$ | $0.2068_{0.2068}^{0.0585}$ |
| 1569 | 7 | $0.001_{0.014}^{0.001}$ | $0.093{ }_{0.093}^{0.058}$ | $25_{25}^{8}$ | $0.1264_{0.1264}^{0.0170}$ |
| 1572 | 1 | $0.010_{0.030}^{0.009}$ | $0.042_{0.096}^{0.022}$ | $10_{10}^{6}$ | $0.2752_{0.3366}^{0.0300}$ |
| 1574 | 23 | $0.014_{0.025}^{0.001}$ | $0.020_{0.032}^{0.010}$ | $8{ }_{25}^{6}$ | $0.1246_{0.2347}^{0.0356}$ |
| 1575 | 11 | $0.030_{0.030}^{0.030}$ | $0.149_{0.149}^{0.149}$ | 55 | $0.0000_{0.0000}^{0.0000}$ |
| 1581 | 11 | $0.001_{0.014}^{0.001}$ | $0.252_{0.252}^{0.158}$ | $25_{25}^{8}$ | $0.1456_{0.1456}^{0.0375}$ |
| 1582 | 25 | $0.001_{0.014}^{0.001}$ | $0.034_{0.034}^{0.021}$ | $25_{25}^{8}$ | $0.1619_{0.1619}^{0.0525}$ |
| 1585 | 3 | $0.004_{0.004}^{0.001}$ | $0.080_{0.084}^{0.080}$ | $15_{25}^{15}$ | $0.0000_{0.0175}^{0.0000}$ |
| 1588 | 7 | $0.030_{0.030}^{0.001}$ | $0.458{ }_{3.183}^{0.457}$ | $5{ }_{25}$ | $0.0000_{0.2634}^{0.0000}$ |
| 1596 | 3 | $0.001_{0.009}^{0.001}$ | $0.008_{0.015}^{0.008}$ | $25_{25}^{4}$ | $0.0245_{0.0687}^{0.0000}$ |
| 1605 | 9 | $0.001_{0.014}^{0.001}$ | $0.016_{0.016}^{0.011}$ | $25_{25}^{8}$ | $0.0743_{0.0744}^{0.0000}$ |
| 1615 | 5 | $0.033_{0.033}^{0.032}$ | $0.013^{0.013}$ | $3{ }_{3}^{0}$ | $0.0000_{0.0000}^{0.0000}$ |
| 1623 | 29 | $0.001_{0.030}^{0.001}$ | $0.034_{0.060}^{0.010}$ | $25_{25}^{4}$ | $0.1347_{0.2166}^{0.0000}$ |
| 1624 | 7 | $0.030_{0.033}^{0.001}$ | $0.187_{1.355}^{0.035}$ | $5_{25}^{4}$ | $0.0768_{0.3543}^{0.0000}$ |
| 1644 | 11 | $0.001_{0.009}^{0.001}$ | $0.013_{0.024}^{0.012}$ | $25_{25}^{4}$ | $0.02911_{0.0765}^{0.0000}$ |
| 1645 | 25 | $0.014_{0.025}^{0.001}$ | $0.006{ }_{0.016}^{0.003}$ | $8{ }_{25}^{4}$ | $0.0552_{0.2092}^{0.0000}$ |
| 1654 | 3 | $0.001_{0.009}^{0.001}$ | $0.046_{0.046}^{0.036}$ | $25_{25}^{10}$ | $0.0696_{0.0696}^{0.0055}$ |
| 1656 | 25 | $0.025_{0.030}^{0.001}$ | $0.015^{0.048}$ | $6_{25}^{5}$ | $0.0087_{0.2084}^{0.0000}$ |
| 1675 | 15 | $0.025_{0.025}^{0.001}$ | $0.114_{0.358}^{0.114}$ | $6_{25}^{6}$ | $0.0038_{0.2023}^{0.0038}$ |

continued

| VCC | TypeID | $\mathrm{Z}_{\text {gas }}$ | $\begin{gathered} \mathrm{SFR} \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\begin{gathered} \tau \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \mathrm{E}(\mathrm{~B}-\mathrm{V}) \\ {[\mathrm{mag}]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] |
| 1678 | 9 | $0.001_{0.004}^{0.001}$ | $0.224_{0.224}^{0.212}$ | $25_{25}^{15}$ | $0.0190_{0.0190}^{0.0000}$ |
| 1685 | 9 | $0.001_{0.007}^{0.001}$ | $0.038_{0.067}^{0.038}$ | $25_{25}^{4}$ | $0.0000_{0.0399}^{0.0000}$ |
| 1686 | 11 | $0.001_{0.009}^{0.001}$ | $0.574_{0.574}^{0.510}$ | $25_{25}^{10}$ | $0.0287_{0.0287}^{0.0000}$ |
| 1690 | 17 | $0.030_{0.030}^{0.001}$ | $4.048{ }_{29.374}^{4.048}$ | $5{ }_{25}^{5}$ | $0.0223_{0.3002}^{0.0223}$ |
| 1696 | 7 | $0.030_{0.030}^{0.001}$ | 0.9356 .9391 | $5_{25}^{5}$ | $0.0048_{0.2832}^{0.0048}$ |
| 1699 | 11 | $0.001_{0.004}^{0.001}$ | $0.199_{0.199}^{0.188}$ | $25_{25}^{15}$ | $0.0275_{0.0275}^{0.0058}$ |
| 1713 | 29 | $0.001_{0.010}^{0.001}$ | $0.049_{0.049}^{0.033}$ | $25_{25}^{10}$ | $0.0632_{0.1500}^{0.0000}$ |
| 1725 | 11 | $0.001_{0.001}^{0.001}$ | $0.225_{0.225}^{0.225}$ | $25_{25}^{25}$ | $0.0730_{0.0731}^{0.0730}$ |
| 1726 | 9 | $0.0077_{0.007}^{0.007}$ | $0.133_{0.133}^{0.133}$ | $4_{4}^{4}$ | $0.0161_{0.0161}^{0.0160}$ |
| 1727 | 17 | $0.033_{0.033}^{0.033}$ | $0.021_{1.058}^{0.000}$ | $3{ }_{4}^{0}$ | $0.0000_{0.0298}^{0.0000}$ |
| 1728 | 3 | $0.001_{0.025}^{0.001}$ | $0.040_{0.040}^{0.014}$ | $25_{25}^{6}$ | $0.15633_{0.1564}^{0.0000}$ |
| 1730 | 7 | $0.033_{0.033}^{0.032}$ | $0.002_{0.003}^{0.000}$ | $33_{3}^{0}$ | $0.0000_{0.0000}^{0.0000}$ |
| 1744 | 1 | $0.010_{0.030}^{0.001}$ | $0.010_{0.025}^{0.005}$ | $10_{25}^{4}$ | $0.0000_{0.0981}^{0.0000}$ |
| 1756 | 25 | $0.0077_{0.014}^{0.004}$ | $0.024_{0.024}^{0.008}$ | $4_{15}^{4}$ | $0.2444_{0.2445}^{0.0921}$ |
| 1757 | 17 | $0.030_{0.033}^{0.001}$ | $0.153_{1.108}^{0.034}$ | $5_{25}^{4}$ | $0.0168_{0.2944}^{0.0000}$ |
| 1758 | 7 | $0.001_{0.025}^{0.001}$ | $0.239_{0.239}^{0.079}$ | $25_{25}^{6}$ | $0.1847_{0.1848}^{0.0000}$ |
| 1760 | 17 | $0.033_{0.033}^{0.032}$ | $0.003_{0.003}^{0.000}$ | $33_{3}^{0}$ | $0.0518_{0.0518}^{0.0452}$ |
| 1776 | 25 | $0.001_{0.033}^{0.001}$ | $0.034_{0.060}^{0.000}$ | $25_{25}^{2}$ | $0.3112_{0.3464}^{0.0000}$ |
| 1778 | 15 | $0.001_{0.025}^{0.001}$ | $0.406_{0.406}^{0.141}$ | $25_{25}^{6}$ | $0.1681_{0.1681}^{0.0000}$ |
| 1784 | 3 | $0.033_{0.033}^{0.001}$ | $0.002_{0.070}^{0.000}$ | $4_{25}^{0}$ | $0.0000_{0.3248}^{0.0000}$ |

[^30]| VCC | TypeID | $\mathrm{Z}_{\text {gas }}$ | $\begin{gathered} \mathrm{SFR} \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\tau$ [Gyr] | $\begin{gathered} \mathrm{E}(\mathrm{~B}-\mathrm{V}) \\ {[\mathrm{mag}]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] |
| 1789 | 3 | $0.001_{0.025}^{0.001}$ | $0.160_{0.160}^{0.053}$ | $25_{25}^{6}$ | $0.1831_{0.1832}^{0.0000}$ |
| 1791 | 11 | $0.007_{0.007}^{0.007}$ | $0.250_{0.250}^{0.250}$ | $4{ }_{4}^{4}$ | $0.0000_{0.0000}^{0.0000}$ |
| 1804 | 29 | $0.001_{0.025}^{0.001}$ | $0.116_{0.116}^{0.037}$ | $25_{25}^{6}$ | $0.2347_{0.2347}^{0.0358}$ |
| 1811 | 7 | $0.001_{0.014}^{0.001}$ | $1.170_{1.170}^{0.732}$ | $25_{25}^{8}$ | $0.1200_{0.1200}^{0.0122}$ |
| 1813 | 17 | $0.033_{0.033}^{0.032}$ | $0.012_{0.012}^{0.000}$ | $3{ }_{3}^{0}$ | $0.0619_{0.0619}^{0.0554}$ |
| 1816 | 3 | $0.001_{0.014}^{0.001}$ | $0.040_{0.074}^{0.035}$ | $25_{25}^{4}$ | $0.0118_{0.0603}^{0.0000}$ |
| 1821 | 11 | $0.009_{0.014}^{0.001}$ | $0.015_{0.019}^{0.012}$ | $10_{25}^{8}$ | $0.1206_{0.1888}^{0.0774}$ |
| 1822 | 3 | $0.001_{0.014}^{0.001}$ | $0.036{ }_{0.036}^{0.023}$ | $25_{25}^{8}$ | $0.0922_{0.0922}^{0.0000}$ |
| 1825 | 25 | $0.030_{0.033}^{0.001}$ | $0.014_{0.100}^{0.003}$ | $5_{25}^{4}$ | $0.0000_{0.2659}^{0.0000}$ |
| 1834 | 19 | $0.033^{0.033} 0$ | $0.060_{0.323}^{0.000}$ | 45 | $0.0000_{0.0825}^{0.0000}$ |
| 1837 | 25 | $0.025_{0.025}^{0.001}$ | $0.015_{0.046}^{0.015}$ | $6_{25}^{6}$ | $0.0206_{0.2205}^{0.0206}$ |
| 1855 | 19 | $0.025_{0.025}^{0.001}$ | $0.048_{0.152}^{0.048}$ | $6{ }_{25}^{6}$ | $0.0283_{0.2271}^{0.0283}$ |
| 1859 | 17 | $0.033_{0.033}^{0.001}$ | $0.098{ }_{3.528}^{0.000}$ | 42 | $0.0000_{0.3322}^{0.0000}$ |
| 1860 | 25 | $0.001_{0.033}^{0.001}$ | $0.023_{0.043}^{0.000}$ | $25^{2}$ | $0.0907_{0.1508}^{0.0000}$ |
| 1868 | 7 | $0.032_{0.033}^{0.032}$ | $0.000_{0.049}^{0.000}$ | $0.5{ }_{4}^{0}$ | $0.0028_{0.0416}^{0.0027}$ |
| 1873 | 11 | $0.004_{0.007}^{0.001}$ | $0.021_{0.039}^{0.021}$ | $15_{25}^{4}$ | $0.0373_{0.1007}^{0.0373}$ |
| 1883 | 19 | $0.033_{0.033}^{0.032}$ | $0.003{ }_{0.003}^{0.000}$ | $3{ }_{3}^{0}$ | $0.0000_{0.0000}^{0.0000}$ |
| 1884 | 25 | $0.001_{0.033}^{0.001}$ | $0.032_{0.050}^{0.000}$ | $25^{2}$ | $0.0000_{0.0000}^{0.0000}$ |
| 1885 | 3 | $0.014_{0.014}^{0.004}$ | $0.046_{0.131}^{0.046}$ | $8{ }_{15}^{4}$ | $0.1338_{0.2822}^{0.1338}$ |
| 1898 | 25 | $0.030_{0.030}^{0.001}$ | $0.020_{0.147}^{0.020}$ | $5_{25}^{5}$ | $0.0205_{0.2965}^{0.0205}$ |

continued

| VCC | TypeID | $\mathrm{Z}_{\text {gas }}$ | $\begin{gathered} \mathrm{SFR} \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\begin{gathered} \tau \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\mathrm{E}(\mathrm{B}-\mathrm{V})$ [mag] |
| :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] |
| 1900 | 25 | $0.001_{0.030}^{0.001}$ | $0.072_{0.125}^{0.012}$ | $25_{25}^{4}$ | $0.2199_{0.2555}^{0.0000}$ |
| 1902 | 19 | $0.033_{0.033}^{0.030}$ | $0.051_{0.262}^{0.000}$ | 45 | $0.0000_{0.0657}^{0.0000}$ |
| 1905 | 25 | $0.007{ }_{0.033}^{0.001}$ | $0.071_{0.071}^{0.000}$ | $4_{25}^{0.5}$ | $0.4220_{0.4220}^{0.0176}$ |
| 1906 | 19 | $0.033_{0.033}^{0.001}$ | $0.004_{0.158}^{0.000}$ | $4_{25}^{0}$ | $0.0000_{0.3273}^{0.0000}$ |
| 1918 | 3 | $0.001_{0.001}^{0.001}$ | $0.021_{0.021}^{0.021}$ | $25_{25}^{25}$ | $0.0049_{0.0050}^{0.0049}$ |
| 1920 | 19 | $0.030_{0.030}^{0.001}$ | $0.025_{0.175}^{0.025}$ | $5{ }_{25}^{5}$ | $0.0000_{0.2700}^{0.0000}$ |
| 1923 | 11 | $0.025_{0.025}^{0.001}$ | $0.522_{1.639}^{0.522}$ | $6{ }_{25}^{6}$ | $0.0068_{0.2052}^{0.0068}$ |
| 1929 | 7 | $0.001_{0.014}^{0.001}$ | $0.671_{0.671}^{0.420}$ | $25_{25}^{8}$ | $0.1585_{0.1585}^{0.0508}$ |
| 1931 | 3 | $0.001_{0.007}^{0.001}$ | $0.056_{0.099}^{0.054}$ | $25_{25}^{4}$ | $0.0151_{0.0572}^{0.0000}$ |
| 1932 | 7 | $0.030_{0.030}^{0.001}$ | $0.279_{2.028}^{0.279}$ | $5{ }_{25}^{5}$ | $0.0077_{0.2860}^{0.0077}$ |
| 1933 | 17 | $0.001_{0.009}^{0.001}$ | $0.027_{0.048}^{0.023}$ | $25_{25}^{4}$ | $0.0476_{0.0919}^{0.0000}$ |
| 1943 | 5 | $0.030_{0.030}^{0.001}$ | $0.5855_{4.194}^{0.584}$ | $5{ }_{25}^{5}$ | $0.0000_{0.2741}^{0.0000}$ |
| 1944 | 31 | $0.033_{0.033}^{0.030}$ | $0.002_{0.009}^{0.000}$ | $4_{5}^{0}$ | $0.0468_{0.1220}^{0.0111}$ |
| 1952 | 3 | $0.007{ }_{0.007}^{0.007}$ | $0.051_{0.051}^{0.051}$ | $4_{4}^{4}$ | $0.0556_{0.0557}^{0.0556}$ |
| 1955 | 30 | $0.001_{0.030}^{0.001}$ | $0.925_{0.925}^{0.146}$ | $25_{25}^{5}$ | $0.2300_{0.2300}^{0.0000}$ |
| 1960 | 29 | $0.001_{0.014}^{0.001}$ | $0.009_{0.009}^{0.006}$ | $25_{25}^{8}$ | $0.1074{ }_{0.1074}^{0.0000}$ |
| 1965 | 3 | $0.001_{0.014}^{0.001}$ | $0.043_{0.043}^{0.027}$ | $25_{25}^{8}$ | $0.1644_{0.1645}^{0.0535}$ |
| 1970 | 3 | $1.000_{1.000}^{1.000}$ | $0.000_{0.000}^{0.000}$ | $0_{0}^{0}$ | $0.0000_{0.0000}^{0.0000}$ |
| 1987 | 7 | $0.001_{0.025}^{0.001}$ | $9.113_{9.115}^{2.901}$ | $25^{6}{ }^{6}$ | $0.2109_{0.2110}^{0.0124}$ |
| 1992 | 3 | $0.001_{0.009}^{0.001}$ | $0.044_{0.072}^{0.043}$ | $25_{25}^{4}$ | $0.0000_{0.0206}^{0.0000}$ |

continued

|  | VCC | $\begin{gathered} \hline \hline \text { Type- } \\ \text { ID } \end{gathered}$ | $\mathrm{Z}_{\text {gas }}$ | $\begin{gathered} \mathrm{SFR} \\ {\left[\mathrm{M}_{\odot} / \mathrm{yr}\right]} \end{gathered}$ | $\begin{gathered} \tau \\ {[\mathrm{Gyr}]} \end{gathered}$ | $\begin{gathered} \hline \text { E(B-V) } \\ {[\mathrm{mag}]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | [1] | [2] | [3] | [4] | [5] | [6] |
|  | 1994 | 25 | $0.033_{0.033}^{0.001}$ | $0.000_{0.072}^{0.000}$ | $3{ }_{25}^{0}$ | $0.0990_{0.4786}^{0.0921}$ |
|  | 1999 | 17 | $0.033_{0.033}^{0.032}$ | $0.000_{0.002}^{0.000}$ | $2{ }_{3}^{0}$ | $0.0182_{0.0212}^{0.0147}$ |
|  | 2006 | 9 | $0.001_{0.025}^{0.001}$ | $0.678_{0.678}^{0.215}$ | $25_{25}^{6}$ | $0.2240_{0.2240}^{0.0250}$ |
|  | 2007 | 29 | $0.001_{0.014}^{0.001}$ | $0.077_{0.077}^{0.048}$ | $25_{25}^{8}$ | $0.1532_{0.1533}^{0.0449}$ |
|  | 2015 | 1 | $0.001_{0.014}^{0.001}$ | $0.050_{0.050}^{0.032}$ | $25_{25}^{8}$ | $0.1258_{0.1258}^{0.0173}$ |
|  | 2023 | 7 | $0.001_{0.014}^{0.001}$ | $0.357_{0.358}^{0.232}$ | $25_{25}^{8}$ | $0.0965_{0.0965}^{0.0000}$ |
|  | 2033 | 1 | $0.001_{0.014}^{10.001}$ | $0.133_{0.133}^{0.083}$ | $25_{25}^{8}$ | $0.1297_{0.1632}^{0.0219}$ |
|  | 2034 | 3 | $0.001_{0.014}^{10.001}$ | $0.079_{0.079}^{0.050}$ | $25_{25}^{8}$ | $0.1247_{0.1248}^{0.0161}$ |
| N | 2037 | 3 | $0.001_{0.025}^{0.001}$ | $0.093_{0.168}^{0.039}$ | $25_{25}^{4}$ | $0.1069_{0.1522}^{0.0000}$ |
|  | 2058 | 7 | $0.030_{0.030}^{0.001}$ | $1.188_{8.610}^{1.185}$ | $5_{25}^{5}$ | $0.0184_{0.2964}^{0.0182}$ |
|  | 2066 | 15 | $0.030_{0.030}^{0.001}$ | $0.597_{4.234}^{0.596}$ | $5_{25}^{5}$ | $0.0000_{0.2706}^{0.0000}$ |
|  | 2070 | 17 | $0.033_{0.033}^{0.032}$ | $0.000_{0.009}^{0.000}$ | $3{ }_{3}^{0}$ | $0.0326_{0.0326}^{0.0263}$ |
|  | 2089 | 31 | $0.033_{0.033}^{0.032}$ | $0.000_{0.000}^{0.000}$ | $2{ }_{3}^{0}$ | $0.1639_{0.1669}^{0.1604}$ |
|  | 2090 | 25 | $0.001_{0.025}^{10.001}$ | $0.100_{0.109}^{0.035}$ | $25_{25}^{6}$ | $0.1955_{0.1955}^{0.0000}$ |
|  | 2094 | 3 | $0.007_{0.014}^{0.001}$ | $0.018_{0.018}^{0.006}$ | $4{ }_{25}^{4}$ | $0.2334_{0.2334}^{0.0827}$ |

[^31]
### 6.4 Results of the GAZELLE surface brightness profiles

In Chapter 4 the time evolution of the SBPs was analysed, using the GALEV/GAZELLEmodels. The following Tab. 6.6 shows the derived parameters from this analysis for the different evolutionary assumptions: galaxy today $(\mathrm{t}=0)$; undisturbed evolution in 2 $\operatorname{Gyr}(\mathrm{t}=2)$ and truncated evolution in $2 \mathrm{Gyr}(\mathrm{t}=2, \operatorname{tru})$.

Table 6.6: Derived parameters of the fits to the SBPs for different evolutionary scenarios.

| VCC | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[\left[^{\prime \prime}\right]\right. \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2) \end{gathered}$ | $(t=2)$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[\left[^{\prime \prime}\right]\right. \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \hline \mathrm{R}_{\mathrm{eff}}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] | [9] | [10] | [11] | [12] | [13] |
| 0001 | 19.070 | 20.187 | 5.338 | 8.956 | 19.326 | 20.443 | 5.354 | 8.985 | 19.428 | 20.544 | 5.410 | 9.079 |
| 0004 | 22.016 | 23.132 | 6.154 | 10.326 | 21.709 | 22.826 | 5.956 | 9.994 | 22.948 | 24.065 | 6.503 | 10.912 |
| 0010 | 19.343 | 20.459 | 4.476 | 7.511 | 19.509 | 20.625 | 4.778 | 8.018 | 19.762 | 20.878 | 4.742 | 7.958 |
| 0015 | 21.562 | 22.678 | 10.692 | 17.941 | 21.405 | 22.522 | 10.410 | 17.468 | 21.700 | 22.816 | 10.652 | 17.874 |
| 0017 | 22.074 | 23.191 | 10.348 | 17.363 | 21.918 | 23.034 | 10.692 | 17.942 | 23.052 | 24.168 | 11.535 | 19.356 |
| 0020 | 22.238 | 23.354 | 3.975 | 6.670 | 22.178 | 23.294 | 4.316 | 7.243 | 22.869 | 23.985 | 5.123 | 8.597 |
| 0022 | 20.285 | 21.401 | 3.695 | 6.200 | 20.193 | 21.310 | 3.592 | 6.027 | 21.030 | 22.147 | 4.456 | 7.477 |
| 0024 | 19.291 | 20.407 | 3.815 | 6.401 | 19.227 | 20.344 | 3.823 | 6.415 | 19.840 | 20.957 | 4.248 | 7.129 |
| 0025 | 18.725 | 19.841 | 9.506 | 15.951 | 18.922 | 20.038 | 10.809 | 18.137 | 18.970 | 20.086 | 8.122 | 13.628 |
| 0026 | 22.251 | 23.367 | 4.333 | 7.270 | 22.040 | 23.156 | 4.325 | 7.258 | 24.223 | 25.340 | 14.201 | 23.829 |
| 0030 | 22.544 | 23.660 | 9.589 | 16.091 | 22.579 | 23.696 | 10.492 | 17.606 | 23.414 | 24.531 | 11.233 | 18.850 |
| 0031 | 19.277 | 20.393 | 3.557 | 5.968 | 19.147 | 20.263 | 3.557 | 5.968 | 19.653 | 20.769 | 3.465 | 5.814 |
| 0034 | 20.963 | 22.079 | 10.228 | 17.163 | 20.833 | 21.949 | 10.229 | 17.163 | 21.512 | 22.628 | 10.097 | 16.943 |
| 0041 | 23.172 | 24.288 | 9.636 | 16.169 | 23.110 | 24.226 | 10.082 | 16.917 | 23.462 | 24.578 | 9.667 | 16.221 |
| 0048 | 21.784 | 22.900 | 15.178 | 25.469 | 21.722 | 22.838 | 15.845 | 26.587 | 22.285 | 23.401 | 14.004 | 23.498 |
| 0052 | 22.537 | 23.653 | 5.500 | 9.229 | 22.503 | 23.619 | 5.692 | 9.551 | 23.089 | 24.206 | 5.306 | 8.904 |

continued

| VCC | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2) \end{gathered}$ | $(\mathrm{t}=2)$ | $\begin{gathered} \hline \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] | [9] | [10] | [11] | [12] | [13] |
| 0058 | 20.367 | 21.483 | 20.407 | 34.243 | 20.471 | 21.588 | 21.594 | 36.234 | 20.645 | 21.761 | 20.438 | 34.294 |
| 0066 | 19.946 | 21.063 | 40.426 | 67.834 | 19.908 | 21.024 | 42.608 | 71.496 | 20.031 | 21.147 | 35.175 | 59.023 |
| 0067 | 21.020 | 22.136 | 20.850 | 34.987 | 20.813 | 21.929 | 19.887 | 33.370 | 21.649 | 22.766 | 22.062 | 37.021 |
| 0073 | 18.014 | 19.130 | 9.325 | 15.647 | 18.247 | 19.364 | 9.269 | 15.554 | 17.811 | 18.927 | 8.545 | 14.339 |
| 0074 | 21.224 | 22.341 | 8.438 | 14.159 | 21.029 | 22.145 | 7.737 | 12.982 | 21.522 | 22.639 | 8.640 | 14.498 |
| 0081 | 21.995 | 23.112 | 13.654 | 22.912 | 22.021 | 23.137 | 14.491 | 24.316 | 22.404 | 23.521 | 12.151 | 20.389 |
| 0083 | 21.134 | 22.250 | 7.637 | 12.814 | 20.927 | 22.043 | 7.305 | 12.258 | 21.745 | 22.862 | 8.630 | 14.481 |
| 0085 | 23.671 | 24.787 | 9.304 | 15.613 | 23.500 | 24.616 | 9.186 | 15.415 | 24.145 | 25.262 | 9.189 | 15.419 |
| 0087 | 20.997 | 22.113 | 12.360 | 20.741 | 20.935 | 22.051 | 12.951 | 21.731 | 21.850 | 22.966 | 13.348 | 22.398 |
| 0089 | 19.541 | 20.657 | 17.679 | 29.665 | 19.582 | 20.698 | 17.943 | 30.108 | 19.699 | 20.815 | 17.251 | 28.947 |
| 0092 | 19.104 | 20.220 | 56.265 | 94.413 | 19.289 | 20.406 | 55.254 | 92.717 | 18.880 | 19.997 | 49.734 | 83.454 |
| 0097 | 19.198 | 20.314 | 16.073 | 26.971 | 19.491 | 20.607 | 17.618 | 29.563 | 19.367 | 20.483 | 16.485 | 27.662 |
| 0099 | 19.918 | 21.034 | 8.715 | 14.624 | 20.216 | 21.333 | 8.907 | 14.947 | 20.218 | 21.334 | 9.443 | 15.845 |
| 0104 | 24.335 | 25.452 | 13.015 | 21.840 | 24.257 | 25.373 | 12.839 | 21.544 | 24.998 | 26.114 | 13.727 | 23.034 |
| 0105 | 21.515 | 22.632 | 19.575 | 32.846 | 21.427 | 22.543 | 19.787 | 33.202 | 21.841 | 22.958 | 19.903 | 33.397 |
| 0113 | 22.513 | 23.630 | 8.345 | 14.002 | 22.119 | 23.235 | 8.005 | 13.432 | 23.680 | 24.796 | 11.296 | 18.954 |
| 0114 | 23.055 | 24.172 | 10.937 | 18.352 | 22.836 | 23.952 | 10.558 | 17.716 | 23.748 | 24.865 | 13.984 | 23.465 |
| 0117 | 21.412 | 22.528 | 5.236 | 8.786 | 21.320 | 22.437 | 5.627 | 9.441 | 22.010 | 23.126 | 4.706 | 7.896 |
| 0119 | 21.586 | 22.703 | 14.925 | 25.044 | 21.548 | 22.665 | 15.853 | 26.601 | 22.163 | 23.279 | 13.626 | 22.864 |
| 0120 | 19.585 | 20.701 | 19.749 | 33.139 | 19.682 | 20.798 | 21.047 | 35.317 | 19.924 | 21.041 | 18.671 | 31.330 |

continued

| VCC | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=0) \end{gathered}$ | $(\mathrm{t}=0)$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[\left[^{\prime \prime}\right]\right. \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2) \end{gathered}$ | $(t=2)$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \mathrm{R}_{\mathrm{efff}}\left[\left[^{\prime \prime}\right]\right. \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] | [9] | [10] | [11] | [12] | [13] |
| 0126 | 21.910 | 23.026 | 17.858 | 29.965 | 21.878 | 22.995 | 18.916 | 31.742 | 22.130 | 23.247 | 16.455 | 27.612 |
| 0130 | 20.751 | 21.867 | 3.751 | 6.294 | 20.627 | 21.743 | 3.725 | 6.251 | 21.596 | 22.712 | 3.946 | 6.621 |
| 0131 | 19.510 | 20.626 | 15.018 | 25.200 | 19.782 | 20.898 | 16.683 | 27.994 | 19.420 | 20.536 | 14.037 | 23.555 |
| 0132 | 23.015 | 24.131 | 10.599 | 17.785 | 22.976 | 24.092 | 11.305 | 18.969 | 24.116 | 25.232 | 14.052 | 23.579 |
| 0137 | 21.443 | 22.560 | 5.356 | 8.988 | 21.840 | 22.956 | 6.012 | 10.088 | 21.850 | 22.966 | 5.684 | 9.539 |
| 0143 | 20.436 | 21.552 | 8.376 | 14.055 | 20.199 | 21.315 | 7.519 | 12.616 | 20.603 | 21.720 | 8.140 | 13.658 |
| 0144 | 18.233 | 19.350 | 2.421 | 4.063 | 19.513 | 20.630 | 3.214 | 5.392 | 20.132 | 21.248 | 3.173 | 5.324 |
| 0145 | 19.844 | 20.960 | 36.347 | 60.990 | 20.088 | 21.205 | 40.001 | 67.121 | 20.067 | 21.183 | 34.726 | 58.271 |
| 0152 | 18.959 | 20.075 | 11.577 | 19.427 | 19.162 | 20.278 | 11.438 | 19.192 | 18.624 | 19.741 | 9.988 | 16.760 |
| 0157 | 18.845 | 19.961 | 20.671 | 34.685 | 18.842 | 19.959 | 19.794 | 33.214 | 18.865 | 19.981 | 19.678 | 33.019 |
| 0159 | 21.470 | 22.586 | 8.131 | 13.644 | 21.262 | 22.378 | 8.378 | 14.058 | 22.369 | 23.485 | 8.416 | 14.123 |
| 0162 | 20.677 | 21.793 | 21.024 | 35.279 | 20.420 | 21.537 | 19.332 | 32.438 | 21.227 | 22.343 | 20.748 | 34.815 |
| 0168 | 22.477 | 23.594 | 8.295 | 13.918 | 22.421 | 23.537 | 8.645 | 14.507 | 23.227 | 24.344 | 12.905 | 21.655 |
| 0169 | 24.464 | 25.581 | 19.009 | 31.898 | 24.130 | 25.246 | 18.328 | 30.754 | 25.709 | 26.825 | 29.418 | 49.363 |
| 0172 | 20.401 | 21.518 | 9.072 | 15.223 | 20.266 | 21.382 | 8.983 | 15.073 | 21.313 | 22.430 | 10.155 | 17.040 |
| 0187 | 20.088 | 21.204 | 24.266 | 40.718 | 19.904 | 21.020 | 23.571 | 39.552 | 19.882 | 20.998 | 22.872 | 38.379 |
| 0199 | 18.886 | 20.003 | 18.644 | 31.284 | 19.065 | 20.181 | 18.864 | 31.654 | 18.463 | 19.580 | 16.592 | 27.842 |
| 0207 | 17.147 | 18.263 | 0.799 | 1.340 | 18.431 | 19.547 | 0.894 | 1.500 | 19.205 | 20.321 | 0.926 | 1.553 |
| 0213 | 19.001 | 20.118 | 5.072 | 8.510 | 18.967 | 20.084 | 4.886 | 8.198 | 19.240 | 20.356 | 5.047 | 8.470 |
| 0217 | 23.509 | 24.625 | 18.516 | 31.070 | 23.504 | 24.621 | 19.978 | 33.524 | 24.136 | 25.252 | 17.936 | 30.097 |

[^32]| VCC | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=0) \end{gathered}$ | $(\mathrm{t}=0)$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[\left[^{\prime \prime}\right]\right. \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2) \end{gathered}$ | $(t=2)$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \mathrm{R}_{\mathrm{eff}}\left[\left[^{\prime \prime}\right]\right. \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] | [9] | [10] | [11] | [12] | [13] |
| 0221 | 19.838 | 20.954 | 11.920 | 20.002 | 19.682 | 20.798 | 11.195 | 18.785 | 20.439 | 21.555 | 13.405 | 22.494 |
| 0222 | 18.712 | 19.829 | 22.792 | 38.245 | 18.889 | 20.005 | 22.980 | 38.560 | 18.346 | 19.463 | 20.590 | 34.550 |
| 0223 | 20.014 | 21.130 | 3.233 | 5.424 | 20.237 | 21.354 | 3.463 | 5.811 | 20.821 | 21.937 | 3.885 | 6.519 |
| 0226 | 18.796 | 19.912 | 15.406 | 25.851 | 19.069 | 20.185 | 15.628 | 26.224 | 18.901 | 20.017 | 14.974 | 25.126 |
| 0234 | 19.615 | 20.731 | 21.877 | 36.709 | 19.762 | 20.878 | 21.569 | 36.194 | 19.450 | 20.566 | 20.055 | 33.653 |
| 0237 | 22.223 | 23.340 | 6.265 | 10.513 | 21.781 | 22.898 | 5.925 | 9.942 | 23.587 | 24.703 | 8.957 | 15.029 |
| 0241 | 21.138 | 22.255 | 17.279 | 28.994 | 20.747 | 21.863 | 17.278 | 28.993 | 22.233 | 23.350 | 17.715 | 29.726 |
| 0247 | 23.273 | 24.390 | 6.948 | 11.659 | 23.235 | 24.352 | 7.418 | 12.448 | 24.403 | 25.520 | 11.799 | 19.799 |
| 0260 | 22.595 | 23.712 | 8.407 | 14.108 | 22.504 | 23.620 | 7.982 | 13.393 | 22.907 | 24.023 | 7.831 | 13.141 |
| 0267 | 21.067 | 22.183 | 18.162 | 30.475 | 20.986 | 22.102 | 18.515 | 31.068 | 21.581 | 22.697 | 19.004 | 31.888 |
| 0274 | 19.273 | 20.389 | 2.447 | 4.106 | 19.453 | 20.570 | 2.499 | 4.193 | 20.223 | 21.339 | 2.585 | 4.337 |
| 0286 | 20.855 | 21.972 | 4.813 | 8.077 | 20.672 | 21.789 | 4.558 | 7.648 | 21.311 | 22.428 | 4.822 | 8.091 |
| 0289 | 19.776 | 20.892 | 10.380 | 17.418 | 19.646 | 20.762 | 10.380 | 17.418 | 20.187 | 21.303 | 9.887 | 16.590 |
| 0307 | 18.944 | 20.060 | 36.430 | 61.130 | 18.814 | 19.930 | 36.431 | 61.131 | 18.959 | 20.075 | 32.428 | 54.413 |
| 0309 | 22.446 | 23.562 | 9.310 | 15.622 | 22.308 | 23.425 | 9.237 | 15.499 | 23.307 | 24.423 | 9.727 | 16.322 |
| 0318 | 20.919 | 22.035 | 14.016 | 23.518 | 20.859 | 21.976 | 14.935 | 25.061 | 21.662 | 22.779 | 14.133 | 23.715 |
| 0320 | 22.866 | 23.982 | 9.709 | 16.291 | 22.628 | 23.745 | 9.323 | 15.644 | 24.390 | 25.507 | 24.994 | 41.941 |
| 0322 | 22.068 | 23.184 | 23.926 | 40.148 | 22.029 | 23.145 | 25.937 | 43.522 | 22.607 | 23.723 | 26.328 | 44.178 |
| 0323 | 18.808 | 19.924 | 8.260 | 13.861 | 19.068 | 20.185 | 8.667 | 14.544 | 18.840 | 19.957 | 8.350 | 14.012 |
| 0324 | 16.611 | 17.727 | 2.640 | 4.430 | 16.777 | 17.893 | 2.680 | 4.497 | 17.773 | 18.889 | 2.855 | 4.790 |

continued

| VCC [1] | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=0) \\ {[2]} \\ \hline \end{gathered}$ | $\begin{gathered} \mu_{\text {eff }} \\ (\mathrm{t}=0) \\ {[3]} \\ \hline \end{gathered}$ | $\begin{gathered} \hline \alpha \\ (\mathrm{t}=0) \\ {[4]} \\ \hline \end{gathered}$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2) \\ {[6]} \\ \hline \end{gathered}$ | $\begin{gathered} \hline \mu_{\mathrm{eff}} \\ (\mathrm{t}=2) \\ {[7]} \\ \hline \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2) \\ {[8]} \\ \hline \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2) \\ {[9]} \\ \hline \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[10]} \\ \hline \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[11]} \\ \hline \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[12]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[13]} \\ \hline \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0329 | 22.624 | 23.740 | 4.542 | 7.621 | 22.564 | 23.681 | 4.887 | 8.201 | 23.575 | 24.691 | 4.803 | 8.060 |
| 0334 | 20.026 | 21.142 | 3.977 | 6.674 | 20.447 | 21.563 | 4.348 | 7.295 | 21.204 | 22.320 | 5.223 | 8.764 |
| 0340 | 19.335 | 20.451 | 5.743 | 9.637 | 18.912 | 20.028 | 4.830 | 8.104 | 20.119 | 21.235 | 6.642 | 11.145 |
| 0341 | 18.960 | 20.077 | 20.924 | 35.110 | 19.137 | 20.253 | 21.087 | 35.385 | 18.734 | 19.850 | 19.776 | 33.184 |
| 0343 | 21.182 | 22.299 | 8.270 | 13.877 | 21.052 | 22.169 | 8.270 | 13.877 | 21.357 | 22.473 | 7.769 | 13.037 |
| 0350 | 23.556 | 24.673 | 13.703 | 22.994 | 23.336 | 24.452 | 13.032 | 21.868 | 24.421 | 25.538 | 15.135 | 25.396 |
| 0358 | 18.636 | 19.752 | 6.752 | 11.330 | 18.809 | 19.926 | 6.787 | 11.388 | 18.649 | 19.765 | 6.792 | 11.396 |
| 0364 | 23.252 | 24.368 | 12.273 | 20.594 | 23.244 | 24.360 | 13.743 | 23.061 | 23.381 | 24.497 | 12.740 | 21.378 |
| 0367 | 24.494 | 25.610 | 34.828 | 58.442 | 23.992 | 25.108 | 31.910 | 53.546 | 24.576 | 25.692 | 45.209 | 75.860 |
| 0381 | 24.156 | 25.273 | 16.287 | 27.330 | 23.608 | 24.724 | 14.088 | 23.640 | 25.622 | 26.739 | 33.604 | 56.388 |
| 0382 | 18.804 | 19.921 | 14.093 | 23.648 | 18.849 | 19.965 | 14.770 | 24.785 | 19.236 | 20.353 | 14.150 | 23.744 |
| 0386 | 19.900 | 21.016 | 10.148 | 17.028 | 20.040 | 21.156 | 10.024 | 16.820 | 19.880 | 20.997 | 9.878 | 16.576 |
| 0393 | 20.160 | 21.276 | 13.459 | 22.584 | 19.917 | 21.033 | 12.490 | 20.957 | 20.310 | 21.426 | 13.063 | 21.920 |
| 0404 | 19.943 | 21.059 | 15.112 | 25.357 | 19.945 | 21.061 | 15.112 | 25.358 | 19.370 | 20.486 | 12.815 | 21.503 |
| 0410 | 20.810 | 21.926 | 2.762 | 4.634 | 21.740 | 22.856 | 3.390 | 5.689 | 22.372 | 23.488 | 3.240 | 5.437 |
| 0413 | 21.986 | 23.102 | 6.074 | 10.192 | 21.909 | 23.025 | 6.051 | 10.153 | 22.689 | 23.806 | 7.450 | 12.501 |
| 0414 | 22.277 | 23.393 | 6.381 | 10.708 | 22.423 | 23.540 | 6.656 | 11.169 | 22.658 | 23.775 | 7.464 | 12.525 |
| 0415 | 21.273 | 22.390 | 11.655 | 19.557 | 21.090 | 22.206 | 11.403 | 19.134 | 21.622 | 22.738 | 12.301 | 20.640 |
| 0423 | 25.051 | 26.167 | 38.178 | 64.062 | 24.990 | 26.107 | 38.905 | 65.283 | 25.088 | 26.204 | 38.828 | 65.154 |
| 0425 | 24.034 | 25.150 | 11.999 | 20.134 | 24.048 | 25.165 | 12.671 | 21.262 | 24.551 | 25.667 | 13.553 | 22.742 |


| VCC | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=0) \end{gathered}$ | $(\mathrm{t}=0)$ | $\begin{gathered} \hline \hline \mathrm{R}_{\text {eff }}\left[\left[^{\prime \prime}\right]\right. \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2) \end{gathered}$ | $(t=2)$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[\left[^{\prime \prime}\right]\right. \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] | [9] | [10] | [11] | [12] | [13] |
| 0428 | 21.501 | 22.617 | 5.401 | 9.063 | 22.621 | 23.737 | 6.139 | 10.301 | 23.256 | 24.372 | 6.167 | 10.349 |
| 0429 | 22.057 | 23.174 | 9.040 | 15.168 | 21.742 | 22.858 | 8.322 | 13.964 | 23.027 | 24.144 | 13.914 | 23.348 |
| 0446 | 21.546 | 22.663 | 9.855 | 16.537 | 21.707 | 22.823 | 10.926 | 18.334 | 22.087 | 23.203 | 10.539 | 17.684 |
| 0448 | 21.298 | 22.414 | 4.082 | 6.849 | 21.114 | 22.231 | 3.984 | 6.684 | 22.304 | 23.421 | 4.943 | 8.295 |
| 0449 | 19.654 | 20.770 | 12.220 | 20.505 | 19.846 | 20.963 | 13.860 | 23.257 | 18.796 | 19.912 | 10.151 | 17.033 |
| 0453 | 20.664 | 21.781 | 6.358 | 10.668 | 20.534 | 21.651 | 6.357 | 10.668 | 20.921 | 22.038 | 6.563 | 11.013 |
| 0459 | 19.453 | 20.569 | 5.226 | 8.769 | 19.335 | 20.452 | 4.907 | 8.235 | 20.377 | 21.493 | 5.681 | 9.532 |
| 0460 | 19.376 | 20.492 | 43.058 | 72.252 | 19.682 | 20.798 | 46.606 | 78.206 | 19.445 | 20.561 | 43.865 | 73.606 |
| 0465 | 19.465 | 20.582 | 21.674 | 36.370 | 19.458 | 20.574 | 22.359 | 37.519 | 20.428 | 21.544 | 25.241 | 42.355 |
| 0472 | 22.865 | 23.981 | 11.038 | 18.521 | 23.255 | 24.371 | 12.899 | 21.644 | 23.808 | 24.924 | 21.181 | 35.541 |
| 0476 | 23.101 | 24.218 | 6.756 | 11.336 | 23.079 | 24.195 | 6.656 | 11.169 | 23.582 | 24.699 | 6.662 | 11.179 |
| 0477 | 23.678 | 24.794 | 12.911 | 21.664 | 23.495 | 24.611 | 13.009 | 21.829 | 23.939 | 25.055 | 11.617 | 19.493 |
| 0479 | 22.473 | 23.589 | 12.943 | 21.718 | 22.319 | 23.435 | 12.569 | 21.090 | 23.521 | 24.637 | 19.999 | 33.559 |
| 0488 | 22.542 | 23.659 | 9.686 | 16.253 | 22.319 | 23.436 | 8.678 | 14.562 | 23.119 | 24.235 | 11.037 | 18.521 |
| 0491 | 19.988 | 21.104 | 14.174 | 23.784 | 19.526 | 20.642 | 13.174 | 22.106 | 21.378 | 22.494 | 19.530 | 32.771 |
| 0509 | 21.597 | 22.714 | 13.639 | 22.887 | 21.566 | 22.683 | 14.387 | 24.141 | 22.064 | 23.180 | 13.096 | 21.975 |
| 0512 | 22.131 | 23.247 | 15.909 | 26.695 | 22.001 | 23.117 | 15.908 | 26.694 | 22.786 | 23.902 | 17.375 | 29.156 |
| 0513 | 19.142 | 20.258 | 3.152 | 5.288 | 19.306 | 20.423 | 3.202 | 5.372 | 19.346 | 20.462 | 3.076 | 5.162 |
| 0514 | 21.803 | 22.919 | 14.270 | 23.946 | 21.980 | 23.096 | 15.806 | 26.523 | 22.011 | 23.127 | 12.880 | 21.613 |
| 0520 | 23.062 | 24.178 | 6.961 | 11.681 | 22.708 | 23.824 | 6.460 | 10.840 | 24.946 | 26.062 | 21.108 | 35.419 |

[^33]| VCC [1] | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=0) \end{gathered}$ <br> [2] | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=0) \\ {[3]} \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=0) \\ {[4]} \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=0) \\ {[5]} \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2) \\ {[6]} \end{gathered}$ | $\begin{gathered} \mu_{\text {eff }} \\ (\mathrm{t}=2) \end{gathered}$ <br> [7] | $\begin{gathered} \alpha \\ (\mathrm{t}=2) \\ {[8]} \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2) \\ {[9]} \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[10]} \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[11]} \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[12]} \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[13]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0522 | 19.960 | 21.077 | 19.888 | 33.373 | 20.153 | 21.269 | 19.704 | 33.064 | 20.138 | 21.255 | 19.430 | 32.603 |
| 0524 | 18.977 | 20.093 | 25.603 | 42.962 | 19.118 | 20.235 | 25.393 | 42.610 | 18.520 | 19.637 | 22.250 | 37.336 |
| 0530 | 23.445 | 24.561 | 23.313 | 39.119 | 23.263 | 24.379 | 23.316 | 39.125 | 24.639 | 25.755 | 41.716 | 70.000 |
| 0534 | 19.885 | 21.002 | 17.186 | 28.838 | 20.011 | 21.128 | 16.775 | 28.148 | 19.815 | 20.932 | 15.956 | 26.774 |
| 0559 | 19.452 | 20.568 | 30.653 | 51.436 | 19.669 | 20.785 | 30.291 | 50.828 | 19.444 | 20.560 | 28.154 | 47.243 |
| 0562 | 18.170 | 19.286 | 2.020 | 3.390 | 18.670 | 19.786 | 2.111 | 3.542 | 19.645 | 20.761 | 2.273 | 3.813 |
| 0566 | 21.198 | 22.314 | 5.696 | 9.557 | 20.893 | 22.010 | 5.695 | 9.557 | 22.188 | 23.304 | 5.681 | 9.533 |
| 0567 | 20.438 | 21.554 | 16.850 | 28.274 | 20.182 | 21.298 | 15.480 | 25.975 | 20.496 | 21.612 | 16.004 | 26.855 |
| 0570 | 19.378 | 20.494 | 22.533 | 37.811 | 19.543 | 20.659 | 22.533 | 37.810 | 19.298 | 20.414 | 21.586 | 36.221 |
| 0574 | 23.481 | 24.598 | 11.322 | 18.998 | 23.793 | 24.910 | 12.608 | 21.156 | 24.255 | 25.371 | 18.924 | 31.754 |
| 0576 | 19.034 | 20.150 | 16.015 | 26.873 | 19.374 | 20.490 | 16.949 | 28.441 | 18.573 | 19.689 | 14.150 | 23.744 |
| 0583 | 22.786 | 23.902 | 20.885 | 35.045 | 22.625 | 23.741 | 20.182 | 33.865 | 22.887 | 24.003 | 21.358 | 35.839 |
| 0585 | 23.897 | 25.013 | 11.434 | 19.186 | 23.762 | 24.878 | 11.877 | 19.930 | 24.679 | 25.795 | 12.999 | 21.812 |
| 0593 | 20.282 | 21.399 | 10.681 | 17.923 | 20.049 | 21.165 | 10.030 | 16.831 | 20.617 | 21.733 | 11.087 | 18.605 |
| 0596 | 19.780 | 20.896 | 58.983 | 98.974 | 19.999 | 21.116 | 66.535 | 111.645 | 19.975 | 21.091 | 59.721 | 100.212 |
| 0618 | 22.178 | 23.294 | 8.464 | 14.203 | 21.960 | 23.077 | 8.462 | 14.200 | 23.258 | 24.374 | 9.636 | 16.168 |
| 0620 | 20.834 | 21.950 | 7.997 | 13.419 | 20.745 | 21.862 | 8.106 | 13.601 | 21.616 | 22.732 | 8.937 | 14.996 |
| 0630 | 20.186 | 21.302 | 41.675 | 69.931 | 20.342 | 21.458 | 42.562 | 71.420 | 20.173 | 21.289 | 39.516 | 66.308 |
| 0641 | 20.431 | 21.547 | 4.407 | 7.395 | 20.015 | 21.131 | 3.725 | 6.250 | 20.827 | 21.944 | 4.019 | 6.744 |
| 0654 | 19.210 | 20.327 | 21.493 | 36.065 | 19.376 | 20.492 | 21.494 | 36.066 | 19.144 | 20.260 | 22.404 | 37.594 |


| VCC | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=0) \end{gathered}$ | $(\mathrm{t}=0)$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2) \end{gathered}$ | $(\mathrm{t}=2)$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[\left[^{\prime \prime}\right]\right. \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] | [9] | [10] | [11] | [12] | [13] |
| 0656 | 18.699 | 19.815 | 10.983 | 18.430 | 18.881 | 19.997 | 10.910 | 18.307 | 18.244 | 19.360 | 9.481 | 15.909 |
| 0664 | 21.296 | 22.412 | 21.712 | 36.433 | 20.968 | 22.084 | 21.106 | 35.416 | 22.552 | 23.668 | 28.128 | 47.198 |
| 0666 | 22.931 | 24.047 | 10.903 | 18.295 | 22.893 | 24.009 | 11.655 | 19.558 | 22.783 | 23.900 | 9.432 | 15.826 |
| 0667 | 20.616 | 21.732 | 17.245 | 28.936 | 20.790 | 21.906 | 17.197 | 28.856 | 20.865 | 21.982 | 17.444 | 29.270 |
| 0679 | 21.420 | 22.536 | 10.412 | 17.471 | 21.290 | 22.406 | 10.412 | 17.471 | 21.584 | 22.701 | 9.854 | 16.534 |
| 0688 | 19.624 | 20.740 | 9.321 | 15.640 | 19.760 | 20.877 | 9.778 | 16.407 | 19.851 | 20.968 | 10.099 | 16.947 |
| 0692 | 20.354 | 21.470 | 19.605 | 32.896 | 20.224 | 21.340 | 19.605 | 32.896 | 20.923 | 22.039 | 21.866 | 36.690 |
| 0697 | 20.829 | 21.945 | 14.281 | 23.964 | 20.962 | 22.078 | 15.088 | 25.319 | 21.004 | 22.120 | 14.230 | 23.879 |
| 0699 | 19.787 | 20.903 | 9.051 | 15.187 | 19.546 | 20.663 | 8.800 | 14.767 | 20.905 | 22.021 | 12.408 | 20.820 |
| 0703 | 22.258 | 23.374 | 5.818 | 9.762 | 22.129 | 23.246 | 5.934 | 9.957 | 23.180 | 24.296 | 6.913 | 11.600 |
| 0713 | 19.940 | 21.057 | 28.930 | 48.545 | 20.106 | 21.223 | 28.400 | 47.655 | 19.706 | 20.822 | 25.777 | 43.254 |
| 0739 | 21.438 | 22.554 | 16.272 | 27.304 | 21.308 | 22.424 | 16.272 | 27.304 | 22.100 | 23.216 | 15.658 | 26.275 |
| 0740 | 21.298 | 22.415 | 7.917 | 13.284 | 21.168 | 22.285 | 7.917 | 13.284 | 21.851 | 22.967 | 7.950 | 13.340 |
| 0768 | 20.152 | 21.269 | 10.370 | 17.401 | 19.919 | 21.035 | 9.739 | 16.343 | 20.796 | 21.912 | 11.647 | 19.544 |
| 0792 | 19.824 | 20.940 | 29.298 | 49.161 | 19.977 | 21.093 | 28.932 | 48.548 | 19.976 | 21.092 | 29.706 | 49.847 |
| 0801 | 18.160 | 19.277 | 9.109 | 15.284 | 18.000 | 19.116 | 8.976 | 15.062 | 18.940 | 20.056 | 11.131 | 18.677 |
| 0802 | 22.064 | 23.180 | 10.017 | 16.809 | 22.250 | 23.366 | 10.495 | 17.610 | 22.838 | 23.955 | 11.422 | 19.166 |
| 0809 | 20.308 | 21.424 | 18.315 | 30.732 | 20.464 | 21.580 | 18.070 | 30.322 | 20.596 | 21.712 | 17.775 | 29.827 |
| 0825 | 22.049 | 23.166 | 8.515 | 14.288 | 21.984 | 23.101 | 8.324 | 13.967 | 22.358 | 23.474 | 8.412 | 14.115 |
| 0826 | 22.617 | 23.734 | 16.737 | 28.085 | 22.507 | 23.624 | 16.782 | 28.160 | 23.357 | 24.473 | 16.952 | 28.446 |

continued

| VCC [1] | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=0) \\ {[2]} \\ \hline \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=0) \end{gathered}$ <br> [3] | $\begin{gathered} \alpha \\ (\mathrm{t}=0) \\ {[4]} \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=0) \\ {[5]} \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2) \end{gathered}$ <br> [6] | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2) \end{gathered}$ <br> [7] | $\begin{gathered} \alpha \\ (\mathrm{t}=2) \\ {[8]} \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2) \\ {[9]} \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[10]} \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ <br> [11] | $\begin{gathered} \alpha \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[12]} \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[13]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0827 | 19.825 | 20.941 | 25.069 | 42.066 | 20.053 | 21.170 | 27.146 | 45.551 | 19.675 | 20.792 | 23.332 | 39.151 |
| 0836 | 19.064 | 20.180 | 33.392 | 56.033 | 19.300 | 20.417 | 33.162 | 55.646 | 19.107 | 20.223 | 30.846 | 51.760 |
| 0841 | 21.126 | 22.243 | 11.232 | 18.848 | 21.126 | 22.242 | 10.691 | 17.939 | 21.620 | 22.737 | 12.669 | 21.259 |
| 0849 | 20.037 | 21.154 | 14.062 | 23.597 | 19.907 | 21.024 | 14.062 | 23.597 | 20.512 | 21.628 | 14.467 | 24.276 |
| 0851 | 20.160 | 21.277 | 21.933 | 36.803 | 20.312 | 21.428 | 21.590 | 36.228 | 20.218 | 21.335 | 20.600 | 34.567 |
| 0857 | 19.638 | 20.755 | 27.475 | 46.102 | 19.945 | 21.061 | 29.684 | 49.810 | 19.960 | 21.077 | 30.927 | 51.895 |
| 0859 | 20.515 | 21.632 | 16.768 | 28.136 | 20.764 | 21.880 | 16.886 | 28.334 | 20.038 | 21.154 | 14.284 | 23.968 |
| 0865 | 20.710 | 21.826 | 36.982 | 62.055 | 20.847 | 21.963 | 43.114 | 72.346 | 21.327 | 22.443 | 35.893 | 60.228 |
| 0869 | 22.758 | 23.874 | 17.940 | 30.103 | 22.883 | 24.000 | 18.754 | 31.470 | 22.767 | 23.883 | 16.586 | 27.832 |
| 0874 | 19.025 | 20.141 | 12.181 | 20.440 | 19.289 | 20.406 | 12.181 | 20.439 | 19.261 | 20.377 | 12.022 | 20.172 |
| 0888 | 21.548 | 22.665 | 12.827 | 21.523 | 21.534 | 22.650 | 13.629 | 22.870 | 22.046 | 23.163 | 12.710 | 21.328 |
| 0890 | 20.367 | 21.484 | 3.041 | 5.103 | 20.887 | 22.003 | 3.327 | 5.583 | 21.617 | 22.733 | 4.167 | 6.992 |
| 0905 | 21.207 | 22.323 | 20.291 | 34.049 | 21.145 | 22.261 | 21.142 | 35.477 | 21.609 | 22.725 | 20.221 | 33.931 |
| 0912 | 19.991 | 21.108 | 20.350 | 34.147 | 19.854 | 20.970 | 18.387 | 30.854 | 20.016 | 21.132 | 19.224 | 32.258 |
| 0938 | 20.053 | 21.169 | 15.267 | 25.618 | 20.265 | 21.381 | 17.314 | 29.054 | 20.253 | 21.369 | 14.400 | 24.164 |
| 0939 | 20.921 | 22.037 | 22.282 | 37.389 | 20.755 | 21.872 | 21.584 | 36.219 | 21.304 | 22.420 | 22.580 | 37.889 |
| 0945 | 21.608 | 22.724 | 12.272 | 20.592 | 21.089 | 22.205 | 11.256 | 18.888 | 22.916 | 24.033 | 15.073 | 25.293 |
| 0950 | 23.059 | 24.176 | 21.339 | 35.807 | 22.823 | 23.940 | 22.334 | 37.476 | 24.152 | 25.268 | 24.740 | 41.514 |
| 0957 | 18.795 | 19.911 | 12.928 | 21.693 | 18.841 | 19.958 | 13.370 | 22.434 | 19.130 | 20.246 | 12.335 | 20.698 |
| 0958 | 17.764 | 18.880 | 17.210 | 28.878 | 17.947 | 19.063 | 17.209 | 28.877 | 17.425 | 18.542 | 15.291 | 25.659 |


| VCC | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \hline \mathrm{R}_{\mathrm{eff}}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] | [9] | [10] | [11] | [12] | [13] |
| 0963 | 22.623 | 23.740 | 8.523 | 14.302 | 22.610 | 23.727 | 9.437 | 15.836 | 23.595 | 24.711 | 12.644 | 21.217 |
| 0979 | 19.798 | 20.914 | 27.931 | 46.869 | 20.002 | 21.118 | 30.489 | 51.160 | 20.056 | 21.173 | 31.887 | 53.506 |
| 0980 | 21.200 | 22.317 | 19.453 | 32.643 | 21.028 | 22.144 | 19.769 | 33.172 | 22.105 | 23.221 | 19.830 | 33.275 |
| 0984 | 18.560 | 19.676 | 16.594 | 27.845 | 18.709 | 19.825 | 16.459 | 27.619 | 18.494 | 19.610 | 15.630 | 26.227 |
| 0989 | 22.022 | 23.138 | 6.149 | 10.319 | 21.700 | 22.816 | 5.743 | 9.636 | 22.961 | 24.077 | 6.949 | 11.660 |
| 0995 | 20.217 | 21.333 | 9.932 | 16.666 | 20.087 | 21.203 | 9.932 | 16.666 | 20.802 | 21.918 | 9.302 | 15.608 |
| 1001 | 22.248 | 23.364 | 6.732 | 11.296 | 22.230 | 23.346 | 6.863 | 11.516 | 23.352 | 24.468 | 10.000 | 16.779 |
| 1011 | 21.188 | 22.305 | 12.798 | 21.476 | 21.100 | 22.216 | 12.963 | 21.753 | 21.464 | 22.580 | 13.190 | 22.133 |
| 1013 | 22.334 | 23.450 | 8.408 | 14.109 | 22.178 | 23.295 | 7.873 | 13.212 | 22.998 | 24.114 | 9.401 | 15.775 |
| 1017 | 23.255 | 24.371 | 26.205 | 43.972 | 23.393 | 24.509 | 28.175 | 47.277 | 23.423 | 24.540 | 25.683 | 43.095 |
| 1020 | 22.086 | 23.202 | 7.801 | 13.090 | 22.337 | 23.453 | 7.811 | 13.107 | 22.470 | 23.586 | 8.290 | 13.910 |
| 1021 | 21.785 | 22.901 | 14.179 | 23.792 | 21.658 | 22.774 | 14.694 | 24.657 | 21.791 | 22.907 | 12.987 | 21.791 |
| 1047 | 18.703 | 19.819 | 11.225 | 18.836 | 18.860 | 19.977 | 11.213 | 18.816 | 18.614 | 19.731 | 10.980 | 18.425 |
| 1060 | 21.340 | 22.456 | 8.522 | 14.300 | 21.191 | 22.307 | 8.810 | 14.784 | 22.429 | 23.546 | 9.732 | 16.330 |
| 1086 | 19.207 | 20.324 | 18.116 | 30.399 | 19.436 | 20.553 | 18.036 | 30.265 | 19.066 | 20.182 | 16.574 | 27.812 |
| 1091 | 19.546 | 20.663 | 8.857 | 14.862 | 19.329 | 20.446 | 8.857 | 14.862 | 20.299 | 21.416 | 8.486 | 14.239 |
| 1098 | 25.427 | 26.543 | 28.277 | 47.449 | 25.134 | 26.250 | 26.498 | 44.464 | 25.833 | 26.949 | 26.913 | 45.160 |
| 1102 | 21.907 | 23.024 | 6.067 | 10.181 | 21.647 | 22.764 | 5.513 | 9.252 | 22.663 | 23.779 | 6.669 | 11.191 |
| 1106 | 22.858 | 23.974 | 8.735 | 14.657 | 23.199 | 24.315 | 9.447 | 15.852 | 23.350 | 24.466 | 9.786 | 16.421 |
| 1110 | 19.121 | 20.237 | 37.754 | 63.351 | 19.228 | 20.345 | 36.523 | 61.285 | 19.056 | 20.172 | 35.681 | 59.872 |

continued

| VCC | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{R}_{\text {eff }}\left[\left[^{\prime \prime}\right]\right. \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{R}_{\text {eff }}\left[{ }^{[\prime \prime}\right] \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{R}_{\mathrm{eff}}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] | [9] | [10] | [11] | [12] | [13] |
| 1114 | 21.996 | 23.112 | 22.779 | 38.223 | 21.785 | 22.902 | 21.284 | 35.715 | 22.189 | 23.305 | 21.764 | 36.520 |
| 1118 | 18.935 | 20.051 | 10.065 | 16.890 | 18.922 | 20.038 | 9.317 | 15.633 | 19.085 | 20.202 | 9.678 | 16.240 |
| 1121 | 23.212 | 24.328 | 10.041 | 16.849 | 23.194 | 24.310 | 10.147 | 17.027 | 23.616 | 24.732 | 12.645 | 21.218 |
| 1126 | 19.406 | 20.522 | 18.989 | 31.863 | 19.639 | 20.756 | 18.988 | 31.862 | 19.482 | 20.598 | 18.161 | 30.474 |
| 1128 | 23.040 | 24.156 | 10.340 | 17.350 | 22.979 | 24.095 | 11.095 | 18.617 | 23.945 | 25.061 | 14.846 | 24.911 |
| 1141 | 20.115 | 21.232 | 3.436 | 5.766 | 19.953 | 21.069 | 3.218 | 5.400 | 20.511 | 21.627 | 3.561 | 5.975 |
| 1145 | 17.864 | 18.980 | 13.524 | 22.693 | 18.053 | 19.170 | 13.656 | 22.915 | 17.945 | 19.061 | 13.338 | 22.381 |
| 1156 | 21.545 | 22.662 | 17.991 | 30.189 | 21.483 | 22.600 | 18.597 | 31.206 | 22.125 | 23.241 | 17.566 | 29.476 |
| 1158 | 18.046 | 19.163 | 12.528 | 21.023 | 18.199 | 19.316 | 12.592 | 21.130 | 17.826 | 18.943 | 12.341 | 20.707 |
| 1165 | 23.599 | 24.715 | 9.552 | 16.029 | 23.878 | 24.994 | 11.129 | 18.674 | 24.203 | 25.320 | 10.743 | 18.028 |
| 1166 | 23.189 | 24.305 | 10.915 | 18.315 | 23.251 | 24.367 | 11.224 | 18.833 | 23.773 | 24.889 | 13.623 | 22.859 |
| 1168 | 21.856 | 22.972 | 5.472 | 9.182 | 21.816 | 22.932 | 5.902 | 9.903 | 22.968 | 24.085 | 11.105 | 18.635 |
| 1169 | 21.974 | 23.090 | 3.924 | 6.584 | 21.736 | 22.853 | 3.843 | 6.448 | 23.745 | 24.862 | 9.021 | 15.137 |
| 1179 | 21.459 | 22.576 | 15.044 | 25.244 | 21.142 | 22.259 | 13.817 | 23.185 | 22.238 | 23.354 | 18.042 | 30.274 |
| 1186 | 23.051 | 24.167 | 6.602 | 11.078 | 22.817 | 23.933 | 6.873 | 11.533 | 22.988 | 24.104 | 5.864 | 9.839 |
| 1188 | 19.002 | 20.119 | 10.233 | 17.171 | 19.255 | 20.371 | 10.216 | 17.143 | 19.268 | 20.385 | 10.145 | 17.024 |
| 1189 | 20.590 | 21.706 | 14.973 | 25.124 | 20.649 | 21.765 | 16.189 | 27.165 | 21.965 | 23.081 | 32.134 | 53.921 |
| 1190 | 18.815 | 19.931 | 25.156 | 42.211 | 18.980 | 20.096 | 25.155 | 42.211 | 18.470 | 19.586 | 22.722 | 38.128 |
| 1193 | 19.843 | 20.959 | 13.218 | 22.179 | 19.702 | 20.818 | 12.284 | 20.612 | 20.106 | 21.222 | 13.013 | 21.836 |
| 1200 | 21.596 | 22.712 | 11.485 | 19.272 | 21.468 | 22.584 | 11.912 | 19.989 | 22.460 | 23.576 | 11.330 | 19.011 |


| VCC | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=0) \end{gathered}$ | $(t=0)$ | $\begin{gathered} \hline \hline \mathrm{R}_{\text {eff }}\left[\left[^{\prime \prime}\right]\right. \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2) \end{gathered}$ | $(\mathrm{t}=2)$ | $\begin{gathered} \hline \hline \mathrm{R}_{\mathrm{eff}}\left[\left[^{\prime \prime}\right]\right. \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \hline R_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] | [9] | [10] | [11] | [12] | [13] |
| 1205 | 19.072 | 20.188 | 10.808 | 18.136 | 19.101 | 20.217 | 10.849 | 18.205 | 19.794 | 20.911 | 12.462 | 20.911 |
| 1208 | 20.323 | 21.440 | 4.614 | 7.742 | 20.193 | 21.310 | 4.614 | 7.742 | 20.835 | 21.952 | 4.358 | 7.313 |
| 1217 | 22.632 | 23.748 | 23.831 | 39.988 | 22.568 | 23.684 | 24.210 | 40.625 | 23.037 | 24.153 | 22.300 | 37.420 |
| 1227 | 24.686 | 25.802 | 18.928 | 31.761 | 25.073 | 26.189 | 23.058 | 38.691 | 24.644 | 25.761 | 17.299 | 29.027 |
| 1237 | 20.926 | 22.042 | 6.184 | 10.376 | 20.864 | 21.981 | 6.473 | 10.862 | 21.251 | 22.367 | 6.199 | 10.402 |
| 1257 | 22.033 | 23.150 | 12.557 | 21.071 | 21.840 | 22.956 | 12.874 | 21.602 | 23.056 | 24.172 | 15.190 | 25.489 |
| 1266 | 21.607 | 22.723 | 9.376 | 15.733 | 21.482 | 22.598 | 10.002 | 16.783 | 22.313 | 23.429 | 8.642 | 14.501 |
| 1273 | 20.962 | 22.079 | 10.002 | 16.783 | 21.096 | 22.212 | 10.002 | 16.783 | 21.196 | 22.312 | 9.916 | 16.639 |
| 1290 | 19.186 | 20.302 | 12.523 | 21.014 | 19.263 | 20.379 | 13.096 | 21.975 | 19.424 | 20.540 | 12.235 | 20.529 |
| 1313 | 19.912 | 21.028 | 2.043 | 3.429 | 21.053 | 22.169 | 2.245 | 3.767 | 21.725 | 22.841 | 2.272 | 3.812 |
| 1326 | 19.475 | 20.592 | 15.262 | 25.610 | 19.721 | 20.837 | 15.382 | 25.810 | 19.786 | 20.903 | 15.639 | 26.243 |
| 1330 | 20.242 | 21.358 | 25.060 | 42.051 | 20.359 | 21.476 | 24.802 | 41.618 | 20.237 | 21.353 | 24.038 | 40.336 |
| 1331 | 23.851 | 24.968 | 12.654 | 21.234 | 23.803 | 24.920 | 13.624 | 22.861 | 24.507 | 25.623 | 15.845 | 26.588 |
| 1356 | 20.016 | 21.132 | 5.848 | 9.812 | 19.730 | 20.847 | 5.599 | 9.395 | 21.029 | 22.146 | 6.710 | 11.260 |
| 1374 | 20.293 | 21.409 | 11.768 | 19.746 | 20.847 | 21.963 | 14.243 | 23.901 | 21.690 | 22.806 | 16.580 | 27.821 |
| 1375 | 20.367 | 21.483 | 26.646 | 44.712 | 20.209 | 21.325 | 26.171 | 43.915 | 20.771 | 21.887 | 25.376 | 42.581 |
| 1377 | 22.915 | 24.031 | 20.849 | 34.984 | 22.766 | 23.882 | 21.648 | 36.326 | 23.390 | 24.506 | 19.088 | 32.031 |
| 1379 | 19.907 | 21.024 | 22.283 | 37.391 | 19.777 | 20.894 | 22.283 | 37.391 | 20.231 | 21.347 | 22.442 | 37.657 |
| 1393 | 20.413 | 21.529 | 16.587 | 27.833 | 20.247 | 21.363 | 15.923 | 26.719 | 20.933 | 22.049 | 18.605 | 31.219 |
| 1397 | 21.889 | 23.005 | 3.173 | 5.324 | 21.806 | 22.922 | 3.311 | 5.556 | 22.228 | 23.345 | 2.813 | 4.720 |

continued

| VCC | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=0) \end{gathered}$ | $(\mathrm{t}=0)$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[\left[^{\prime \prime}\right]\right. \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2) \end{gathered}$ | $(\mathrm{t}=2)$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \mu_{\text {eff }} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \mathrm{R}_{\mathrm{eff}}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] | [9] | [10] | [11] | [12] | [13] |
| 1403 | 23.141 | 24.257 | 11.092 | 18.612 | 23.582 | 24.699 | 13.501 | 22.654 | 23.716 | 24.832 | 13.517 | 22.682 |
| 1408 | 23.364 | 24.481 | 8.700 | 14.599 | 23.542 | 24.658 | 9.615 | 16.134 | 23.560 | 24.676 | 9.485 | 15.915 |
| 1410 | 19.924 | 21.041 | 8.983 | 15.074 | 19.794 | 20.911 | 8.983 | 15.074 | 20.259 | 21.375 | 8.650 | 14.515 |
| 1411 | 18.176 | 19.292 | 2.977 | 4.995 | 18.042 | 19.159 | 2.952 | 4.954 | 18.921 | 20.037 | 3.169 | 5.317 |
| 1412 | 18.323 | 19.440 | 18.928 | 31.762 | 18.455 | 19.572 | 18.592 | 31.197 | 18.091 | 19.208 | 18.141 | 30.440 |
| 1413 | 23.324 | 24.440 | 8.557 | 14.358 | 23.265 | 24.381 | 8.498 | 14.259 | 24.432 | 25.548 | 26.908 | 45.151 |
| 1419 | 20.136 | 21.252 | 15.146 | 25.416 | 20.473 | 21.589 | 16.189 | 27.166 | 20.438 | 21.554 | 15.950 | 26.764 |
| 1426 | 22.346 | 23.462 | 13.533 | 22.708 | 22.623 | 23.739 | 16.302 | 27.354 | 22.662 | 23.778 | 15.284 | 25.647 |
| 1427 | 20.984 | 22.100 | 9.586 | 16.086 | 20.763 | 21.879 | 8.681 | 14.567 | 21.433 | 22.550 | 10.040 | 16.847 |
| 1435 | 21.841 | 22.957 | 21.964 | 36.856 | 21.640 | 22.756 | 20.354 | 34.155 | 21.782 | 22.899 | 20.015 | 33.586 |
| 1437 | 18.878 | 19.995 | 3.589 | 6.023 | 18.760 | 19.877 | 3.436 | 5.765 | 19.279 | 20.395 | 3.886 | 6.521 |
| 1442 | 20.664 | 21.780 | 20.256 | 33.989 | 20.430 | 21.547 | 19.125 | 32.092 | 21.122 | 22.239 | 20.878 | 35.034 |
| 1448 | 22.159 | 23.276 | 28.591 | 47.976 | 22.332 | 23.449 | 31.965 | 53.638 | 22.123 | 23.239 | 23.930 | 40.154 |
| 1450 | 20.376 | 21.493 | 18.439 | 30.940 | 20.293 | 21.410 | 18.675 | 31.337 | 21.299 | 22.416 | 21.816 | 36.608 |
| 1455 | 21.875 | 22.991 | 8.703 | 14.603 | 21.665 | 22.781 | 7.964 | 13.364 | 22.556 | 23.672 | 10.440 | 17.518 |
| 1459 | 21.531 | 22.647 | 6.944 | 11.652 | 21.384 | 22.500 | 6.239 | 10.469 | 21.675 | 22.791 | 6.861 | 11.512 |
| 1465 | 22.717 | 23.834 | 16.592 | 27.842 | 22.326 | 23.443 | 16.593 | 27.843 | 23.792 | 24.908 | 16.595 | 27.846 |
| 1468 | 22.252 | 23.368 | 12.029 | 20.185 | 22.061 | 23.177 | 12.295 | 20.631 | 23.170 | 24.287 | 11.950 | 20.052 |
| 1507 | 21.615 | 22.731 | 12.181 | 20.440 | 21.388 | 22.504 | 11.459 | 19.228 | 21.839 | 22.955 | 12.374 | 20.764 |
| 1508 | 19.816 | 20.932 | 18.906 | 31.725 | 19.747 | 20.863 | 19.450 | 32.637 | 20.204 | 21.320 | 17.656 | 29.626 |


| VCC | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=0) \end{gathered}$ | $(t=0)$ | $\begin{gathered} \hline \hline \mathrm{R}_{\text {eff }}\left[\left[^{\prime \prime}\right]\right. \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2) \end{gathered}$ | $(\mathrm{t}=2)$ | $\begin{gathered} \hline \hline \mathrm{R}_{\mathrm{eff}}\left[\left[^{\prime \prime}\right]\right. \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \hline R_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] | [9] | [10] | [11] | [12] | [13] |
| 1515 | 22.538 | 23.655 | 5.037 | 8.453 | 22.353 | 23.469 | 4.831 | 8.106 | 23.946 | 25.062 | 9.581 | 16.076 |
| 1516 | 19.967 | 21.083 | 30.404 | 51.017 | 19.959 | 21.076 | 31.366 | 52.632 | 20.062 | 21.179 | 30.415 | 51.036 |
| 1524 | 21.489 | 22.606 | 24.958 | 41.880 | 21.430 | 22.547 | 26.770 | 44.920 | 22.229 | 23.346 | 24.605 | 41.287 |
| 1529 | 21.700 | 22.816 | 15.284 | 25.647 | 21.492 | 22.609 | 14.560 | 24.431 | 22.259 | 23.376 | 17.672 | 29.653 |
| 1532 | 20.628 | 21.744 | 16.891 | 28.343 | 20.752 | 21.869 | 17.465 | 29.305 | 21.080 | 22.197 | 16.992 | 28.512 |
| 1540 | 18.830 | 19.947 | 39.122 | 65.647 | 19.079 | 20.195 | 39.119 | 65.642 | 18.244 | 19.361 | 33.214 | 55.733 |
| 1552 | 19.632 | 20.749 | 23.021 | 38.630 | 19.766 | 20.882 | 22.721 | 38.125 | 19.762 | 20.878 | 22.669 | 38.038 |
| 1554 | 19.099 | 20.215 | 17.874 | 29.993 | 18.637 | 19.754 | 16.659 | 27.954 | 20.195 | 21.312 | 20.384 | 34.204 |
| 1555 | 20.023 | 21.139 | 54.960 | 92.222 | 20.100 | 21.216 | 57.789 | 96.969 | 20.056 | 21.172 | 50.378 | 84.535 |
| 1557 | 20.135 | 21.251 | 17.878 | 30.000 | 20.005 | 21.121 | 17.878 | 29.999 | 20.455 | 21.571 | 20.139 | 33.794 |
| 1562 | 19.758 | 20.874 | 53.957 | 90.540 | 19.862 | 20.979 | 57.379 | 96.282 | 19.590 | 20.706 | 48.079 | 80.677 |
| 1566 | 21.678 | 22.795 | 14.151 | 23.745 | 21.672 | 22.788 | 15.559 | 26.108 | 22.342 | 23.458 | 17.383 | 29.169 |
| 1569 | 21.907 | 23.023 | 12.823 | 21.517 | 21.681 | 22.797 | 11.929 | 20.017 | 22.688 | 23.804 | 15.671 | 26.296 |
| 1572 | 21.637 | 22.754 | 9.601 | 16.110 | 21.598 | 22.714 | 9.127 | 15.316 | 22.027 | 23.143 | 9.808 | 16.458 |
| 1574 | 20.930 | 22.046 | 3.430 | 5.755 | 21.073 | 22.190 | 3.642 | 6.112 | 20.905 | 22.021 | 3.254 | 5.461 |
| 1575 | 20.260 | 21.376 | 14.429 | 24.212 | 20.545 | 21.661 | 14.796 | 24.827 | 20.698 | 21.815 | 15.504 | 26.015 |
| 1581 | 22.149 | 23.265 | 19.841 | 33.293 | 22.110 | 23.226 | 21.416 | 35.935 | 22.481 | 23.598 | 19.695 | 33.049 |
| 1582 | 23.314 | 24.431 | 11.039 | 18.524 | 23.190 | 24.307 | 10.929 | 18.338 | 23.792 | 24.908 | 11.406 | 19.139 |
| 1585 | 21.804 | 22.920 | 14.821 | 24.870 | 21.655 | 22.772 | 15.423 | 25.879 | 22.540 | 23.656 | 14.383 | 24.134 |
| 1588 | 19.720 | 20.836 | 18.646 | 31.288 | 19.582 | 20.699 | 16.848 | 28.271 | 19.705 | 20.821 | 17.162 | 28.797 |

continued

| VCC [1] | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=0) \\ {[2]} \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=0) \end{gathered}$ <br> [3] | $\begin{gathered} \alpha \\ (\mathrm{t}=0) \end{gathered}$ <br> [4] | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=0) \\ {[5]} \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2) \end{gathered}$ <br> [6] | $\begin{gathered} \mu_{\text {eff }} \\ (\mathrm{t}=2) \end{gathered}$ [7] | $\begin{gathered} \hline \alpha \\ (\mathrm{t}=2) \\ {[8]} \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2) \\ {[9]} \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[10]} \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[11]} \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[12]} \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[13]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1596 | 22.290 | 23.406 | 6.137 | 10.298 | 21.829 | 22.946 | 5.503 | 9.235 | 24.011 | 25.127 | 15.231 | 25.558 |
| 1605 | 22.084 | 23.200 | 10.176 | 17.076 | 21.977 | 23.093 | 10.447 | 17.530 | 22.561 | 23.678 | 10.114 | 16.971 |
| 1615 | 19.555 | 20.671 | 41.796 | 70.133 | 19.688 | 20.804 | 40.798 | 68.458 | 19.472 | 20.588 | 38.910 | 65.292 |
| 1623 | 21.427 | 22.544 | 6.135 | 10.295 | 21.365 | 22.482 | 6.418 | 10.770 | 21.871 | 22.987 | 6.254 | 10.494 |
| 1624 | 19.432 | 20.549 | 13.906 | 23.335 | 19.677 | 20.793 | 13.843 | 23.229 | 19.365 | 20.482 | 12.740 | 21.378 |
| 1644 | 21.857 | 22.973 | 7.813 | 13.110 | 21.773 | 22.890 | 8.244 | 13.834 | 22.676 | 23.792 | 8.201 | 13.761 |
| 1654 | 21.472 | 22.588 | 10.045 | 16.856 | 21.365 | 22.482 | 10.315 | 17.308 | 22.377 | 23.493 | 12.085 | 20.279 |
| 1656 | 22.884 | 24.001 | 11.990 | 20.119 | 23.269 | 24.385 | 13.393 | 22.474 | 23.125 | 24.241 | 11.243 | 18.866 |
| 1675 | 21.387 | 22.504 | 17.205 | 28.869 | 21.299 | 22.415 | 17.478 | 29.329 | 21.675 | 22.792 | 18.687 | 31.356 |
| 1678 | 21.790 | 22.906 | 19.000 | 31.882 | 21.535 | 22.651 | 18.103 | 30.377 | 22.520 | 23.636 | 19.086 | 32.026 |
| 1685 | 21.879 | 22.995 | 15.856 | 26.607 | 21.749 | 22.865 | 15.856 | 26.607 | 22.704 | 23.820 | 15.822 | 26.549 |
| 1686 | 21.032 | 22.148 | 26.213 | 43.986 | 20.902 | 22.018 | 26.213 | 43.985 | 21.752 | 22.868 | 26.563 | 44.572 |
| 1690 | 19.127 | 20.243 | 57.660 | 96.753 | 19.479 | 20.595 | 64.098 | 107.556 | 19.152 | 20.269 | 54.846 | 92.031 |
| 1696 | 20.162 | 21.278 | 31.657 | 53.121 | 20.438 | 21.555 | 34.070 | 57.170 | 20.397 | 21.513 | 32.118 | 53.895 |
| 1699 | 20.829 | 21.945 | 13.771 | 23.108 | 20.588 | 21.705 | 13.376 | 22.445 | 21.777 | 22.893 | 16.497 | 27.682 |
| 1713 | 21.553 | 22.670 | 8.511 | 14.281 | 21.188 | 22.304 | 7.792 | 13.075 | 22.534 | 23.650 | 10.580 | 17.753 |
| 1725 | 21.490 | 22.606 | 16.776 | 28.150 | 20.981 | 22.097 | 14.298 | 23.991 | 22.676 | 23.793 | 23.234 | 38.987 |
| 1726 | 22.188 | 23.304 | 14.686 | 24.643 | 21.600 | 22.717 | 13.019 | 21.847 | 23.556 | 24.673 | 18.608 | 31.224 |
| 1727 | 18.839 | 19.956 | 38.366 | 64.378 | 18.991 | 20.108 | 37.644 | 63.166 | 18.814 | 19.931 | 36.242 | 60.814 |
| 1728 | 22.400 | 23.517 | 8.709 | 14.613 | 22.169 | 23.286 | 7.982 | 13.393 | 23.219 | 24.335 | 10.535 | 17.677 |

[^34]| VCC | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2) \end{gathered}$ | $(\mathrm{t}=2)$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] | [9] | [10] | [11] | [12] | [13] |
| 1730 | 19.379 | 20.495 | 17.712 | 29.721 | 19.591 | 20.707 | 17.998 | 30.201 | 19.487 | 20.603 | 17.673 | 29.656 |
| 1744 | 21.741 | 22.857 | 6.419 | 10.771 | 22.654 | 23.770 | 9.034 | 15.159 | 23.504 | 24.620 | 12.172 | 20.425 |
| 1756 | 22.780 | 23.896 | 6.619 | 11.106 | 22.746 | 23.862 | 6.927 | 11.623 | 22.616 | 23.732 | 6.001 | 10.069 |
| 1757 | 20.010 | 21.126 | 13.479 | 22.617 | 20.296 | 21.413 | 13.724 | 23.029 | 20.363 | 21.479 | 13.962 | 23.428 |
| 1758 | 20.232 | 21.348 | 18.967 | 31.826 | 20.102 | 21.218 | 18.967 | 31.826 | 20.421 | 21.537 | 18.772 | 31.499 |
| 1760 | 19.431 | 20.548 | 21.219 | 35.606 | 19.587 | 20.703 | 21.155 | 35.498 | 19.196 | 20.312 | 19.297 | 32.380 |
| 1776 | 24.639 | 25.755 | 16.550 | 27.771 | 24.857 | 25.973 | 19.786 | 33.201 | 24.409 | 25.525 | 13.565 | 22.762 |
| 1778 | 20.362 | 21.478 | 13.462 | 22.590 | 20.232 | 21.348 | 13.462 | 22.590 | 20.863 | 21.980 | 15.311 | 25.692 |
| 1784 | 23.598 | 24.714 | 16.578 | 27.818 | 23.540 | 24.656 | 16.593 | 27.843 | 24.057 | 25.173 | 17.715 | 29.726 |
| 1789 | 20.868 | 21.985 | 9.408 | 15.786 | 20.806 | 21.922 | 9.827 | 16.490 | 21.139 | 22.255 | 9.924 | 16.652 |
| 1791 | 21.183 | 22.299 | 15.155 | 25.431 | 20.789 | 21.905 | 14.659 | 24.598 | 22.404 | 23.521 | 17.738 | 29.764 |
| 1804 | 21.568 | 22.685 | 11.327 | 19.006 | 21.572 | 22.689 | 11.484 | 19.270 | 21.978 | 23.094 | 12.682 | 21.280 |
| 1811 | 19.363 | 20.479 | 13.298 | 22.314 | 19.233 | 20.349 | 13.298 | 22.314 | 19.820 | 20.937 | 13.756 | 23.082 |
| 1813 | 19.185 | 20.302 | 29.944 | 50.246 | 19.351 | 20.467 | 29.943 | 50.245 | 19.000 | 20.116 | 28.613 | 48.014 |
| 1816 | 22.985 | 24.102 | 15.513 | 26.031 | 22.755 | 23.872 | 14.527 | 24.376 | 24.041 | 25.157 | 18.706 | 31.389 |
| 1821 | 21.597 | 22.714 | 4.584 | 7.691 | 21.534 | 22.651 | 4.768 | 8.000 | 22.426 | 23.542 | 7.210 | 12.099 |
| 1822 | 22.097 | 23.213 | 9.642 | 16.179 | 21.872 | 22.988 | 8.955 | 15.026 | 22.980 | 24.096 | 11.550 | 19.382 |
| 1825 | 21.947 | 23.064 | 8.901 | 14.935 | 22.143 | 23.259 | 8.973 | 15.056 | 22.212 | 23.329 | 8.700 | 14.599 |
| 1837 | 22.453 | 23.569 | 10.959 | 18.388 | 22.705 | 23.821 | 12.563 | 21.080 | 22.456 | 23.573 | 9.605 | 16.118 |
| 1859 | 19.509 | 20.626 | 21.029 | 35.287 | 19.743 | 20.859 | 21.029 | 35.286 | 19.730 | 20.847 | 20.934 | 35.128 |

continued

| VCC [1] | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=0) \\ {[2]} \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=0) \end{gathered}$ <br> [3] | $\begin{gathered} \alpha \\ (\mathrm{t}=0) \\ {[4]} \end{gathered}$ | $\begin{gathered} \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=0) \\ {[5]} \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2) \\ {[6]} \end{gathered}$ | $\begin{gathered} \mu_{\text {eff }} \\ (\mathrm{t}=2) \end{gathered}$ [7] | $\begin{gathered} \alpha \\ (\mathrm{t}=2) \\ {[8]} \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2) \\ {[9]} \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[10]} \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[11]} \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[12]} \end{gathered}$ | $\begin{gathered} \hline \hline \mathrm{R}_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2, \mathrm{tru}) \\ {[13]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1868 | 19.912 | 21.028 | 25.320 | 42.487 | 19.957 | 21.074 | 24.033 | 40.327 | 19.195 | 20.311 | 20.074 | 33.684 |
| 1873 | 21.697 | 22.813 | 7.396 | 12.410 | 21.499 | 22.615 | 7.334 | 12.307 | 22.760 | 23.877 | 9.042 | 15.172 |
| 1883 | 18.392 | 19.509 | 12.532 | 21.028 | 18.539 | 19.655 | 12.361 | 20.742 | 18.485 | 19.601 | 12.178 | 20.434 |
| 1884 | 24.116 | 25.232 | 18.953 | 31.802 | 24.291 | 25.407 | 21.466 | 36.020 | 24.524 | 25.640 | 19.840 | 33.292 |
| 1885 | 22.826 | 23.942 | 15.864 | 26.620 | 22.864 | 23.981 | 15.734 | 26.402 | 22.951 | 24.067 | 15.759 | 26.443 |
| 1898 | 21.651 | 22.767 | 9.816 | 16.471 | 21.886 | 23.003 | 9.745 | 16.352 | 21.894 | 23.010 | 9.616 | 16.135 |
| 1900 | 23.021 | 24.137 | 13.774 | 23.112 | 22.952 | 24.068 | 13.692 | 22.975 | 23.305 | 24.421 | 14.304 | 24.002 |
| 1905 | 25.731 | 26.848 | 25.584 | 42.930 | 25.719 | 26.836 | 24.383 | 40.915 | 26.723 | 27.840 | 135.524 | 227.409 |
| 1918 | 21.986 | 23.102 | 7.760 | 13.021 | 21.879 | 22.996 | 7.964 | 13.364 | 22.892 | 24.008 | 8.325 | 13.970 |
| 1923 | 19.803 | 20.920 | 15.871 | 26.631 | 19.848 | 20.965 | 16.704 | 28.030 | 20.219 | 21.336 | 18.158 | 30.469 |
| 1929 | 20.527 | 21.643 | 20.820 | 34.936 | 20.433 | 21.549 | 19.807 | 33.237 | 21.012 | 22.129 | 21.913 | 36.770 |
| 1931 | 21.976 | 23.093 | 10.731 | 18.006 | 21.580 | 22.697 | 10.000 | 16.779 | 22.968 | 24.085 | 11.989 | 20.118 |
| 1932 | 19.296 | 20.413 | 25.333 | 42.508 | 19.180 | 20.296 | 25.247 | 42.364 | 19.107 | 20.223 | 23.610 | 39.617 |
| 1933 | 21.995 | 23.111 | 6.392 | 10.726 | 21.798 | 22.915 | 6.084 | 10.209 | 23.622 | 24.738 | 11.868 | 19.914 |
| 1943 | 18.785 | 19.901 | 13.682 | 22.958 | 18.998 | 20.114 | 14.842 | 24.906 | 18.770 | 19.887 | 12.649 | 21.225 |
| 1944 | 20.394 | 21.511 | 2.714 | 4.554 | 20.678 | 21.794 | 2.791 | 4.683 | 20.237 | 21.354 | 2.671 | 4.481 |
| 1952 | 22.139 | 23.256 | 9.775 | 16.402 | 21.946 | 23.062 | 10.015 | 16.805 | 23.078 | 24.194 | 10.370 | 17.401 |
| 1960 | 21.305 | 22.421 | 3.064 | 5.141 | 21.090 | 22.206 | 2.987 | 5.013 | 22.738 | 23.854 | 6.175 | 10.361 |
| 1965 | 21.895 | 23.011 | 7.842 | 13.159 | 21.788 | 22.905 | 7.255 | 12.175 | 22.333 | 23.450 | 7.918 | 13.286 |
| 1987 | 19.444 | 20.561 | 34.531 | 57.942 | 19.549 | 20.666 | 36.608 | 61.429 | 19.657 | 20.774 | 34.593 | 58.046 |


| VCC | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=0) \end{gathered}$ | $(t=0)$ | $\begin{gathered} \hline \hline \mathrm{R}_{\text {eff }}\left[\left[^{\prime \prime}\right]\right. \\ (\mathrm{t}=0) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2) \end{gathered}$ | $(\mathrm{t}=2)$ | $\begin{gathered} \hline \hline \mathrm{R}_{\mathrm{eff}}\left[\left[^{\prime \prime}\right]\right. \\ (\mathrm{t}=2) \end{gathered}$ | $\begin{gathered} \mu_{0} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \mu_{\mathrm{eff}} \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \alpha \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ | $\begin{gathered} \hline R_{\text {eff }}\left[{ }^{\prime \prime}\right] \\ (\mathrm{t}=2, \mathrm{tru}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [1] | [2] | [3] | [4] | [5] | [6] | [7] | [8] | [9] | [10] | [11] | [12] | [13] |
| 1992 | 22.347 | 23.463 | 13.099 | 21.981 | 22.029 | 23.145 | 12.320 | 20.674 | 23.395 | 24.511 | 14.546 | 24.409 |
| 1994 | 22.958 | 24.074 | 8.394 | 14.085 | 23.393 | 24.509 | 10.432 | 17.505 | 24.074 | 25.191 | 19.962 | 33.496 |
| 1999 | 18.816 | 19.932 | 10.947 | 18.368 | 18.982 | 20.098 | 11.039 | 18.524 | 19.022 | 20.139 | 11.571 | 19.416 |
| 2006 | 20.634 | 21.750 | 23.173 | 38.884 | 20.535 | 21.652 | 22.874 | 38.383 | 20.808 | 21.925 | 23.890 | 40.088 |
| 2007 | 20.944 | 22.060 | 6.164 | 10.344 | 20.814 | 21.930 | 6.164 | 10.344 | 21.424 | 22.540 | 6.574 | 11.030 |
| 2015 | 21.282 | 22.399 | 6.786 | 11.386 | 21.569 | 22.686 | 7.437 | 12.479 | 22.285 | 23.401 | 9.119 | 15.302 |
| 2023 | 20.432 | 21.548 | 15.358 | 25.772 | 20.163 | 21.279 | 14.100 | 23.660 | 21.089 | 22.205 | 16.254 | 27.274 |
| 2033 | 20.243 | 21.359 | 5.829 | 9.780 | 20.529 | 21.646 | 6.353 | 10.660 | 21.162 | 22.279 | 7.584 | 12.726 |
| 2034 | 21.962 | 23.079 | 10.522 | 17.656 | 21.755 | 22.872 | 10.031 | 16.832 | 22.385 | 23.501 | 10.516 | 17.646 |
| 2037 | 22.731 | 23.848 | 24.428 | 40.990 | 22.601 | 23.718 | 24.427 | 40.989 | 23.398 | 24.514 | 26.513 | 44.488 |
| 2058 | 20.318 | 21.434 | 39.745 | 66.691 | 20.645 | 21.762 | 45.935 | 77.079 | 20.729 | 21.845 | 47.431 | 79.589 |
| 2066 | 19.080 | 20.196 | 20.078 | 33.691 | 19.110 | 20.226 | 18.785 | 31.521 | 19.250 | 20.366 | 19.252 | 32.304 |
| 2070 | 18.764 | 19.880 | 22.826 | 38.303 | 19.002 | 20.118 | 23.702 | 39.772 | 18.728 | 19.845 | 23.363 | 39.203 |
| 2089 | 19.271 | 20.388 | 2.891 | 4.852 | 19.413 | 20.529 | 2.884 | 4.840 | 18.766 | 19.882 | 2.768 | 4.645 |
| 2094 | 23.474 | 24.590 | 6.102 | 10.239 | 23.849 | 24.965 | 7.341 | 12.318 | 23.457 | 24.574 | 5.432 | 9.116 |

Derived parameters of the fits to the SBPs for different evolutionary scenarios.


[^0]:    ${ }^{1}$ Even more confusing: Ferguson and Binggeli (1994) used a limit of $\mathrm{M}_{\mathrm{B}}=-16 \mathrm{mag}$, but also $M_{B}=-18 \mathrm{mag}$, to define dwarf elliptical galaxies.

[^1]:    ${ }^{2}$ parsec $(\mathrm{pc})=3.09 \cdot 10^{16} \mathrm{~m}$ (Jones and Lambourne, 2004).

[^2]:    ${ }^{3}$ Intensities in Fig. 1.5 are given in arbitrary units.

[^3]:    ${ }^{1}$ The X indicates the extreme properties of these BCDs.
    ${ }^{2}$ Note: The LSB-component of BCDs should not be mixed up with the galaxy type commonly called "Low Surface Brightness galaxies". To the contrary, a high central surface brightness $\left(\mu_{0}(B)<22\right.$ $\mathrm{mag} / \mathrm{arcsec}^{2}$ ) has been found to be a characteristic property of the LSB host of BCDs (Papaderos et al., 1996a; Gil de Paz and Madore, 2005) implying a much higher central mass density than in dIs, and the more so, in dSphs and genuine ( $\mu_{0} \geq 23.5 \mathrm{mag} / \mathrm{arcsec}^{2}$ ) LSB galaxies.
    ${ }^{3}$ The Holmberg radius defines the radius, where $\mu(\mathrm{B})=26.5 \mathrm{mag} / \mathrm{arcsec}^{2}$.

[^4]:    ${ }^{4}$ Using the online Image List Tool of the SDSS, http://skyserver.sdss3.org/dr8/en/tools/chart/list.asp.

[^5]:    ${ }^{5}$ Please do not mix up $\eta$ of the Sérsic law with the $\eta$ of the $\eta$-function.

[^6]:    ${ }^{6}$ Note that the $d I s$ of our sample and the Ims of the VCC are handled as the same morphological class of dwarf irregular galaxies and we therefore use both notations.

[^7]:    ${ }^{7}{ }^{V_{L G}}$ defines the velocities of the galaxies with respect to the local group (LG) and was calculated via $v_{\text {LG }}=v_{\text {helio }}+220 \mathrm{~km} / \mathrm{s}$.

[^8]:    ${ }^{8}$ It should be pointed out that the findings of Drinkwater and Hardy (1991) particular based on a sample of the most compact BCDs in Virgo.

[^9]:    ${ }^{9}$ Results of uncertain morphological types are taken from Section 2.4.8.

[^10]:    ${ }^{10}$ European Extremely Large Telescope (e.g. Kissler-Patig et al., 2009)
    ${ }^{11}$ James Webb Space Telescope (e.g. Gardner et al., 2006)
    ${ }^{12}$ Next Generation Virgo Cluster Survey (e.g. Ferrarese et al., 2012)

[^11]:    ${ }^{1}$ Also the bulges show a variety in morphology: classical, boxy and peanut bulges (e.g. Athanassoula, 2005).

[^12]:    ${ }^{2}$ Objects with brighter magnitudes than the dwarf galaxies.
    ${ }^{3}$ Please note that the classes of dI and Im will treated equally in this study.

[^13]:    ${ }^{4}$ Web-page of the GOLDMine data base: http://goldmine.mib.infn.it/
    ${ }^{5}$ In an updated version of the GOLDMine data base (check: 04.10.2012), VCC-0323 has a velocity of $\mathrm{v}_{\text {helio }}=2756 \mathrm{~km} / \mathrm{s}$.

[^14]:    ${ }^{6}$ For a better contrast we recommend the digital version of this thesis.

[^15]:    ${ }^{7}$ http://iraf.noao.edu/

[^16]:    ${ }^{8}$ The colours are not affected by different distances, since we are still at redshifts close to zero.

[^17]:    ${ }^{9}$ Citation from Sandage and Brucato (1979).

[^18]:    ${ }^{10}$ The GOLDMine distances were only applied to the late-type galaxies and not to the early-types.

[^19]:    ${ }^{1}$ Of course, not only massive stars, but also low- and intermediate mass stars are formed (see initial mass function in Section 4.2.3).

[^20]:    ${ }^{2}$ Negative values of $\tau$ will lead to increasing SFRs.

[^21]:    ${ }^{3}$ The $\chi^{2}$ algorithm of GAZELLE is slightly modified from the classical definition.

[^22]:    ${ }^{4}$ Even the age estimates of the most metal-poor BCD I Zw18 changes over the last years, making it older with every new and deeper observation (e.g. Hunt et al., 2003; Aloisi et al., 2007).

[^23]:    ${ }^{5}$ Note that the mean effective surface brightness $\langle\mu\rangle_{\mathrm{eff,r}}$ and the effective surface brightness at $\mathrm{R}_{\mathrm{eff}}$ are different quantities.

[^24]:    ${ }^{6}$ In this context misclassified means that the classification by blue sensitive photo-plates differs from the classification by our parameter-space.

[^25]:    ${ }^{7}$ Except for a face-on galaxy, the inclination angle between the galaxy and ICM only plays a minor role (Quilis et al., 2000; Marcolini et al., 2003).

[^26]:    ${ }^{8}$ Citation from Thuan (1985).

[^27]:    ${ }^{9}$ The spiral arm on the left is bluer than the right one (see lower panel of Fig. 2.16)

[^28]:    ${ }^{10}$ Abbreviation for "Wisconsin Indiana Yale NOAO" telescope.

[^29]:    ${ }^{11}$ http://www.wiyn.org/Observe/wiynfilters.html.
    ${ }^{12}$ http://www.noao.edu/kpno/manuals/whirc/filters.html.

[^30]:    continued

[^31]:    Results from GALEV/GAZELLE runs (Part II)

[^32]:    continued

[^33]:    continued

[^34]:    continued

